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Title: 2020 Virtual Theoretical Division Lightning Talk Series

Author(s): Ramsey, Marilyn Leann; Albright, Apollo James Yin; Andrews, Elizabeth Marie; Csicsery-Ronay, Jennifer Ann; Darwin, Leah Jewell; Fedik, Nikita Sergeyevich; Forde, Aaron Arthur; George, Andrew; Gonzales, Theodore Gary; Kolade, Beauty Ayomide; Kulichenko, Maksim; Kupets, William Redpath; Mallela, Abhishek; Mora Perez, Carlos Adrian; Morency, Marie Yamilee Natalie; Patrick, Amanda C.; Sakano, Michael Neil; Sud, James Patrick; Tkachenko, Nikolay; Wang, Ying; Wu, Shanshan; et al.

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## 2020 Virtual Theoretical Division Lightning Talk Series

Marilyn L. Ramsey

Albright, A.  
Andrews, E.  
Csicsery-Ronay, J.  
Darwin, L.  
Fedik, N.  
Forde, A.  
George, A.  
Gonzales, T.  
Kupets, W.  
Kolade, B.  
Kulichenko, M.  
Kupets, W.  
Mallela, A.  
Mora Perez, C.  
Morency, Y.  
Patrick, A.  
Sakano, M.  
Sud, J.  
Tkachenko, N.  
Wang, Y.  
Wu, S.

## 2020 Virtual Theoretical Division Lightning Talk Series

LANL Host: Marilyn L. Ramsey

**WebEx Meeting Link:** <https://lanl-us.webex.com/lanl-us/j.php?MTID=me40cde05abed7de3e8413cc6fb64af13>

**WebEx Meeting ID:** 133 665 4525

**WebEx Password:** F4NiNEn443P

All members of the T Division Community, students, staff members, group leaders, division management and other interested individuals are **invited to come and support** the following student(s) during their Lightning Talk presentation.

**\*The Lightning Talk Series is a forum for T Division students to present information about their current projects**

**Tuesday, July 14, 2020, 10:00 a.m.–11:00 a.m.**

**Webex Meeting ID:** 133 854 0103

**Webex Password:** \_A3Ayv3cr2G

**Student Speakers:** James Sud, UGS

**School Affiliation:** University of California, Berkeley

**Group:** T-4, Physics of Condensed Matter and Complex Systems

**Mentors:** Lukasz Cincio (T-4), Yu Zhang (T-1), Pavel Dub (C-IIAC)

**Student Speakers:** Nikolay Tkachenko, GRA

**School Affiliation:** Utah State University

**Group:** T-4, Physics of Condensed Matter and Complex Systems

**Mentors:** Lukasz Cincio (T-4), Yu Zhang (T-1), Pavel Dub (C-IIAC)

**Title:** *Reducing Complexity of Variational Quantum Eigensolver for Quantum Chemistry*

**Abstract:** The solution of electronic structure problem is thought to be the nearest future implementation of small, noisy quantum computers. While the complexity of the exact solution of electronic structure problems on classical computers grows exponentially with the number of electrons, we can use the advantage of quantum calculations to make the problem complexity polynomial with respect to the number of simulated particles. Due to the noise inherent in quantum computers, one approach is to use variational quantum-classical algorithms (VQCs), although a concern for VQCs is the large number of measurements needed for convergence, especially to within chemical accuracy. Our group uses two techniques to reduce the number of measurements. The first is Rosalin, a shot-frugal algorithm that measures qubits with probabilities based on the coefficients of corresponding terms in the Hamiltonian. The second is permuting qubits in the fermionic Hamiltonian to make the qubit assignment more appropriate for a given architecture and/or ansatz function, which also decreases the required circuit depth. In our study we attempt to combine the powers of these two methods to push VQE algorithms to molecules previously incalculable due to exceedingly high measurement counts and circuit depths.

**Tuesday, August 4, 2020, 10:00 a.m –11:00 a.m** **Student Speaker:** Andrew George, GRA

**Group:** T-5, Material and Physical Data

**School affiliation:** University of Houston, PhD. Student

**Mentors:** Nathan Garland(T-5) and Mark Zammit(T-1)

**Title:** *Uncertainty Quantification of Atomic Data for Collisional-radiative Modeling of Fusion Plasmas*

**Abstract:** Knowing plasma properties such as ion populations and radiated power are crucial in both standard, steady-state and non-standard, instability or disruption scenarios of tokamak plasmas. A superconfiguration collisional-radiative model which is predicated on simple fit cross sections can be used to describe radiative processes and also allow effective completeness in state space. The approximate cross sections have an

inherit uncertainty due to the approximations utilized in cross section formulas. Dakota, an uncertainty quantification package can assess these uncertainties by performing a range of UQ studies on multiple collision processes for various atomic targets. The results of the analysis will allow us to understand and approximate CR models by comparing simple fit cross section formulas to the ones derived from fine structure models.

**Student Speaker:** Willem Kupets\*, GRA

**Group:** \*T-1: Physics and Chemistry of Materials and T-5: Materials and Physical Data

**School Affiliation:** St. John's College Santa Fe

**Mentors:** Mark Zammit\*, Nathan Garland^, James Colgan

**Title:** *Electron Stopping Powers and Mean Excitation Energy of Atoms and Ions using the Los Alamos Atomic Physics Suite of Codes*

**Abstract:** Electron-impact excitation and ionization cross sections of atoms and ions are crucial for modeling plasmas that are in non local thermodynamic equilibrium during tokamak disruption events. Here we utilize cross sections calculated from the Los Alamos Atomic Physics Suite of Codes to calculate electron stopping powers and mean excitation energies for a variety of atoms and ions. We also utilize the Bragg's additive rule to approximate stopping power of molecules. We validate the calculated stopping powers against the convergent close-coupling data for H and H<sub>2</sub> before being integrated into the Los Alamos Suite of Codes.

**Student Speaker:** Leah Darwin, UGS

**Group:** T-6, Theoretical Biology and Biophysics

**School Affiliation:** Arizona State University

**Mentors:** Chris Neale

**Title:** *Computational Search for Cancer Treatments*

**Abstract:** Cancer is a class of diseases where cells grow uncontrollably and form tumors. Untreated tumors can be lethal, making cancer a leading cause of death worldwide. Up to 30% of all human cancers are found to have a mutation in a protein involved in cell growth known as Ras. Mutations to Ras can cause it to be stuck in an on state, where it continuously signals that more cells need to be made, resulting in uncontrollable cell growth. My research focuses on repairing the mechanism that allows Ras to switch to an off state using new drugs. I use molecular dynamics simulations to emulate drugs interacting with Ras and a well-studied model system. After I optimize our system, I will search large chemical libraries to find drug candidates that will be tested at Frederick National Lab to determine their potential to become new treatments for cancer.

**Student Speaker:** Amanda Patrick, GRA

**Group:** T-6, Theoretical Biology and Biophysics

**School Affiliation:** University of Texas at Arlington, PhD. Student

**Mentors:** Nick Hengartner and Imelda Trejo

**Title:** *Nonparametric estimation of force of infection in a generalized SEIR model*

**Abstract:** The rate at which susceptible individuals become infected is known as the rate of infection. The rate of infection will be assumed as a function of time. The foundation of this estimation will be explored using a generalized SEIR model with known distributions for time to recovery and exposure time. Coronavirus data will be applied to the system to obtain numerical results. How the rate of infection can be estimated will be explored.

Thursday, August 6, 2020, 10:00 a.m. –11:00 a.m

**Student Speaker: Carlos Mora Perez, GRA**  
**School affiliation: University of Southern California (USC) Department of Chemistry**  
**Group: T-1, Physics and Chemistry of Materials**  
**Mentor(s): Amanda Neukirch & Dibyajyoti Ghosh**

**Title:** TDDFT Characterization of a CsPbBr<sub>3</sub> Clusters' Optical Properties

**Abstract:** CsPbBr<sub>3</sub> perovskite quantum dots (QDs) have recently become promising candidates for photovoltaic applications such as light-emitting diodes (LEDs) but the nature of their optical properties have yet to be fully characterized. In this work, we utilize Time-Dependent Density Functional Theory (TDDFT) within the Gaussian16 package to characterize the excitation states transition energies of a small CsPbBr<sub>3</sub> cluster. Several charged vacancy defects and a charge build-up in the system were modeled to gain insight into their perturbative effect on electronic properties. The simulated UV-Vis spectrums outline the extent to which Br vacancies and charge systems (+1 or -1) lower absorption, and why they should be avoided for high performing QDs. The findings provide a guide for the tunability of CsPbBr<sub>3</sub> QDs optical properties by understanding the nature of the excitations.

**Student Speaker: Apollo Albright, UGS**  
**School Affiliation: Reed College**  
**Group: T-5, Applied Mathematics and Plasma Physics**  
**Mentor(s): Nathan Lemons**

**Student Speaker: Theodore Gonzales**  
**School Affiliation: University of Colorado-Boulder**  
**Group: T-5, Applied Mathematics and Plasma Physics**  
**Mentor(s): Nathan Lemons**

**Title:** *Detecting Anomalies in Graphs*

**Abstract:** We present a graph-based mathematical framework to do high dimensional, combinatorial hypothesis testing. The modeling framework has many potential applications, including water testing, outbreak detection, anomalous computer detection, and more. Within this framework, we review three specific tests over lattices that generalize to other graphs.

**Thursday, August 13, 2020, 10:00 a.m.–11:00 a.m.**

**Student Speaker: Jennifer A. Csicsery-Ronay, UGS**  
**School Affiliation: University of Houston, School of Natural Sciences & Mathematics**  
**Group: T-6, Theoretical Biology and Biophysics**  
**Mentor: William S. Hlavacek**

**Title:** *de novo* Assembly of a Reference Genome Sequence for *Pterophyllum* spp. & the Search for Evolutionary Insights

**Abstract:** *Pterophyllum* is the genus of fish commonly known as angelfish, an iconic member of the aquatic hobbyist's aquarium. By generating a *Pterophyllum* reference-based genome sequence and subsequently the sequences of additional wild and domestic individuals of varying populations, we will identify how the genome structure of the domestic angelfish is affected by anthropogenic hybridization, as it is believed to be a hybrid of closely related species such as *P. scalare* and *P. leopoldi*. Comparisons of the genomes will lead to identifying the original sources of the domestic angelfish in the wild, as well as pondering a deeper understanding of how hybridization affects its morphological features.

**Student Speaker:** Beauty Kolade, UGS  
**School affiliation:** CUNY, Lehman College  
**Group:** T-CNLS, Center for Non-Linear Studies and B-11  
**Mentor(s):** Jacob Miner & Robert F. Williams

**Title: Validating Experimental Results of toxins using Cheminformatics and Quantum Chemistry**

**Abstract:** Cardiac glycosides are classes of drugs used in the treatment of heart failure. Developing experimental and computational techniques for predicting the range of conformers that can be adopted by these molecules can provide valuable insight into their chemical potency. Previous studies, including efforts by our group, have focused on using ion-mobility mass spectrometry (IM-MS) and molecular dynamics (MD) to describe arrival times average collisional cross sections (CCS) values associated with many different classes of biomolecules, including cardiac glycoside such as Digitoxin. Our goal in this project is to develop a new computational pipeline for generating conformers for validation against experimental results. Using cheminformatics software (RDKit) and open-source quantum chemistry software (psi4) we are able to generate ensembles of digitoxin, analyze results, and generate CCS values more rapidly than we could previously with molecular dynamics. Using geometric criteria of these structures, and those of simple buckyballs and nanotubes, we are able to identify the most salient geometric features that affect calculated CCS values. We aim to apply this approach to other cardiac glycosides so that we can create a library of ion-mobility-based data for cardiac glycosides that other scientists can use to carry out toxicological investigations of unknown compounds.

**Student Speaker:** Ying Wang, GRA  
**School affiliation:** University of Southern California  
**Group:** T-1, Physics and Chemistry of Materials  
**Mentor:** Amanda Neukirch

**Title: Electron-Hole recombination of Hybrid Dion-Jacobson 2D Lead Iodide Perovskite: Nonadiabatic Molecular Dynamics Analysis**

**Abstract:** 2D hybrid organic-inorganic halide perovskites attracts increasing research attention because of the higher moisture stability and tunabilities of electronic and optical properties. Recently, Dion-Jacobson 2D lead iodide perovskites have been applied to solar cells, exhibiting considerable application prospects. In this project, we focus on two basic hybrid Dion-Jacobson 2D lead iodide perovskite structures,

(3AMP)(MA)<sub>3</sub>Pb<sub>4</sub>I<sub>13</sub> and (4AMP)(MA)<sub>3</sub>Pb<sub>4</sub>I<sub>13</sub> with power conversion efficiency (PCE) as solar cells, 7%, 5%, respectively. Experimental work has indicated that the less distorted inorganic lattice

in (3AMP)(MA)<sub>3</sub>Pb<sub>4</sub>I<sub>13</sub> results in a narrower band gap. The shortest I···I interlayer distance among reported 2D lead iodide perovskites leads to the enhancement in charge transport performance. In order to investigate optoelectrical behavior of these materials and improve the PCE, we adopt ab initio calculation and nonadiabatic molecular dynamics simulation. Through ab initio calculation, we investigate the relationship between optoelectric properties and structures of these materials in ground state, especially the high carrier mobility along the stacking direction. The comparison between radiative lifetimes and nonradiative lifetimes of these perovskites rationalizes their PCE difference. Our study helps us have a better understanding of 2D hybrid organic-inorganic halide perovskites from a structure to property aspect and guide the direction for solar cells design with improved PCE.

**Student Speaker:** Shanshan Wu, GRA  
**School affiliation:** University of Illinois at Chicago  
**Group name:** T-1, Physics and Chemistry of Materials.  
**Mentors:** Danny Perez, Ping Yang

**Title: Long-timescale simulations of morphological changes in ThO<sub>2</sub> nanoparticles**

**Abstract:** Thorium dioxide (ThO<sub>2</sub>) nanomaterials have gained increasing attention due to their various application, including catalysis and ceramics. As early actinide oxide, ThO<sub>2</sub> is also of interest for nuclear applications, in and of itself, or as a surrogate for later actinides. Understanding the morphological and dynamical properties of ThO<sub>2</sub> is therefore very important. Molecular dynamics (MD) simulations are an ideal tool to do so, as they provide fully-resolved atomistic trajectories. My research focuses on studying morphological changes of ThO<sub>2</sub> nanoparticles using Accelerated MD, a class of methods that significantly extend the time scale of MD. In this talk, I will first introduce the method I used, and then discuss the morphological and dynamical properties of very small ThO<sub>2</sub> nanoparticles, such as those that would form soon after nucleation. In particular, I will show that some nanoparticles experience very slow shape transitions between different classes of structures.

**Tuesday, August 18, 2020, 10:00 a.m.–11:00 a.m.**

**Student Speaker: Yamilee Morency, GRA**

**School Affiliation: University of Pennsylvania**

**Group: T-1, Physics and Chemistry of Materials**

**Mentor: Justin Smith**

**Title: *Physics-based Machine Learning Potential Report Card***

**Abstract:** Potential energy surfaces are essential to solving many problems in chemistry. They are however limited by costly ab initio methods. Machine learning potentials provide a cost-efficient alternative to ab initio methods with little to no loss in accuracy. However, these models have no physical basis. To evaluate their ability to reproduce complex physical phenomena, we propose a physics-based machine learning report card. This report card will be a test suite made up of archetypal problems in chemistry it will also attempt to answer the question of whether failure of the model is caused by the poor parameterization or model descriptors. This talk is a first step towards this report card. It is an analysis of the short-range dependency of the necessary cutoff radius of ANI models in terms of conjugated hydrocarbons.

**Student Speaker: Nikita Fedik, GRA**

**Group: T-1, Physics and Chemistry of Materials**

**School: Utah State University, Chemistry and Biochemistry Department**

**Mentors: Benjamin Nebgen, Sergei Tretiak**

**Title: *Machine learning parametrization of empirical and semiempirical methods for improving transition metal chemistry***

**Abstract:** Transition metals (TM) and their complexes are of paramount importance for organometallics, environmental study, catalysis, and industrial applications. However, large scale modeling of TM complexes is very limited by the performance of contemporary methods such as Density Functional Theory (DFT). Fast semiempirical quantum mechanical methods (SEQM), which are of two order of magnitude faster than DFT, often produce errors about 100 kcal/mol for simple organometallic complexes. Additionally, many such SEQM methods completely lack parameters for TM, sometimes neglecting d-orbitals entirely. In this talk, we aim to develop classic machine learning (ML) model for transition metals using the HIPNN architecture. Additionally, we attempt to generate a SEQM parameterization scheme utilizing ML dynamic parameterization to simultaneously improve the accuracy of SEQM methods and rapidly parameterize previously unstudied transition metal.

**Student Speaker:** Maksim Kulichenko  
**School Affiliation:** Utah State University  
**Group:** T-1, Physics and Chemistry of Materials  
**Mentors:** Justin Smith, Sergei Tretiak

**Title:** *Diversification of ML Datasets Using Bias Potential as Function of “Uncertainty”*

**Abstract:** One of the central issues in machine learning is a dataset. Its diversity, size, and preparation time are the variables which researchers always aim to optimize. The purpose of this work is to capture as much chemical space as possible using active learning and biased potentials embedded in the molecular dynamics. Active learning is the technique which allows to keep the dataset as diverse as possible using the ensemble of neural networks. Within MD framework this means augmentation of dataset with only those MD steps where ensemble of neural networks feels uncertain. By constructing biased potentials which are functions of the “uncertainty”, we want to make poorly sampled regions more energetically favorable. Thus, artificially forcing the MD to avoid well sampled regions, bias potentials within active learning may significantly accelerate datasets preparation time while increasing the data diversity.

**Student Speaker:** Yamilee Morency, GRA  
**School Affiliation:** University of Pennsylvania  
**Group:** T-1, Physics and Chemistry of Materials  
**Mentor:** Justin Smith

**Title:** *Physics-based Machine Learning Potential Report Card*

**Abstract:** Potential energy surfaces are essential to solving many problems in chemistry. They are however limited by costly ab initio methods. Machine learning potentials provide a cost-efficient alternative to ab initio methods with little to no loss in accuracy. However, these models have no physical basis. To evaluate their ability to reproduce complex physical phenomena, we propose a physics-based machine learning report card. This report card will be a test suite made up of archetypal problems in chemistry it will also attempt to answer the question of whether failure of the model is caused by the poor parameterization or model descriptors. This talk is a first step towards this report card. It is an analysis of the short-range dependency of the necessary cutoff radius of ANI models in terms of conjugated hydrocarbons.

**Wednesday, August 19, 2020, 10:00 a.m. –11:00 a.m.**

**Student Speaker:** Elizabeth Andrews, GRA  
**Group:** T-CNLS, Center for Non-Linear Studies and EES-16  
**School affiliation:** PhD Candidate at Colorado School of Mines  
**Mentors:** Jeffrey Hyman, Matthew Sweeney, David Moulton

**Title:** *How many years can a mountain exist before it's washed to the sea: Investigating weathering of fractured rocks*

**Abstract:** Weathering plays a critical role in our daily lives and the global climate cycle. It provides the soils which we depend on to grow food, and over longer timescales, it consumes atmospheric CO<sub>2</sub> which regulates the Earth's climate. Many studies have been completed to understand the rate that weathering occurs at the Earth's surface, including field studies and laboratory studies. However, there is a consistent discrepancy between rates measured in the lab and those measured in the field, sometimes the difference is several orders of magnitude. If the rates measured in the lab were widespread in field systems, we might expect that mountains would have fully weathered away. Despite many attempts to reconcile this difference in rates, we still do not fully understand the mechanisms that create the observed large discrepancy. This study investigates how heterogeneous flow paths in fractured bedrock control the weathering rates observed at the hillslope scale using reactive transport simulations. Fractures are modeled discretely using dfnWorks, coupled

to PFLOTRAN to simulate weathering in 3D domains. We will use these simulations to understand the spatial distribution of weathering rates in fractured rock systems. We will also determine how fracture density impacts measured weathering rates at the hillslope scale.

**Student Speaker:** Michael Sakano, GRA  
**Group:** T-1, Physics and Chemistry of Materials  
**School Affiliation:** Purdue University, PhD Student  
**Mentors:** Edward Kober

**Title:** *Comparing RDX and HMX reduced chemistry models applied to 1-D hot spot simulations*

**Abstract:** Both shock-to-detonation and deflagration-to-detonation transitions are important phenomena in the field of high-energy materials research as their behaviors are still not well understood after decades of experimental and theoretical work. Previous efforts focused on connecting gas-phase energetics with shock and thermal experiments to develop simplified, yet plausible chemistry models. However, it is relatively unknown how well these reaction schemes generalize the overall chemical decomposition process, or if they are even applicable to condensed-phase materials. In this study, we refine our preliminary “reduced order” chemistry models for condensed-phase RDX and HMX single crystals obtained from reactive molecular dynamics simulation with the ReaxFF-Ig force field. Data was collected from various small-scale temperature and pressure homogeneous isothermal simulations and these results were analyzed using the non-negative matrix factorization algorithm. The computed chemistry components, which were in good agreement with published literature, added greater detail and insight over prevailing theories. We mapped these components onto larger 1-D hot spot simulations to understand how the system evolves from reactants to gaseous products both spatially and temporally. We observed trends in the relationship between variations in chemistry and loading conditions, and eventually plan to implement these simplified models into continuum-level simulation codes.

**Student Speaker:** Abhishek Mallela  
**School Affiliation:** Department of Mathematics, University of California Davis  
**Group name:** T-6, Theoretical Biology and Biophysics  
**Mentor name:** William S. Hlavacek

**Title:** *Bayesian Uncertainty Quantification for COVID-19 Situational Awareness*

**Abstract:** To increase situational awareness and support evidence-based policy-making, we formulated two types of mathematical models for COVID-19 transmission within a regional population. One is a fitting function that can be calibrated to reproduce an epidemic curve with two timescales (e.g., fast growth and slow decay). The other is a compartmental model that accounts for quarantine, self-isolation, social distancing, a non-exponentially distributed incubation period, asymptomatic individuals, and mild and severe forms of symptomatic disease. Using Bayesian inference, we have been calibrating our models daily for consistency with new reports of confirmed cases from the 15 most populous metropolitan statistical areas in the United States and quantifying uncertainty in parameter estimates and predictions of future case reports. This online learning approach allows for early identification of new trends despite considerable variability in case reporting. We infer new significant upward trends for five of the metropolitan areas starting between 19-April-2020 and 12-June-2020.

**Thursday, August 27, 2020, 10:00 a.m. –11:00 a.m.**

**Student Speaker:** Aaron Forde (GRA)  
**School Affiliation:** Materials and Nanotechnology, North Dakota State University  
**Group name:** T-1, Physics and Chemistry of Materials  
**Mentor:** Amanda J. Neukirch, Dibya Ghosh



**Title: *Chiroptical Properties Induced into a Lead-Halide Perovskite Cluster Through Surface Chemistry***

**Abstract:** Chiroptical properties of materials is of interest for various applications, such as structure determination, photo-detectors for polarized light, highly polarized photoluminescence, and spintronics. Chiroptical activity of organic molecules has been a long studied subject due the enantiomer structures which break mirror symmetry, allowing for distinct magnetic-dipole interactions. Inducing chiroptical activity into semiconductor materials has been a challenge due to the difficulty of creating crystal structures that break mirror symmetry. One proposed method to induce chirality into semiconductor materials is to use chiral organic molecules as capping ligands for inorganic nanocrystals. Experimentally it is observed that chiral capped nanocrystals show circular dichroism and polarized photoluminescence, but the atomistic details as to how these molecules impart chiral-optical signatures is not well understood. Is it suspected that the chiral signatures are activated from chiral templating of the nanocrystal surface or through charge-transfer excitations. Here we implement atomistic modeling using time-dependent density functional theory calculations to explore the photo-physical pathways that give rise to circular dichroism for a  $\text{CsPbCl}_3$  lead-halide perovskite cluster that is passivated with chiral methylbenzylamine and diaminocyclohexane molecules. We find that when the chiral molecules are anchored to the cluster surface there are chiroptical signatures which are distinct from the molecule and cluster themselves. We also find the chirpticalo signatures to be contributions of charge-transfer excitations from the molecule to the cluster and from cluster-to-cluster transitions. This suggests there are competing photo-physical pathways giving rise to circular dichroism in chiral passivated clusters.

# Detecting Anomalies in Graphs

Los Alamos National Laboratory

Mentor: Nathan Lemons

Apollo Albright<sup>1</sup>    Teddy Gonzales<sup>2</sup>

<sup>1</sup>Reed College

<sup>2</sup>University of Colorado Boulder

August 06, 2020

# Can You Find the Path of Elevated Mean?

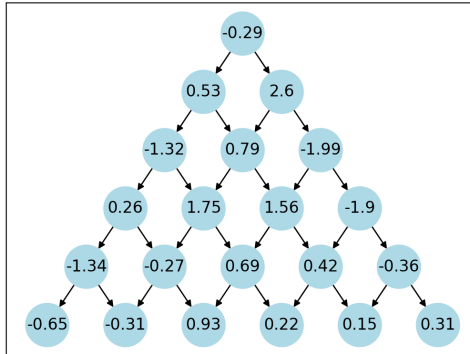


Figure: Lattice With Path of Elevated Mean

# Can You Find the Path of Elevated Mean?

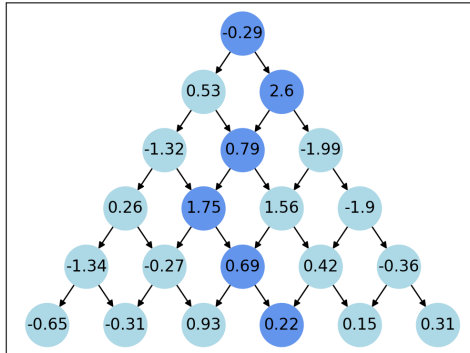


Figure: Lattice With Path of Elevated Mean

# Introduction To Framework: Motivation

## Potential Applications:

- Chemical spill
- Image patterns
- Anomalous Computers
- Disease outbreak
- Anomalous event in temporal data

# Introduction To Framework: Water Quality

1. Sensors are placed at each branching in a river system.
2. At each node  $i$ , a sensor detects the concentration  $X_i$  of an unwanted chemical  $X$ .
3. Using this data, a test or series of tests determines if a path of elevated chemical concentration exists in the river network.

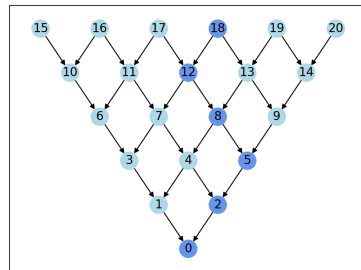


Figure: Example of a Contaminated Stream

# Introduction to Framework: In General

- Given  $\mathbf{X} \subset \{0, 1, 2, \dots, n-1\}^m$  with some graph structure and with each component a random variable, we would like to distinguish between [1]:
  - $H_0$  : Every  $x \in \mathbf{X}$  comes from  $N(0, 1)$
  - $H_1$  : Some subset  $P \subseteq \mathbf{X}$  from a class  $\mathcal{C}$  of subsets has an elevated mean  $N(\mu, 1)$
- As  $n \rightarrow \infty$ , we would like to find a threshold for  $\mu$ , above which, some test  $T$  can distinguish between  $H_0$  and  $H_1$ , and below which, no test or sequence of tests can distinguish  $H_0$  and  $H_1$ .

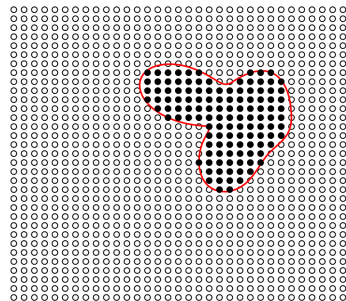


Figure: More General Cluster

# Introduction To Framework: Notation

- From here on, we focus on the case of a two-dimensional lattice.
- The lattice has depth (number of levels)  $m$ .
- We are interested in  $P \in \mathcal{P}_m$ , where  $\mathcal{P}_m$  is all paths of length  $m$  within the lattice.
- A test  $T : \mathbf{X} \rightarrow \{0, 1\}$  distinguishes between the presence or absence of a path with elevated mean.
  - $H_0$  : Every  $x \in \mathbf{X} \sim N(0, 1)$
  - $H_1$  : A path  $P \in \mathcal{P}_m \sim N(\mu, 1)$ ; all  $x \notin P \sim N(0, 1)$

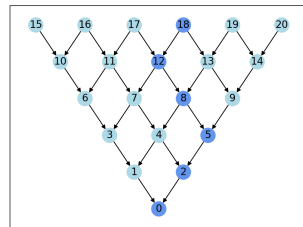


Figure: Lattice with  $m = 6$



# Three Questions

- What values of  $\mu$  allow us to distinguish between  $H_0$  and  $H_1$ ?
- Which tests can work for this?
- What are the computational complexities of these tests?

# Types of Tests: Averaging Test

- The Averaging Test is perhaps the simplest [1]:
  - Reject  $H_0$  if  $\sum_{x \in \mathbf{x}} \geq \frac{1}{2} \mu m$
- Benefits:
  - The Averaging Test is computationally very fast, making it the easiest to use on very large graphs.
  - Can be used to detect any pattern in the graph, not just paths.
- Drawbacks:
  - Its detection threshold is far from optimal, and the value of  $\mu$  required to detect anomalies can grow quite quickly with some graph types.
  - Cannot find the exact location of a path, only its presence.

# Types of Tests: Strip Test

- Strip Test [3]:
  - Reject  $H_0$  if  $\sum_{v \in V} X_v > \frac{1}{2}\mu m$  with  $V$  the set of nodes which are within the strip of radius  $2\sqrt{m} + 1$  around the root node.
- Benefits:
  - Checks significantly fewer nodes ( $\approx 4m^{3/2}$ ) than the averaging test ( $\approx m^2/2$ ).
- Drawbacks:
  - Same drawbacks as the averaging test.
  - Limited to 2 Dimensions
  - Poor performance if path is outside strip, susceptible to adversarial attacks.

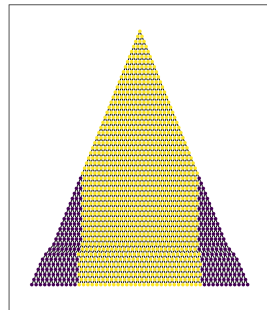


Figure: Strip Over a Lattice

# Types of Tests: Bayes Test

- Bayes Test [3]:
  - Assume paths are generated according to some prior distribution  $\pi$
  - Reject  $H_0$  if  $L_m(X) = P(H_1|\pi)/P(H_0|\pi) > 1$ .
- Benefits:
  - The Bayes Test can be computed relatively quickly using a recursive algorithm on any acyclic directed graph.
  - If the distribution of paths is known, it has an optimal detection threshold for  $\mu$  [3].
- Drawbacks:
  - Detecting paths of unknown length can be slow.
  - It assumes paths are generated according to a given prior distribution and is susceptible to adversarial attacks.

# Types of Tests: GLRT

- Generalized Likelihood Ratio Test [3]:
  - Reject  $H_0$  if  $\max_{p \in \mathcal{P}_m} (\sum_{i \in p} X_i) \geq m\sqrt{2 \log(2)}$
- Benefits:
  - Can be applied to find more general shapes in many graphs
  - Path distribution not required
  - Can be used to determine path location
- Drawbacks:
  - Detecting paths of unknown length can be slow.
  - Threshold is not clear when generalizing to other graph and path structures.
  - Often computationally difficult to scan over all “windows”.

# Comparison of Tests over 2-D Lattices

- Because the value of  $\mu$  required to separate  $H_0$  and  $H_1$  can change based on  $m$ , we define  $\mu_m$  to be the value of  $\mu$  for a given  $m$ .

Test	Separable	Inseparable	Runtime	$\mu_{0.95}$
Averaging	$\mu \rightarrow \infty$	$\mu \rightarrow 0$	$O(m^2)$	2.4
Strip	$\mu_m m^{1/4} \rightarrow \infty$	$\mu_m m^{1/4} \rightarrow 0$	$O(m^{3/2})$	2.25
GLRT	N/A	N/A	$O(m^2)$	1.25
Bayes	$\mu_m m^{1/4} \rightarrow \infty$	$\mu_m m^{1/4} \rightarrow 0$	$O(m^2)$	0.85

**Table:** Bounds for separability and inseparability, runtime scaling as a function of  $m$ , and the value of  $\mu$  required for 95% accuracy with  $m = 50$ .

# Comparison of Tests Over Lattices



**Figure:** Success Rate vs  $\mu$  for Various Tests,  $m = 100$

- In our testing, with probability 0.5, a path  $p$  was picked uniformly at random.
- Under these circumstances, the Bayes test performs the best, followed by the GLRT.
- The Strip and Averaging Tests perform very similarly.

# More General Methods

- Generalizations:
  - These methods apply to more than just a 2-D lattice: random geometric graphs, binary trees, and higher dimensional lattices [3, 2].
  - Similar approaches can be used if the random weights are not normally distributed [3].
  - We worked on a few methods of detecting paths of unknown length.
  - These methods, particularly the GLRT, can be applied to detect more general clusters [2].
  - However, without some form of structure to the data, it remains impossible or impractical to detect such clusters as the size of the data set increases.



# Thank You for Your Time!

Thank you to our mentor Nathan Lemons who has helped us with this project throughout the summer!

# References



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*Annals of Statistics - ANN STATIST*, 39, 01 2010.



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Searching for a trail of evidence in a maze.

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# HOW MANY YEARS CAN A MOUNTAIN EXIST BEFORE IT'S WASHED TO THE SEA: INVESTIGATING WEATHERING OF FRACTURED ROCK SYSTEMS

ELIZABETH ANDREWS<sup>1,2,3</sup>, MATT  
SWEENEY<sup>2</sup>, DAVID MOULTON<sup>4</sup>,  
ALEXIS NAVARRE-SITCHLER<sup>1</sup>,  
JEFFREY HYMAN<sup>2</sup>

1. COLORADO SCHOOL OF MINES

2. EES-16

3. CNLS

4. T-5





# WHAT IS WEATHERING?

## Chemical - dissolution



## Physical – breaking of rock

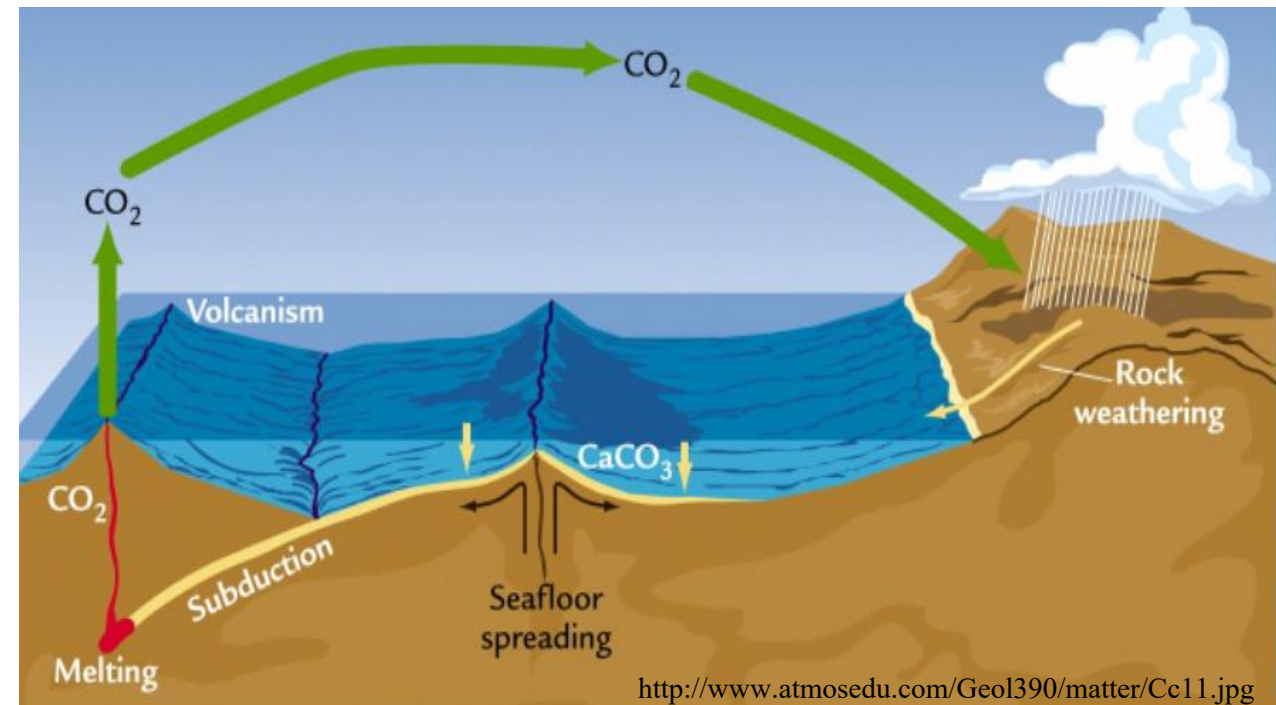




# WHY IS WEATHERING IMPORTANT?

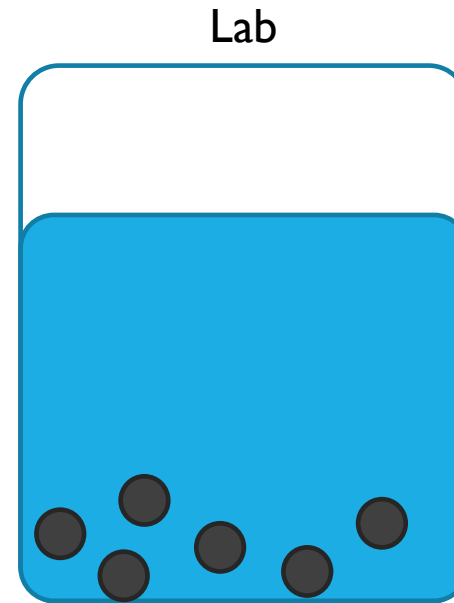


- Agriculture depends on soil
- Silicate weathering is a long-term climate regulator



# WEATHERING RATES MEASURED IN THE LAB ARE SIGNIFICANTLY FASTER THAN FIELD RATES

- Lab conditions?
  - High water:rock ratio?
  - Solution far from equilibrium?
  - Too much surface area?
- Or, is it because of the heterogeneity in the natural system?
  - Heterogeneous flow paths through fractured rock?



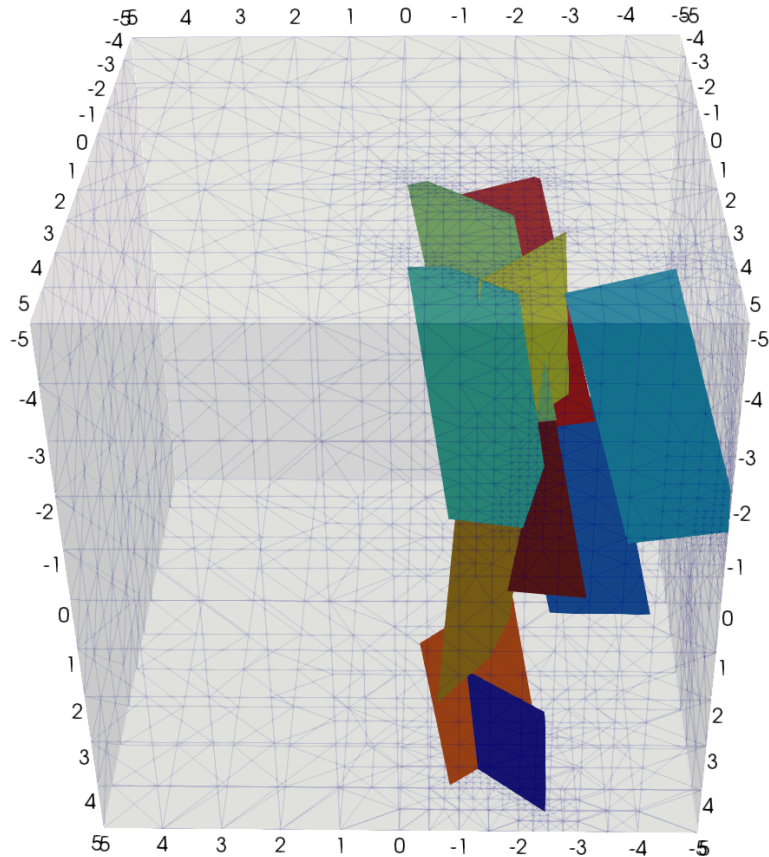


## RESEARCH QUESTIONS:

Can flow through fractures explain the difference between lab rates and field rates?

How does fracture density impact the dissolution rate?

# DOMAIN SETUP

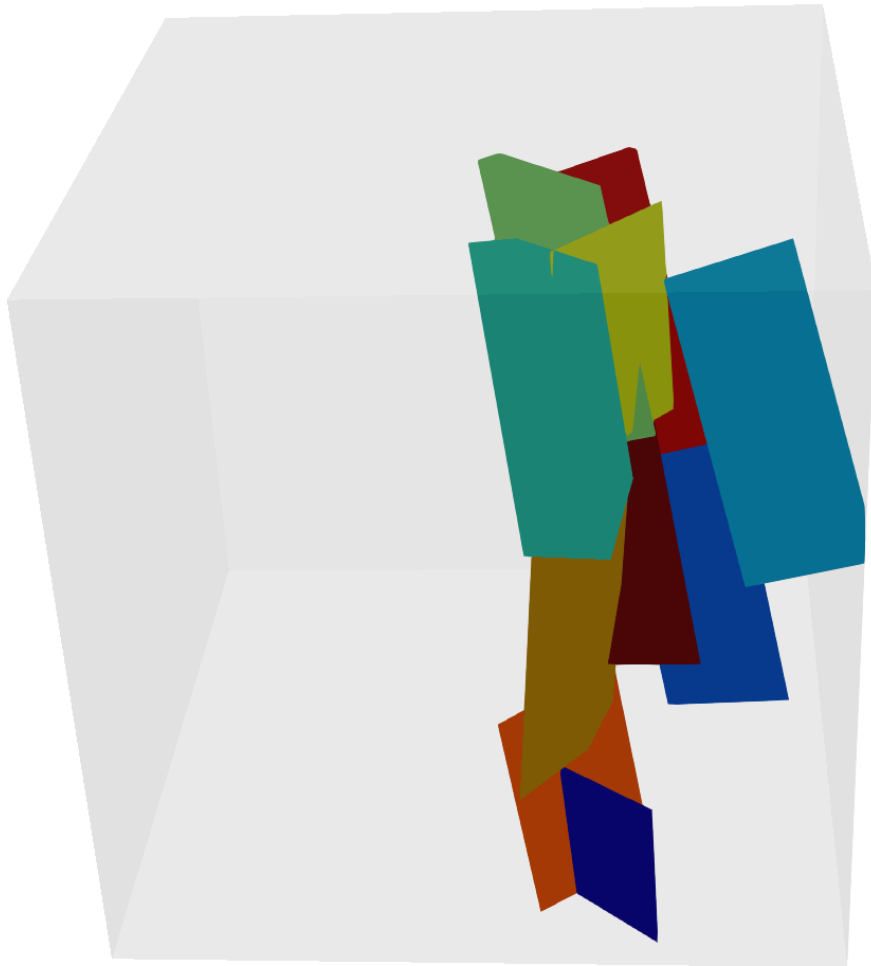


- dfnWorks coupled to PFLOTRAN
- Octree refined domain
- Fluid enters and exits the domain through fractures
- Reactions: Albite dissolution, kaolinite precipitation
  - $\text{Albite (NaAlSi}_3\text{O}_8) + 4\text{H}^+ \rightarrow \text{Na}^+ + \text{Al}^{3+} + 3\text{SiO}_2 + 2\text{H}_2\text{O}$
  - $2\text{Al}^{3+} + 2\text{SiO}_2 + 5\text{H}_2\text{O} \rightarrow \text{Kaolinite (Al}_2\text{Si}_2\text{O}_5(\text{OH})_4) + 6\text{H}^+$

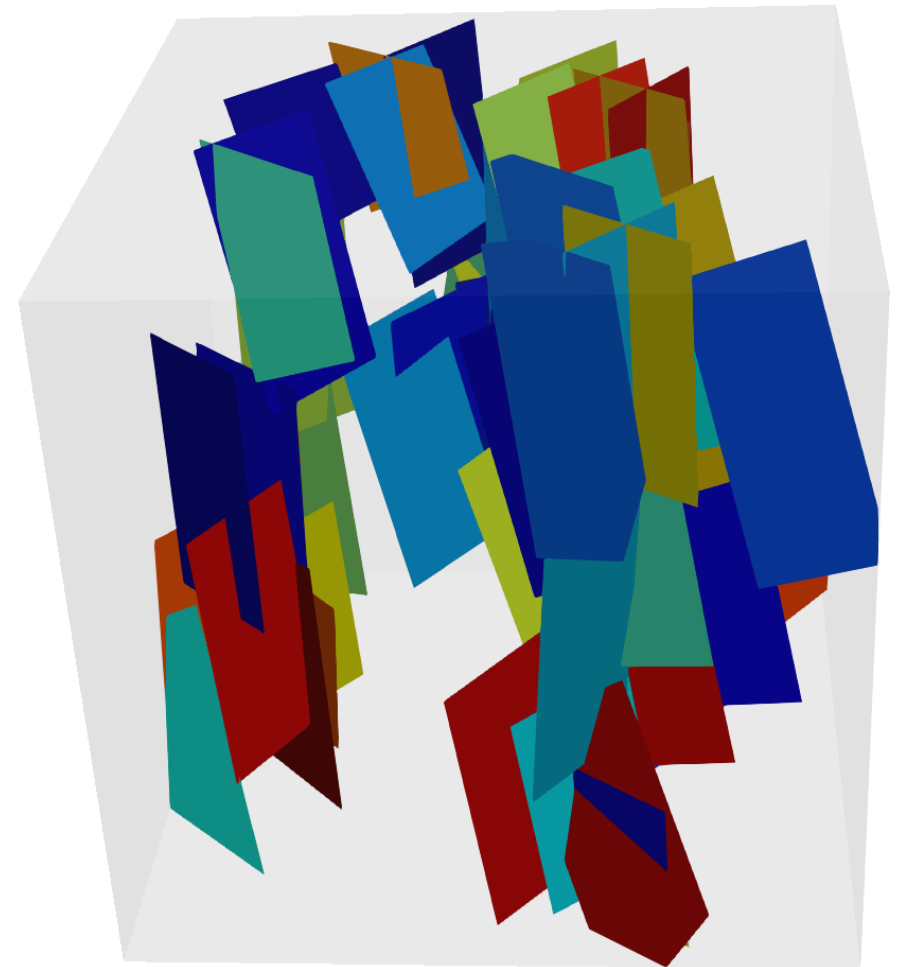


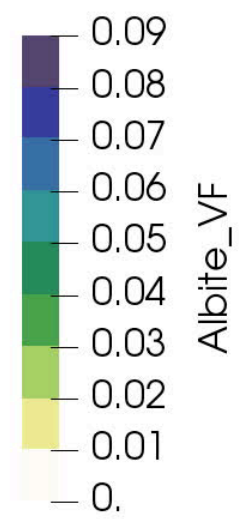
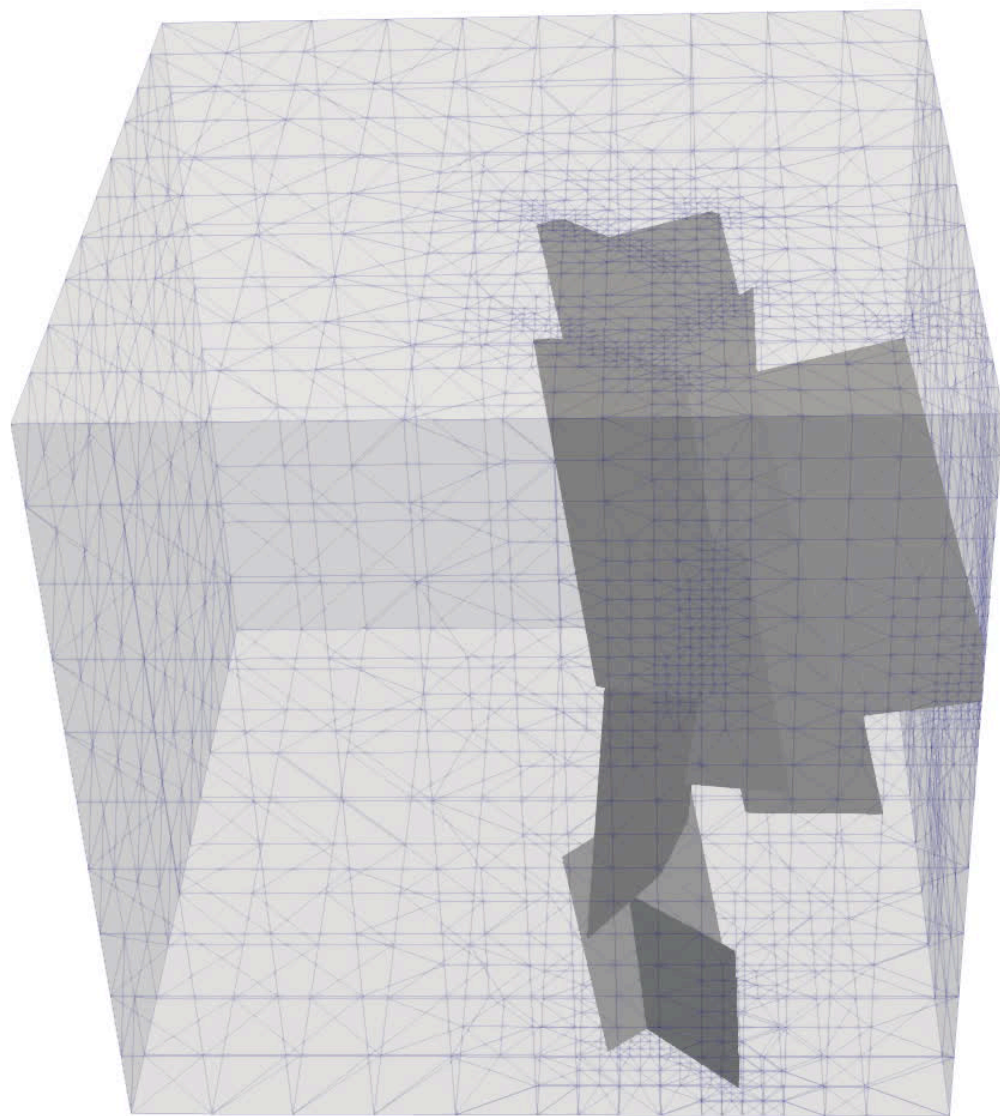
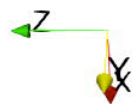
# SIMULATIONS: 2 DOMAINS

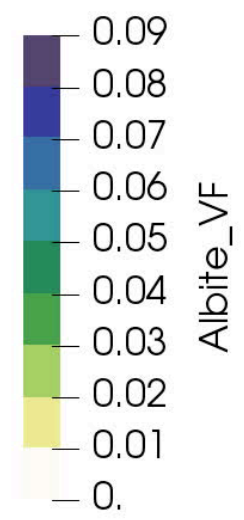
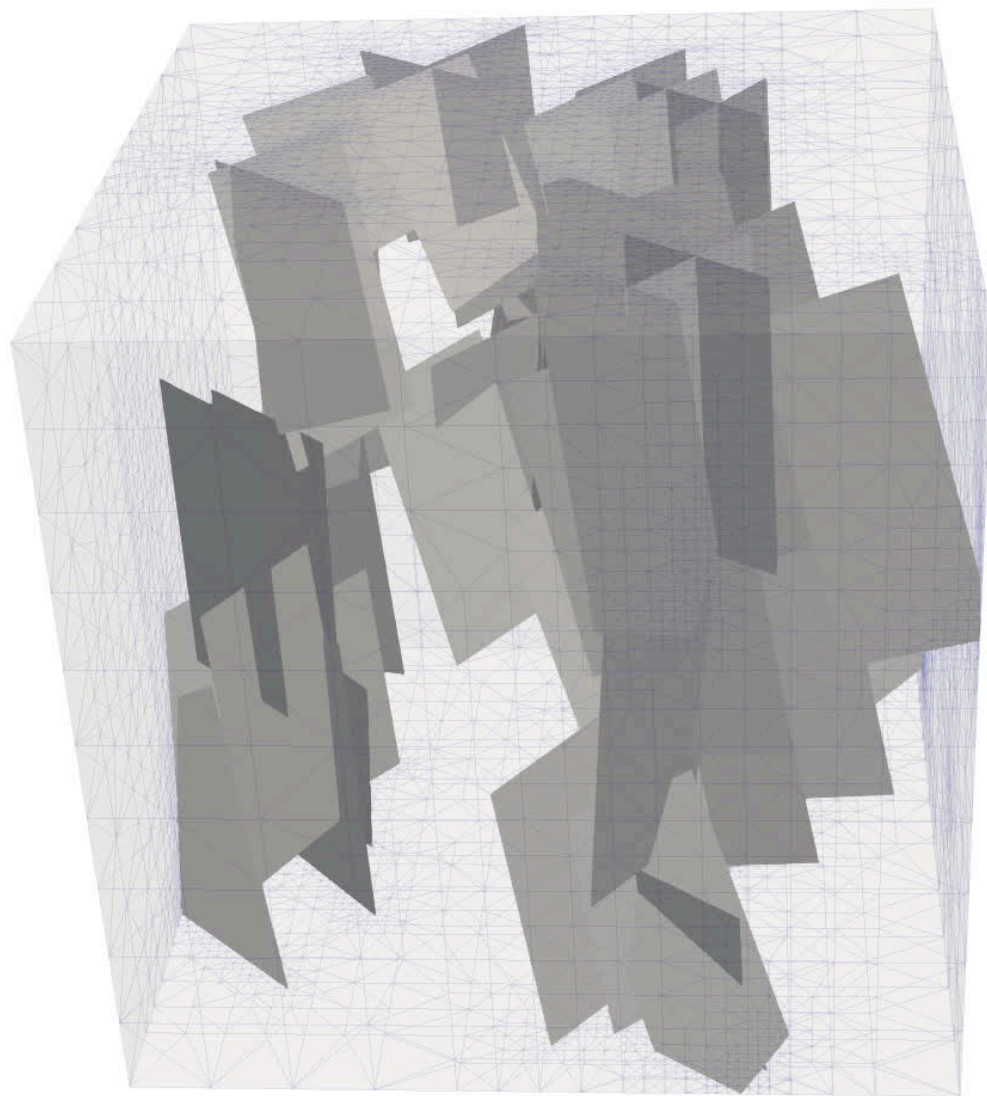
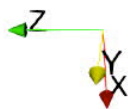
Low density: 11 fractures



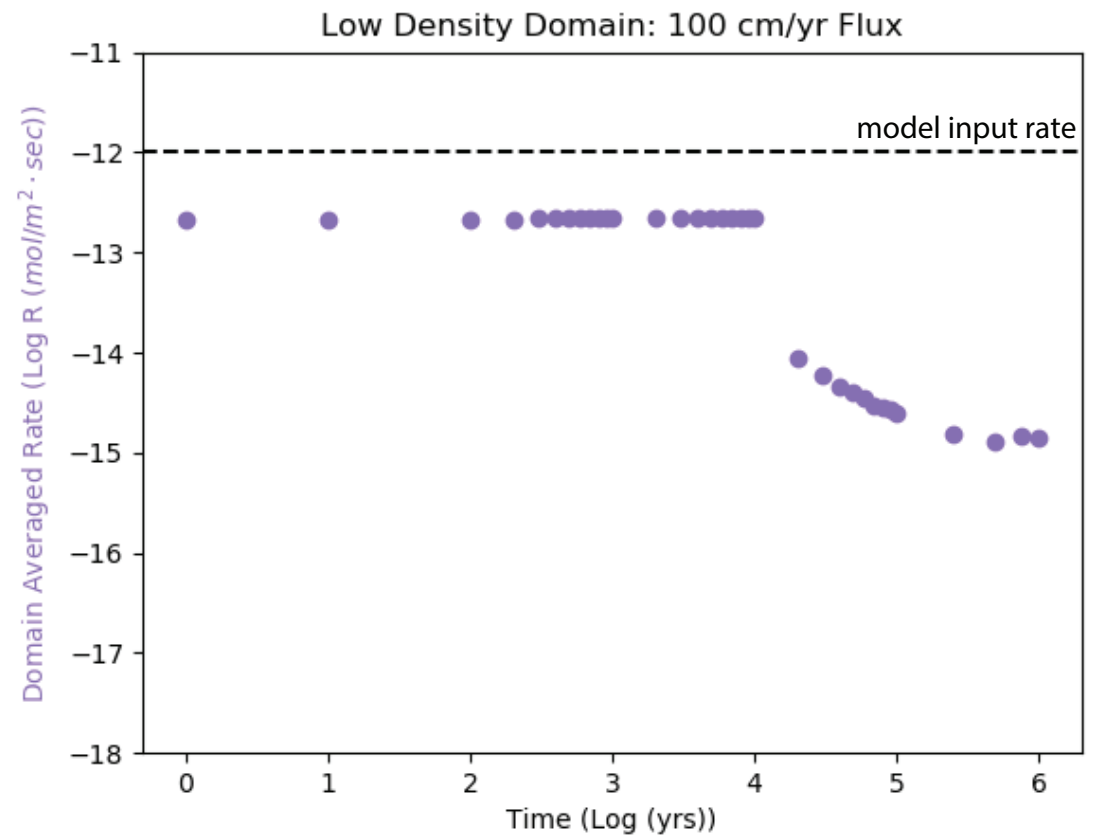
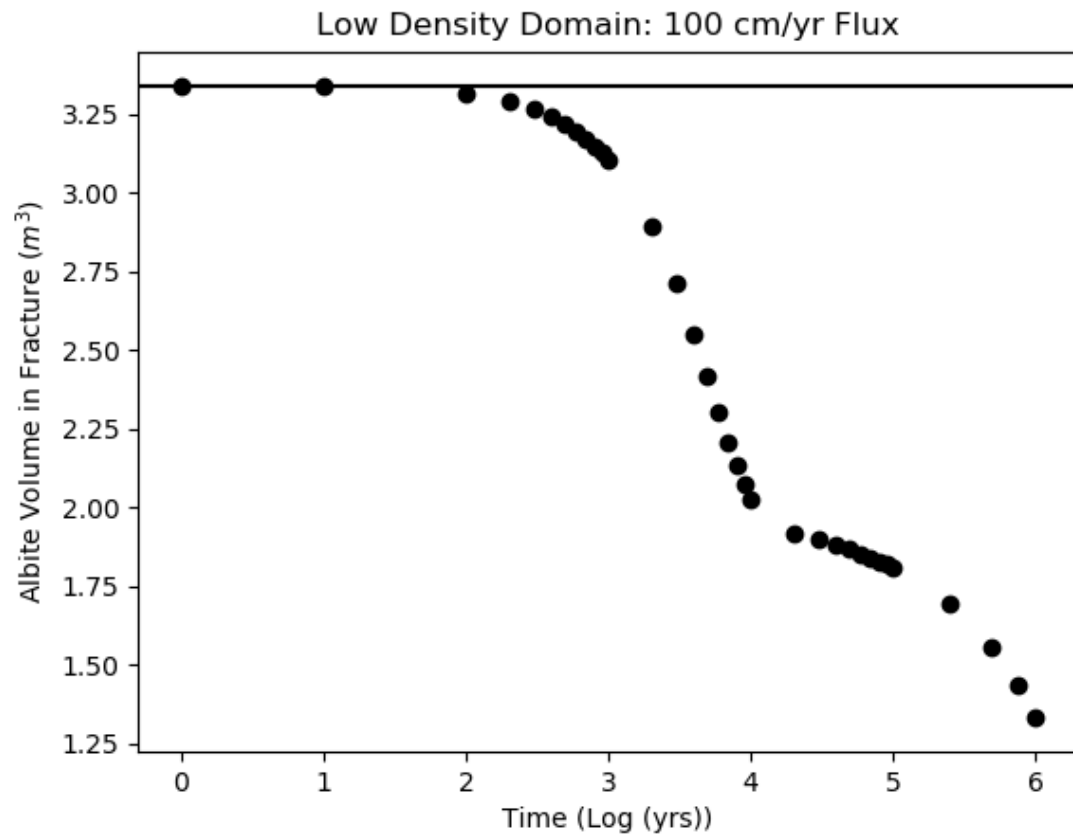
High density: 71 fractures



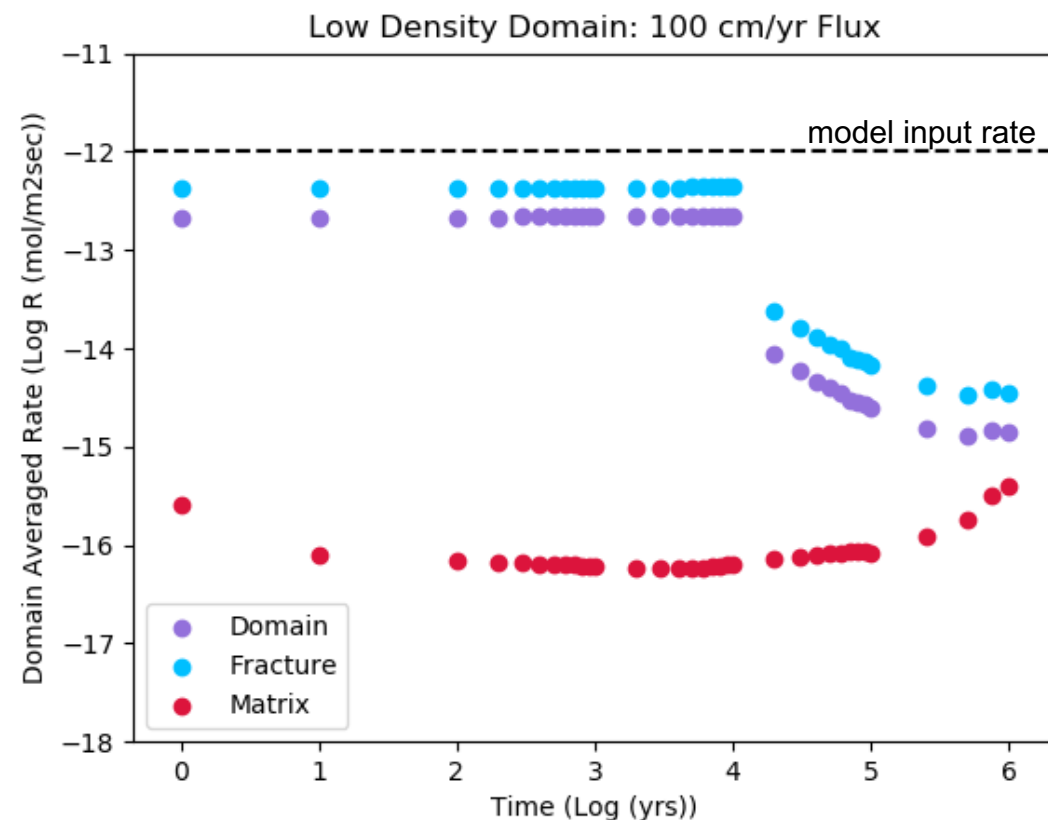
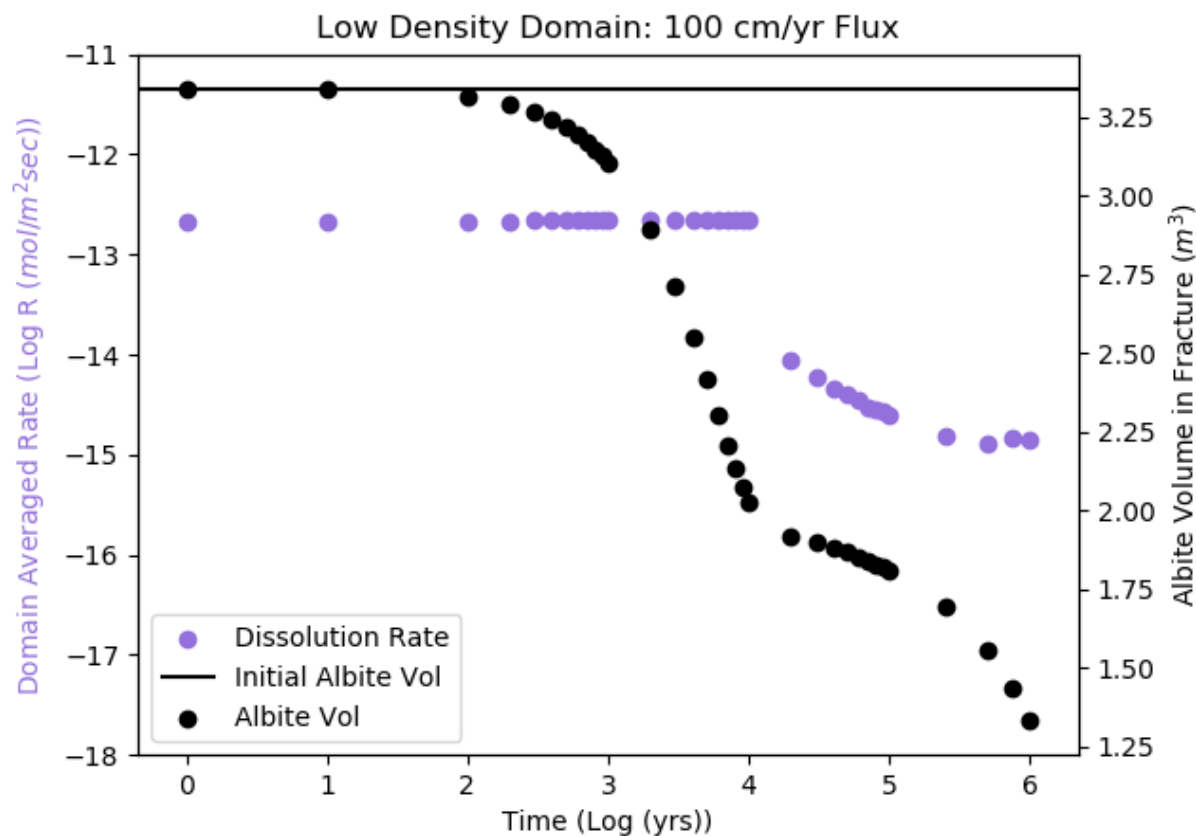




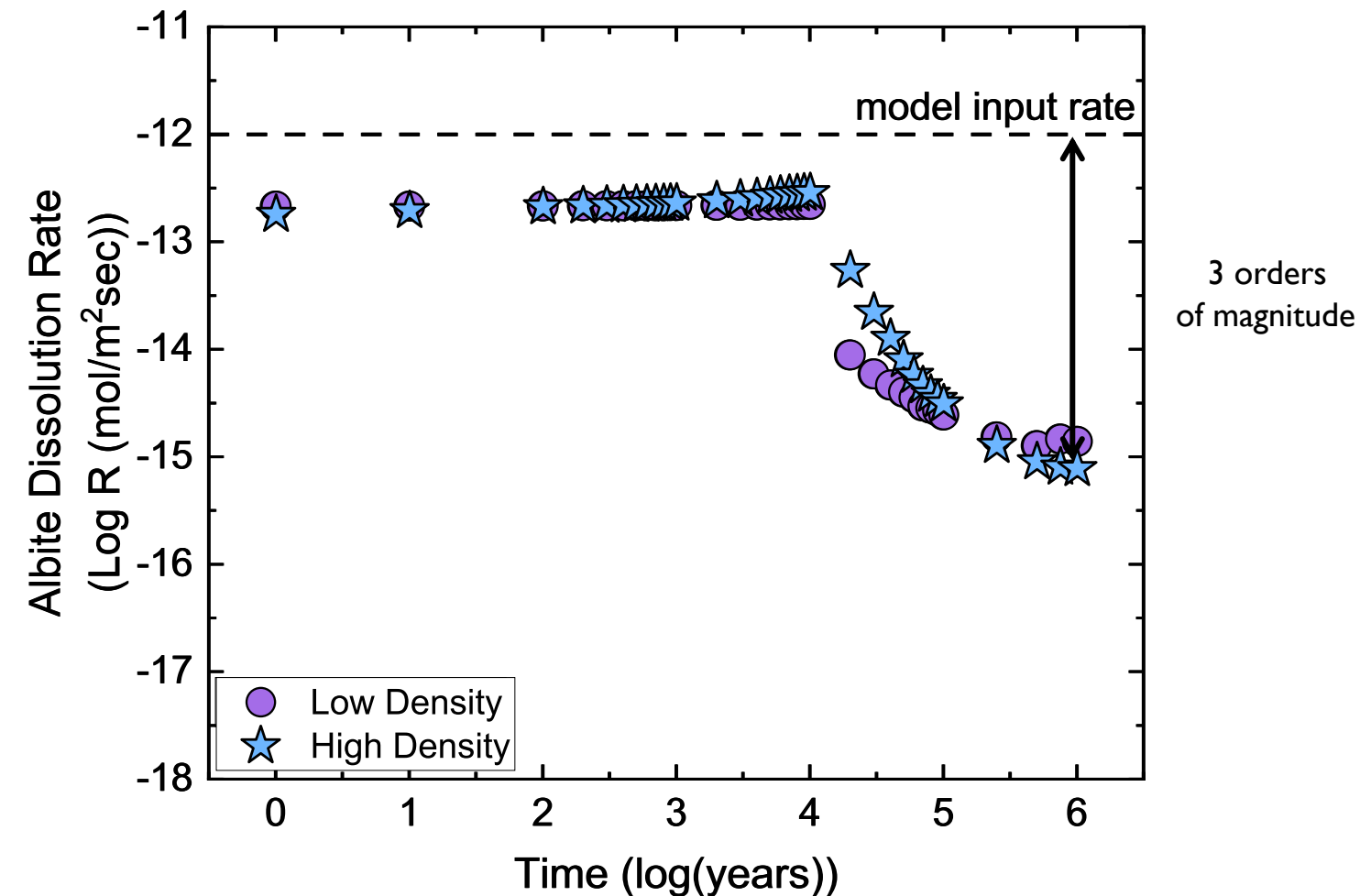
# RESULTS



# RESULTS: FAST WEATHERING IS DOMINATED BY REACTION IN FAST FLOWING FRACTURES

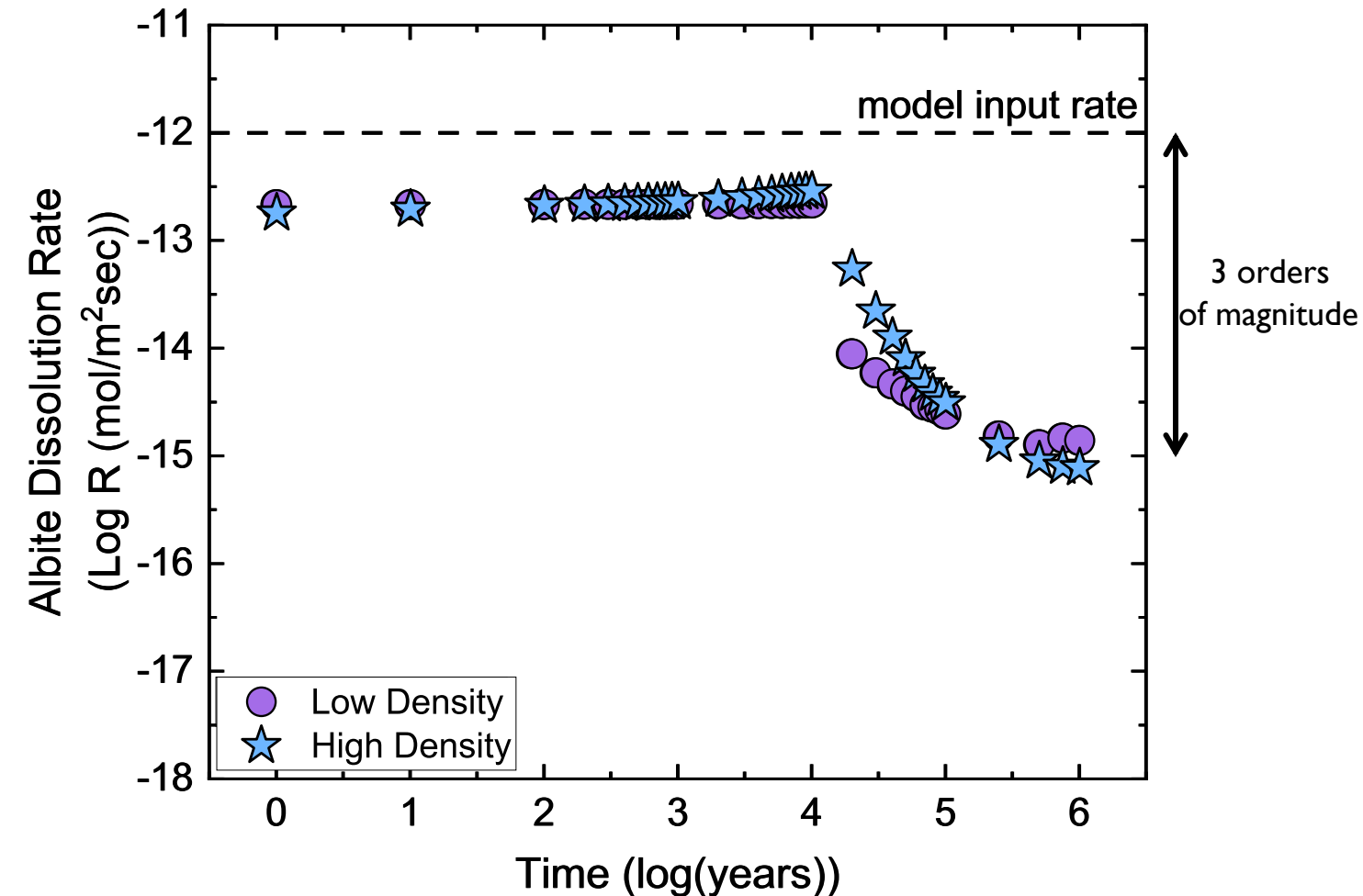


# DISCUSSION



- At early times, no difference in rates between fracture density variations
- At late times, rate behavior varies
- After 100,000 years, the output weathering rate is 3 orders of magnitude smaller than the input rate
- The processes associated with flow through fractures could potentially explain the lab vs. field rate conundrum

# THANKS! QUESTIONS?



- At early times, no difference in rates between fracture density variations
- At late times, rate behavior varies
- After 100,000 years, the output weathering rate is 3 orders of magnitude smaller than the input rate
- The processes associated with flow through fractures could potentially explain the lab vs. field rate conundrum

*de novo* Assembly of a  
Reference Genome Sequence for  
*Pterophyllum*, spp.

...

Jennifer A. Csicsery-Ronay  
University of Houston  
Theoretical Biology & Biophysics  
Center for Nonlinear Studies  
Mentor: William S. Hlavacek



# Outline

- Specimen selection, Tissue Collection, and DNA Extraction
- Hybrid *de novo* assembly
- Quality Assessment of the Draft Genome
- Strategy for Evolutionary Analysis

# Specimen Selection, Tissue Collection, & DNA Extraction



Black Lace Angelfish

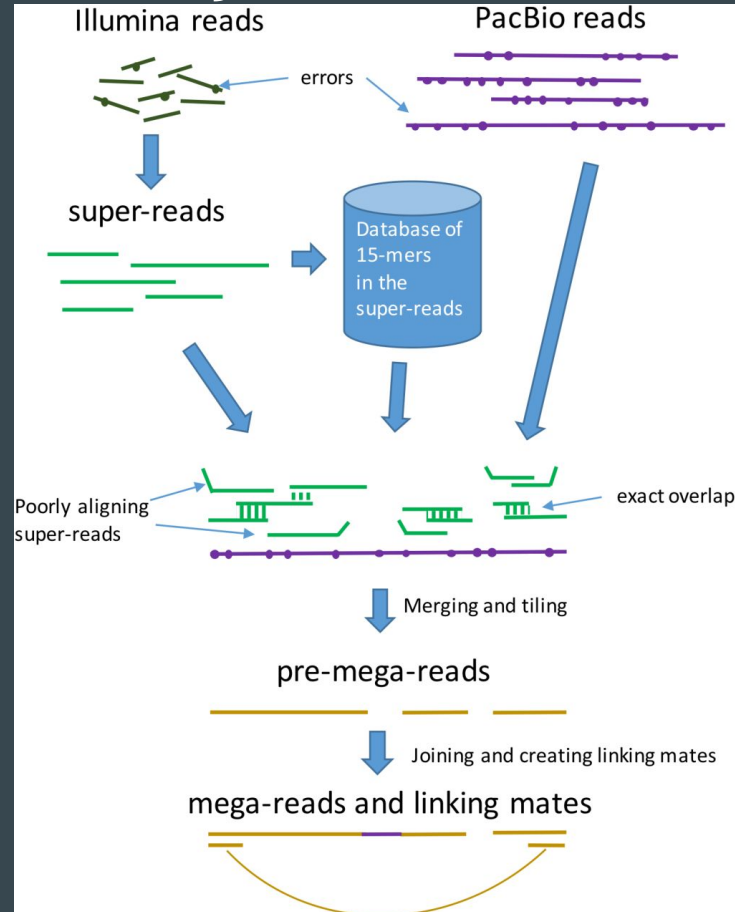


Pterophyllum scalare

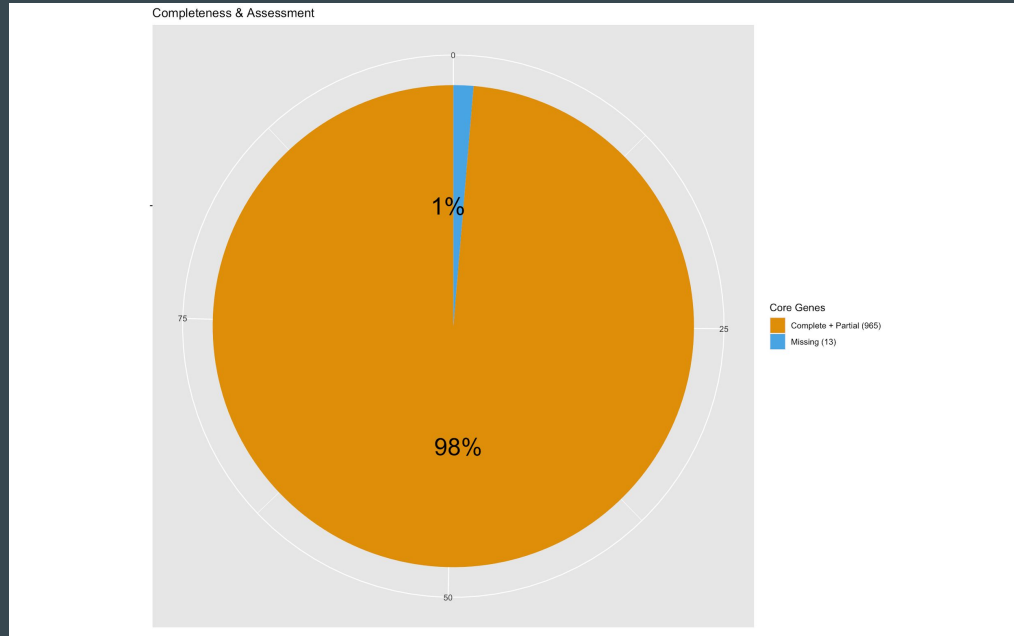
## SEQUENCING RESULTS:

Sequencing Type	Library	Coverage	Sample
Illumina	800 bp	57x	3ME2
PacBio	long reads	70x	3ME1

# Hybrid *de novo* Assembly



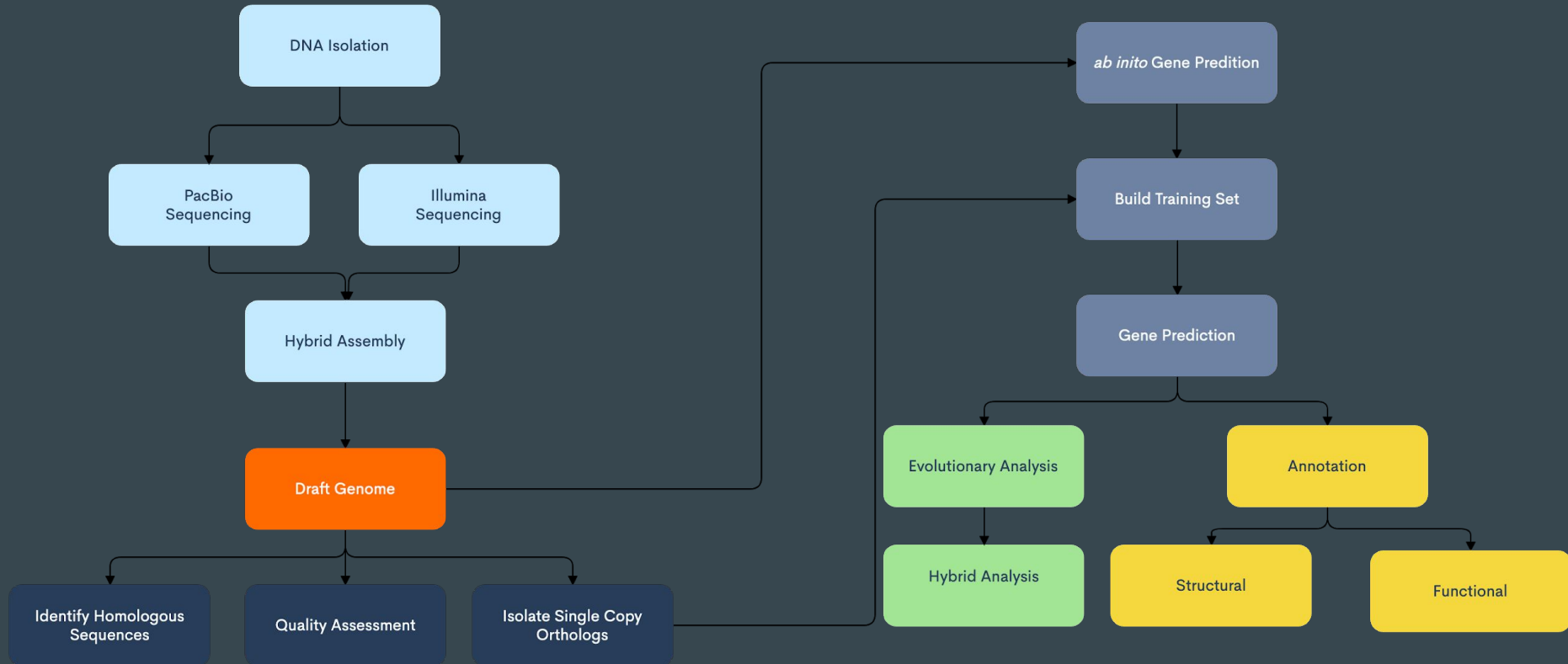
# Quality Assessment of the Draft Genome: BUSCO



## COMPLETENESS ASSESSMENT RESULTS:

Total number of core genes queried	978
Number of core genes detected	
Complete	959 (98.06%)
Complete + Partial	965 (98.67%)
Number of missing core genes	13 (1.33%)
Average number of orthologs per core genes	1.05
% of detected core genes that have more than 1 ortholog	4.59
Scores in BUSCO format	C:98.1%[S:93.6%,D:4.5%],F:0.6%,M:1.3%

# Strategy for Evolutionary Analysis



# Next Steps

- Complete evolutionary analysis
- Deposit annotated reference sequence
- Re-sequence wild type individuals collected from geographically diverse locations
- Compare genomes to identify which wild types have contributed to the angelfish origin



# Acknowledgments

Funding from CNLS

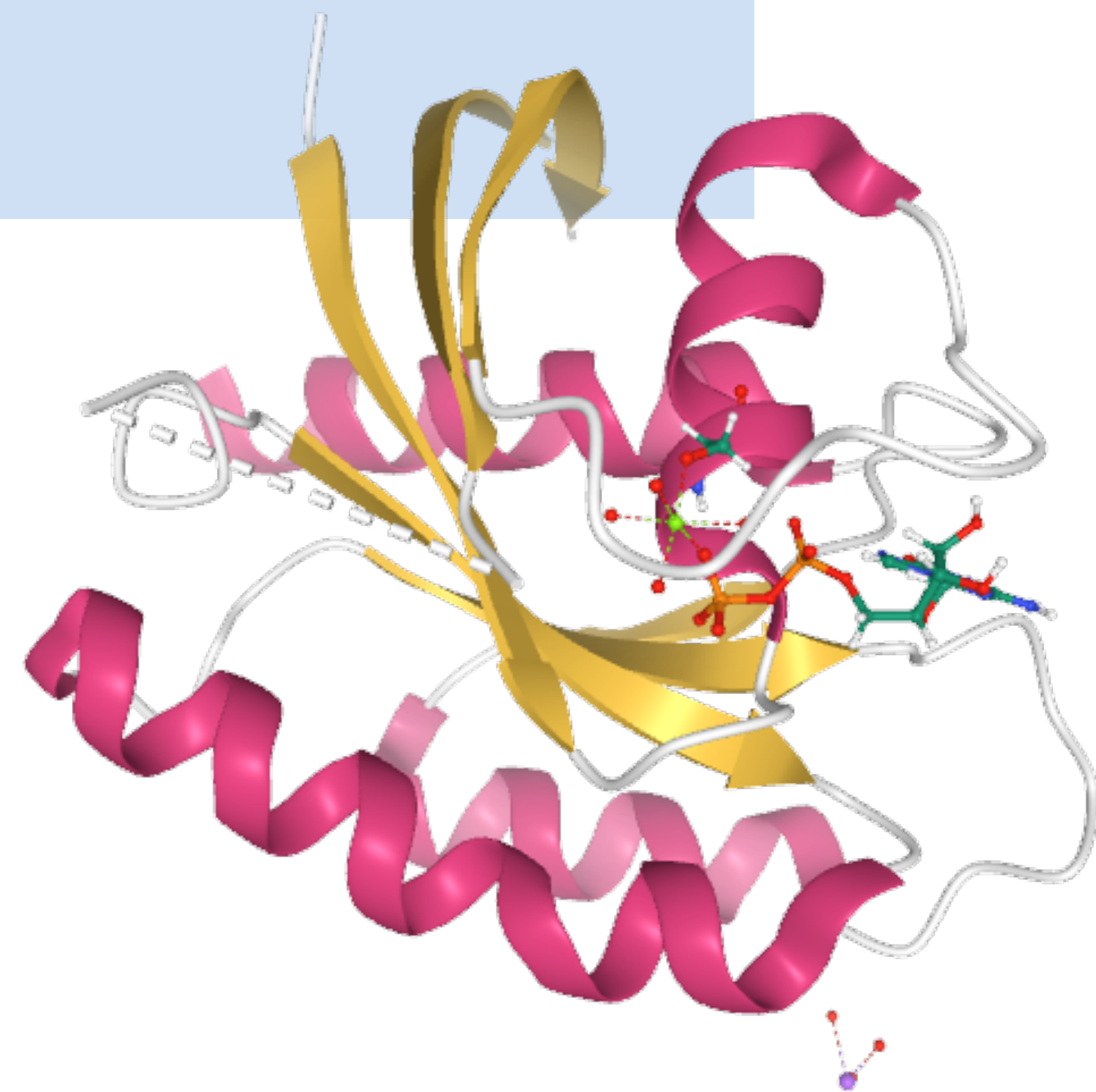
Mentoring:

William S. Hlavacek, T-6

# Computational Search for Cancer Treatments

**Los Alamos Theoretical Division T-6**

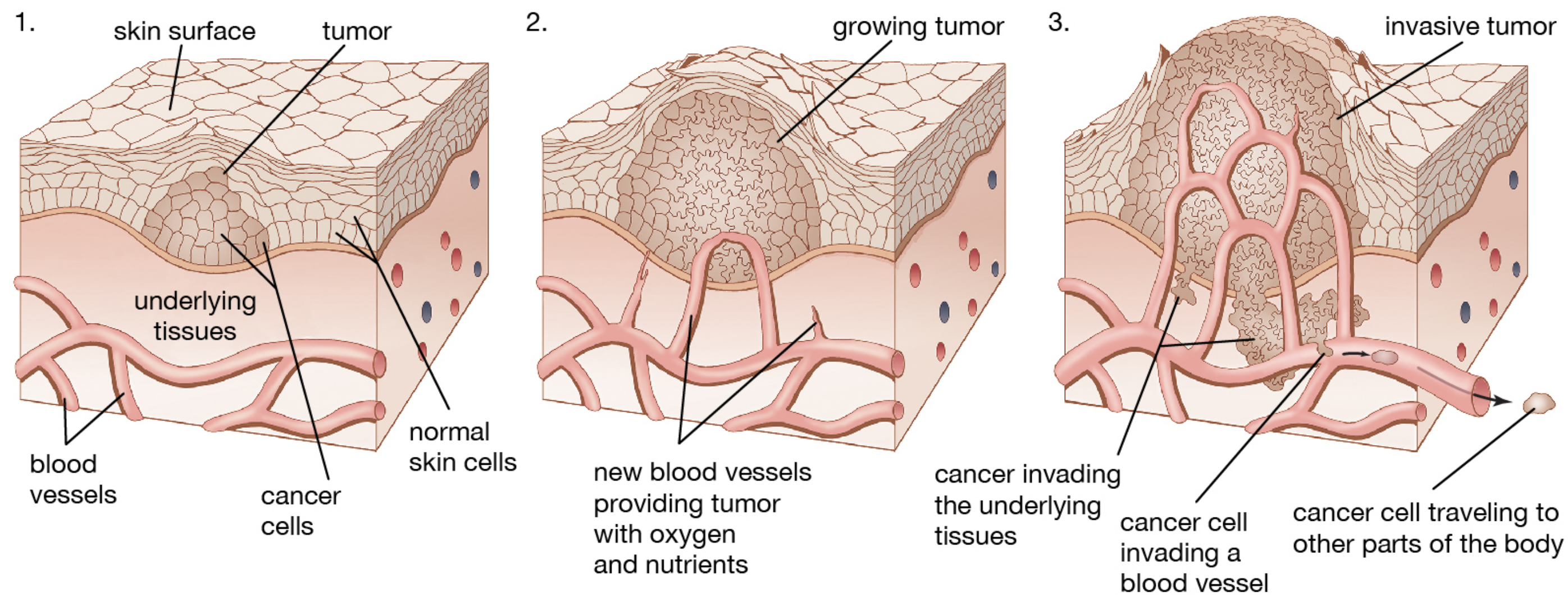
**Leah Darwin**  
**Mentor: Chris Neale**





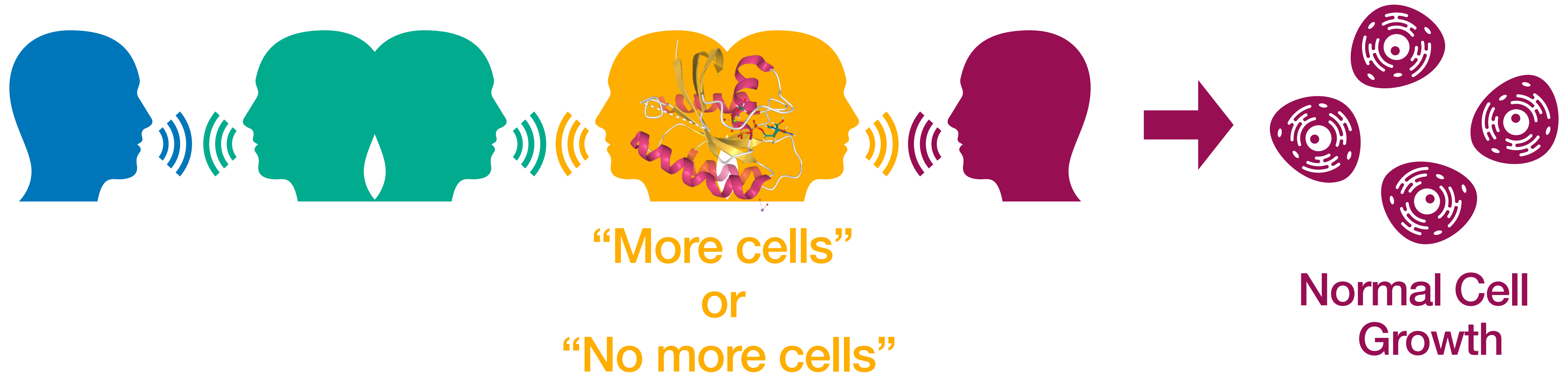
## What is cancer?

- Cancer is a leading cause of death world wide
- Uncontrollable cell growth leads to the formation of **tumors**



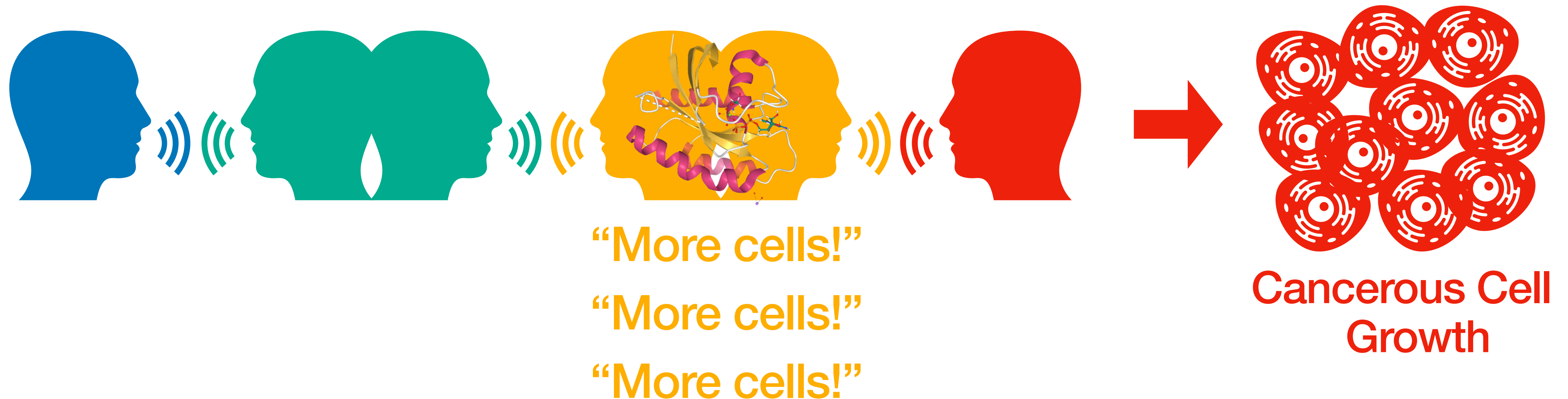
## Ras Mutations

- The Ras protein is involved in cell growth
- Ras works in a protein signaling pathway



## Ras Mutations

- Mutations to Ras can eliminate the on-off function
- Ras becomes stuck in the **on state**

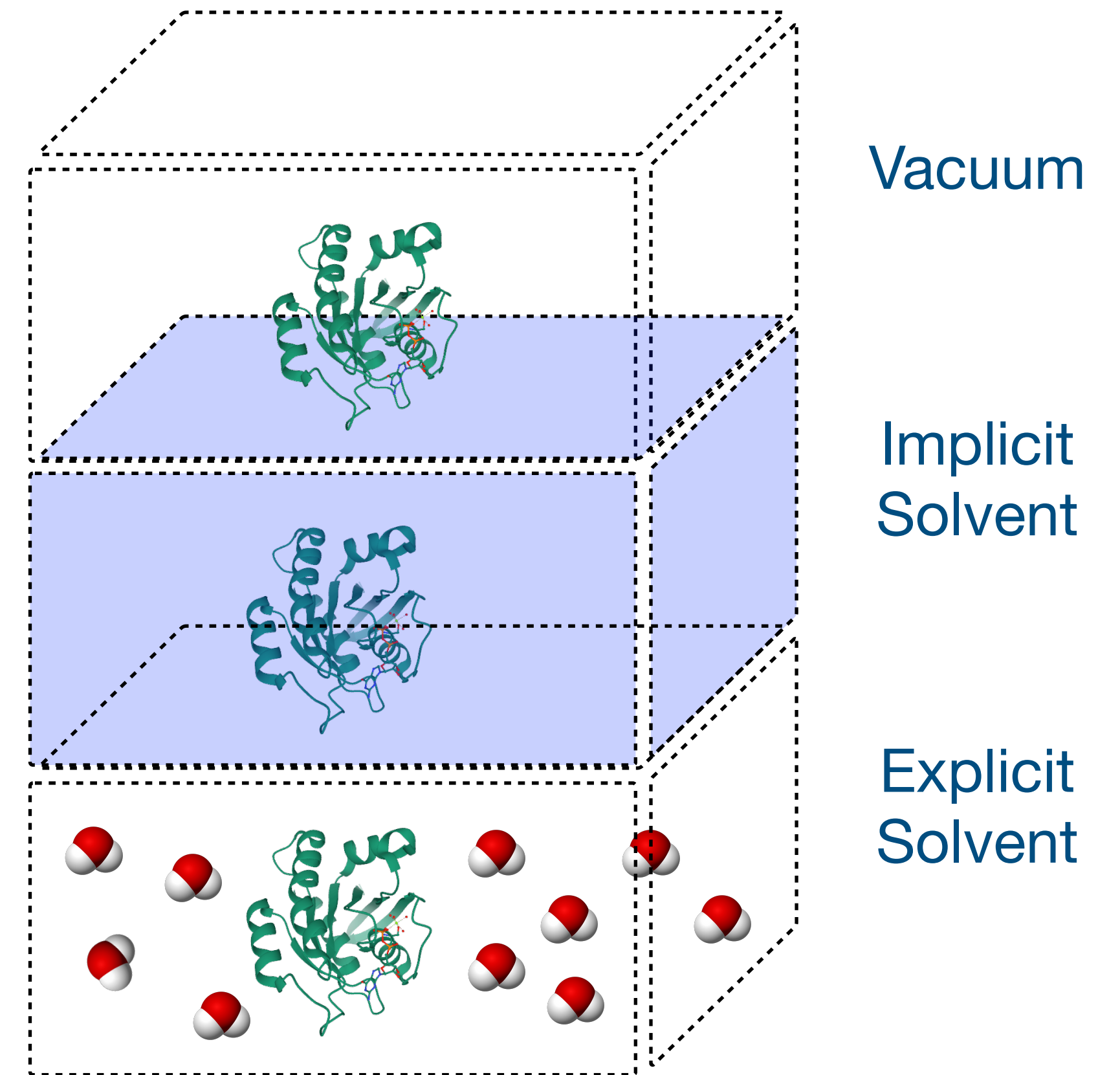


**Goal: Find drug candidates that will restore the on-off mechanism of Ras**



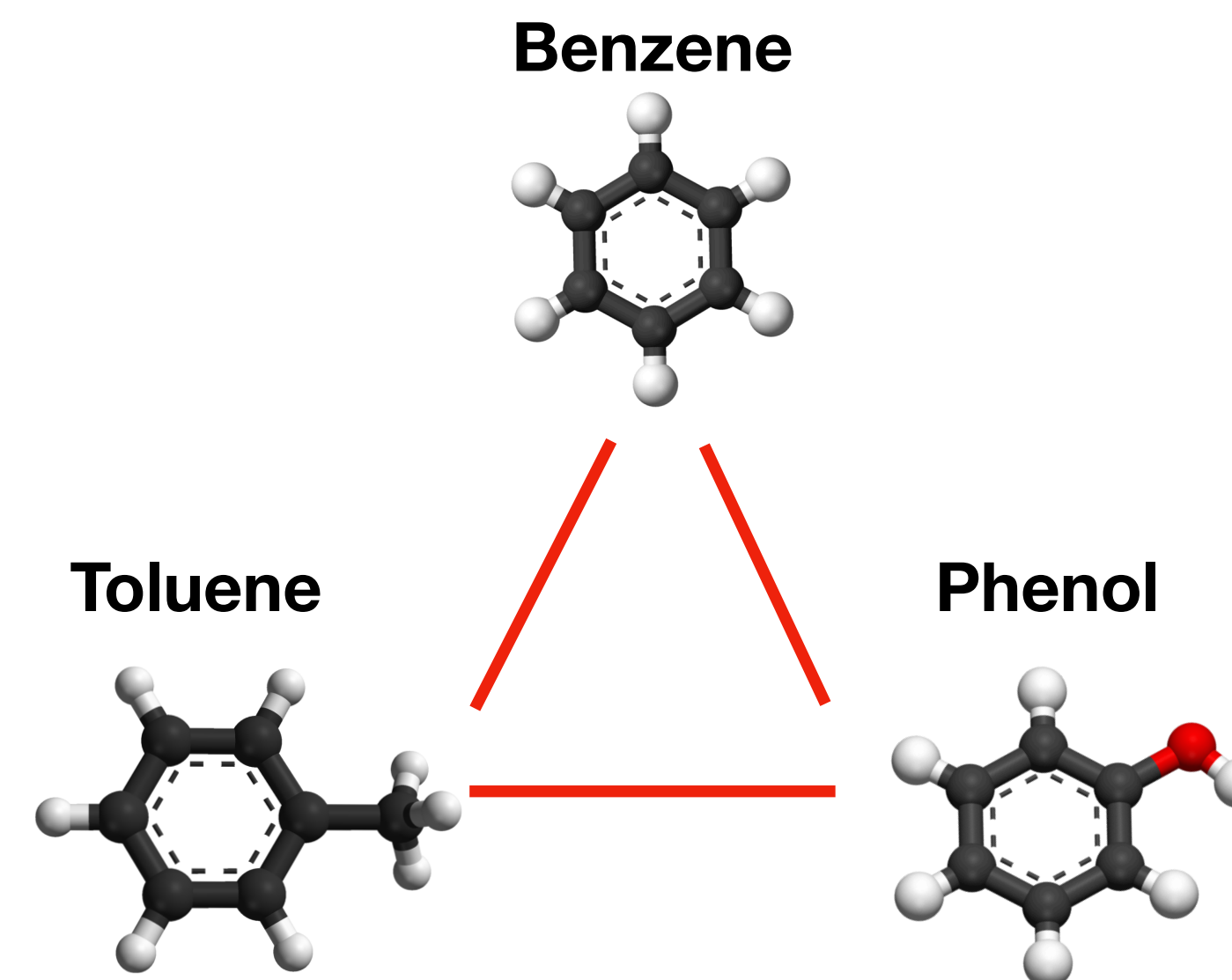
## Molecular Dynamics (MD)

- Emulate the conditions of the human body through computer simulation
- The protein is able to interact with a chemical for a finite period of time
- Newton's equations of motion determine trajectories

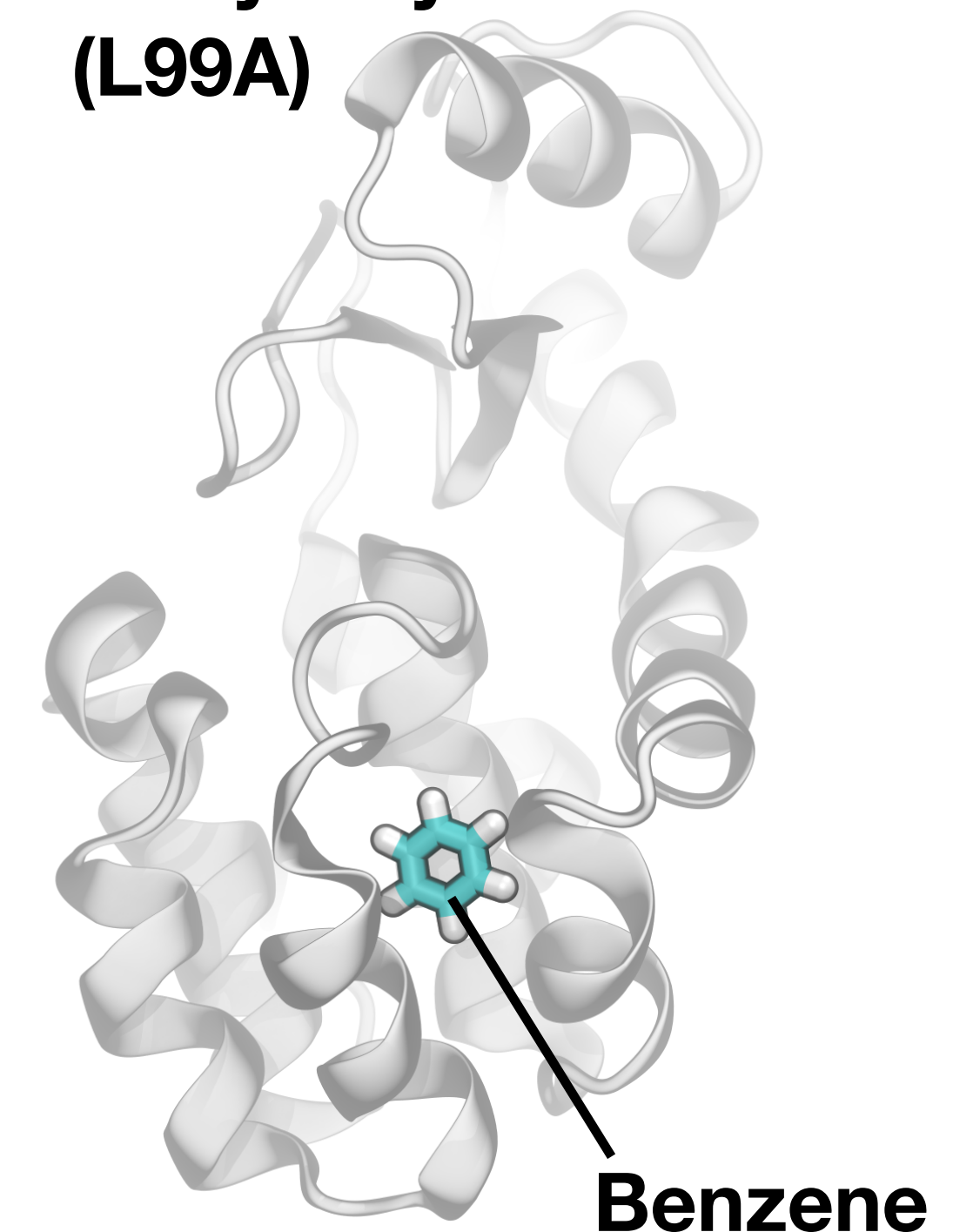


## MD-Monte Carlo and T4 Lysozyme Model System

- **T4 Lysozyme** is used as a model system
- Three chemicals were used as they have well studied binding properties for T4 Lysozyme
- Monte Carlo moves allow chemicals to change into close structural relatives



**T4 Lysozyme  
(L99A)**



## Alignment Algorithm Version 1



- First version of algorithm has issues with alignment
- Wrong configuration
- Placement of hydroxyl group is not random

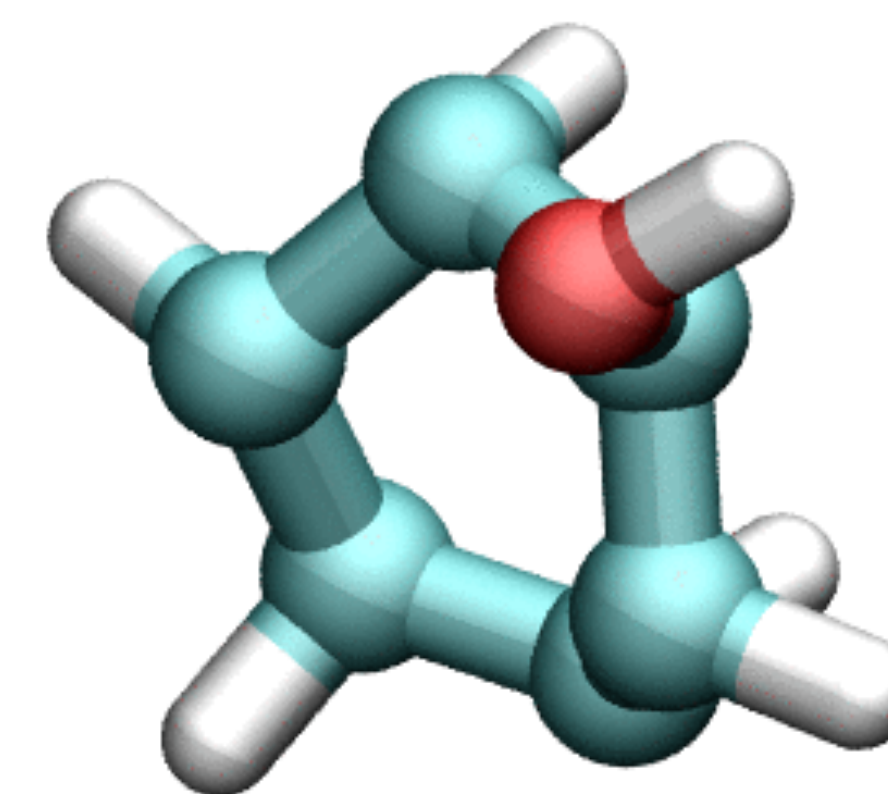
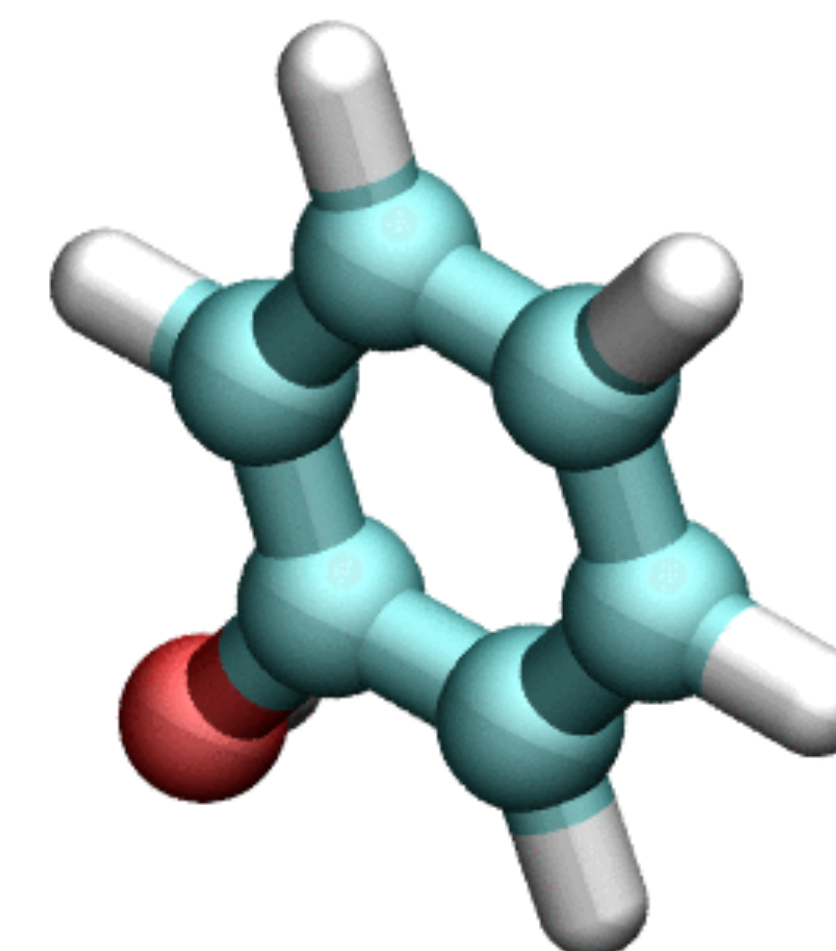


Image courtesy of: Lara A. Patel

## Alignment Algorithm Version 2

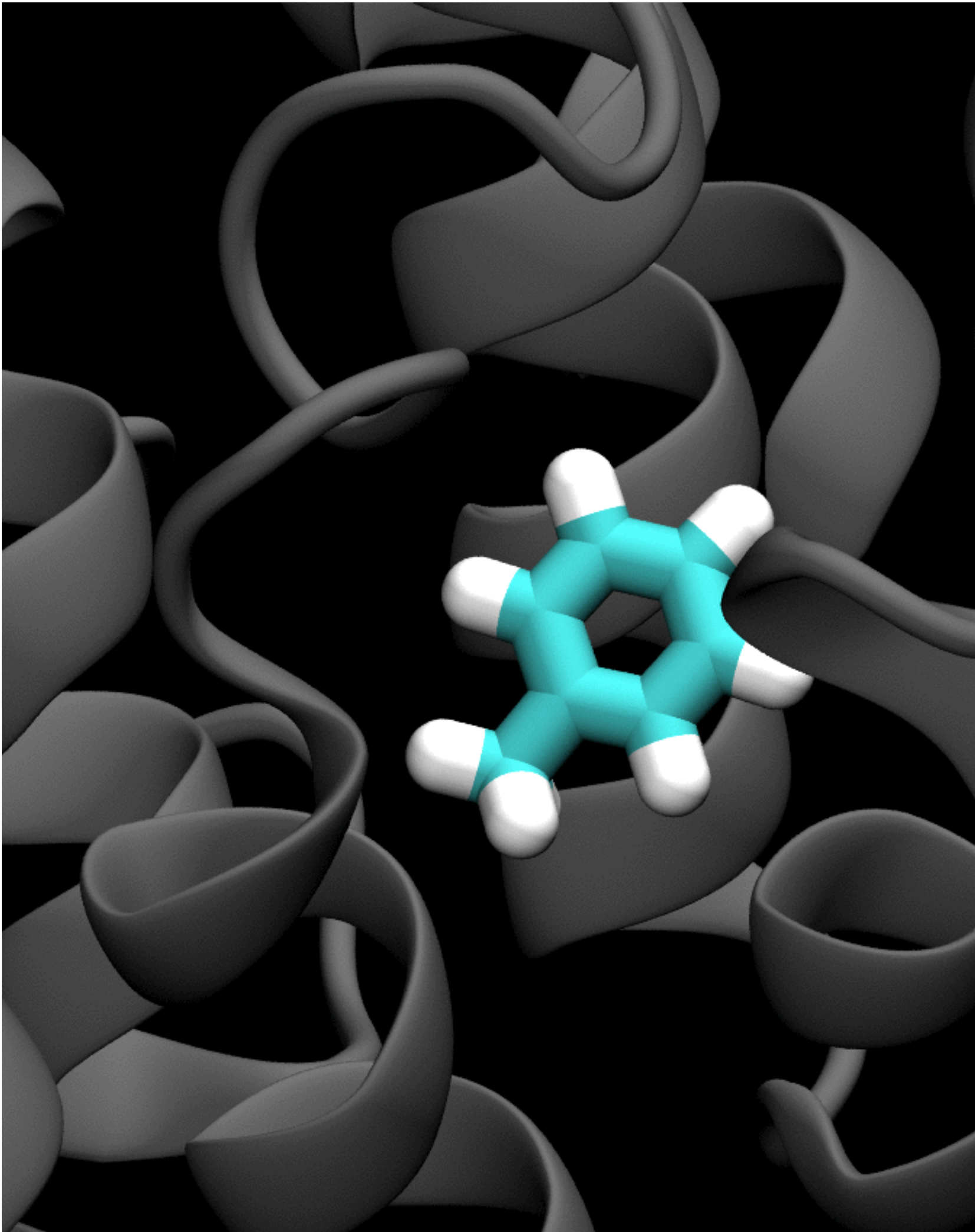


- Updated version of alignment algorithm fixes major issues with the first version
  - Now in correct configuration
  - Random placement of hydroxyl group





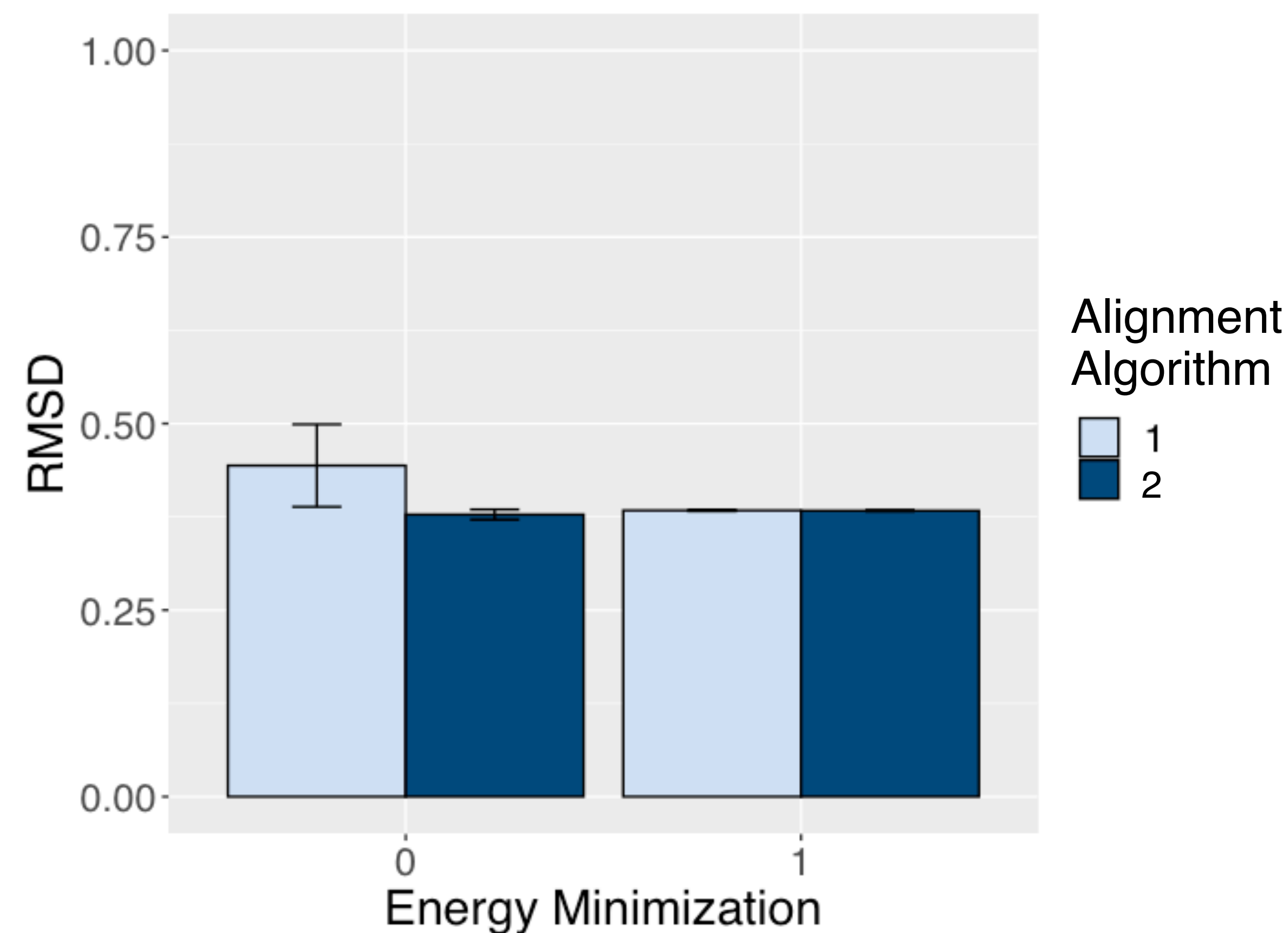
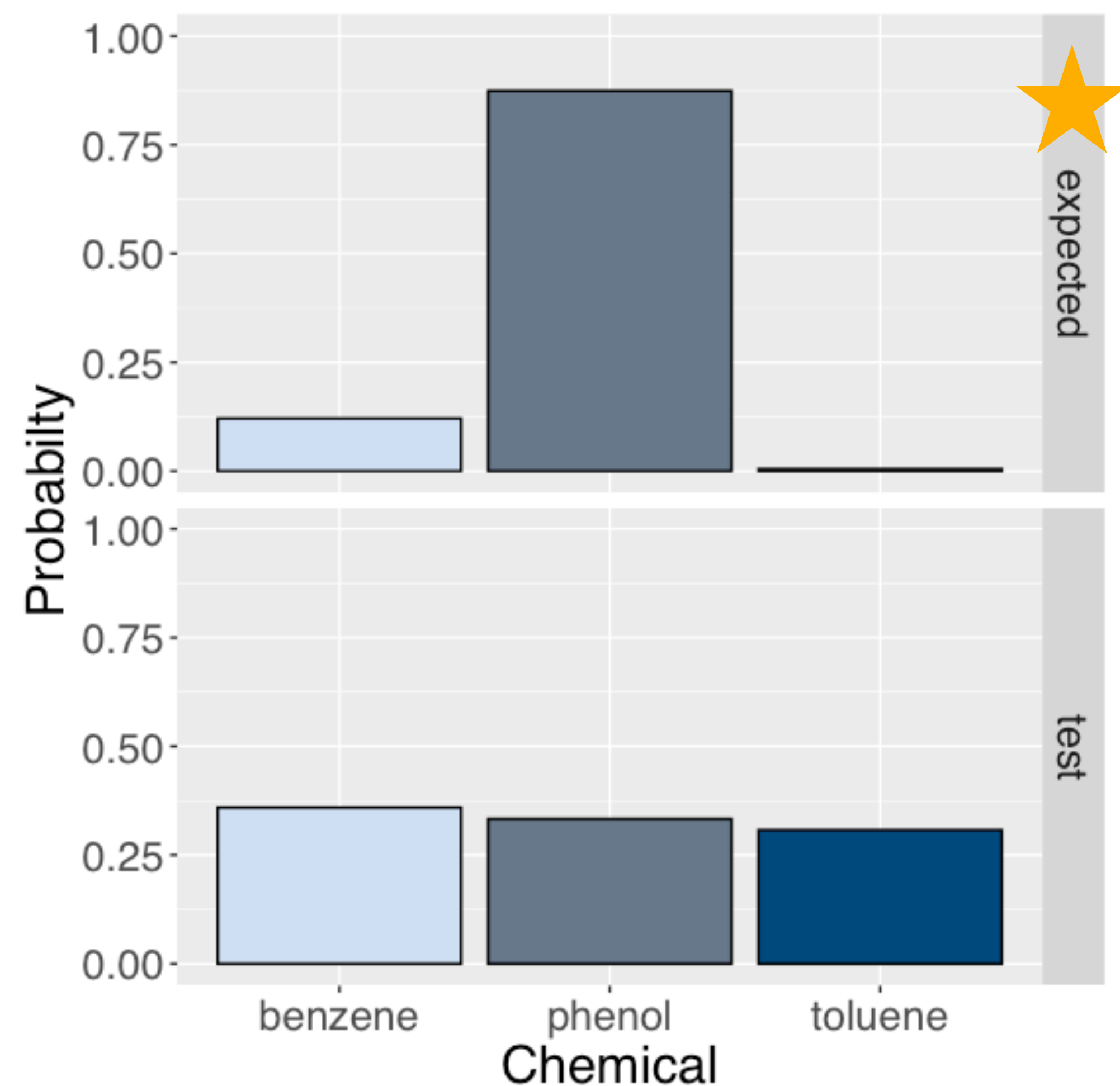
Fixed Protein in Vacuum



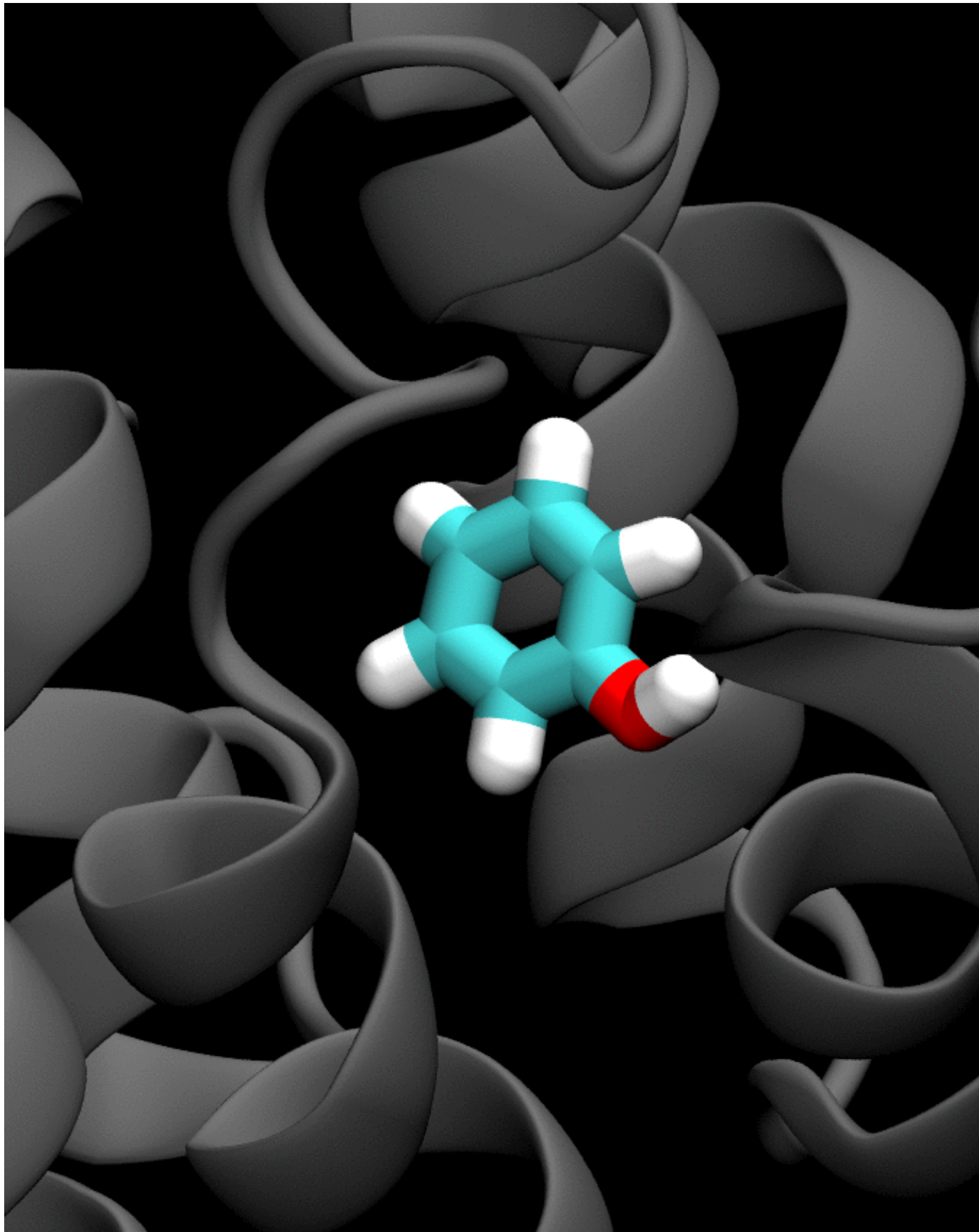
Parameter	Values
Environment	<b>Vacuum</b> , Implicit Solvent
Energy Minimization	ON, OFF
Alignment Algorithm	Version 1, Version 2



## Fixed Protein in Vacuum

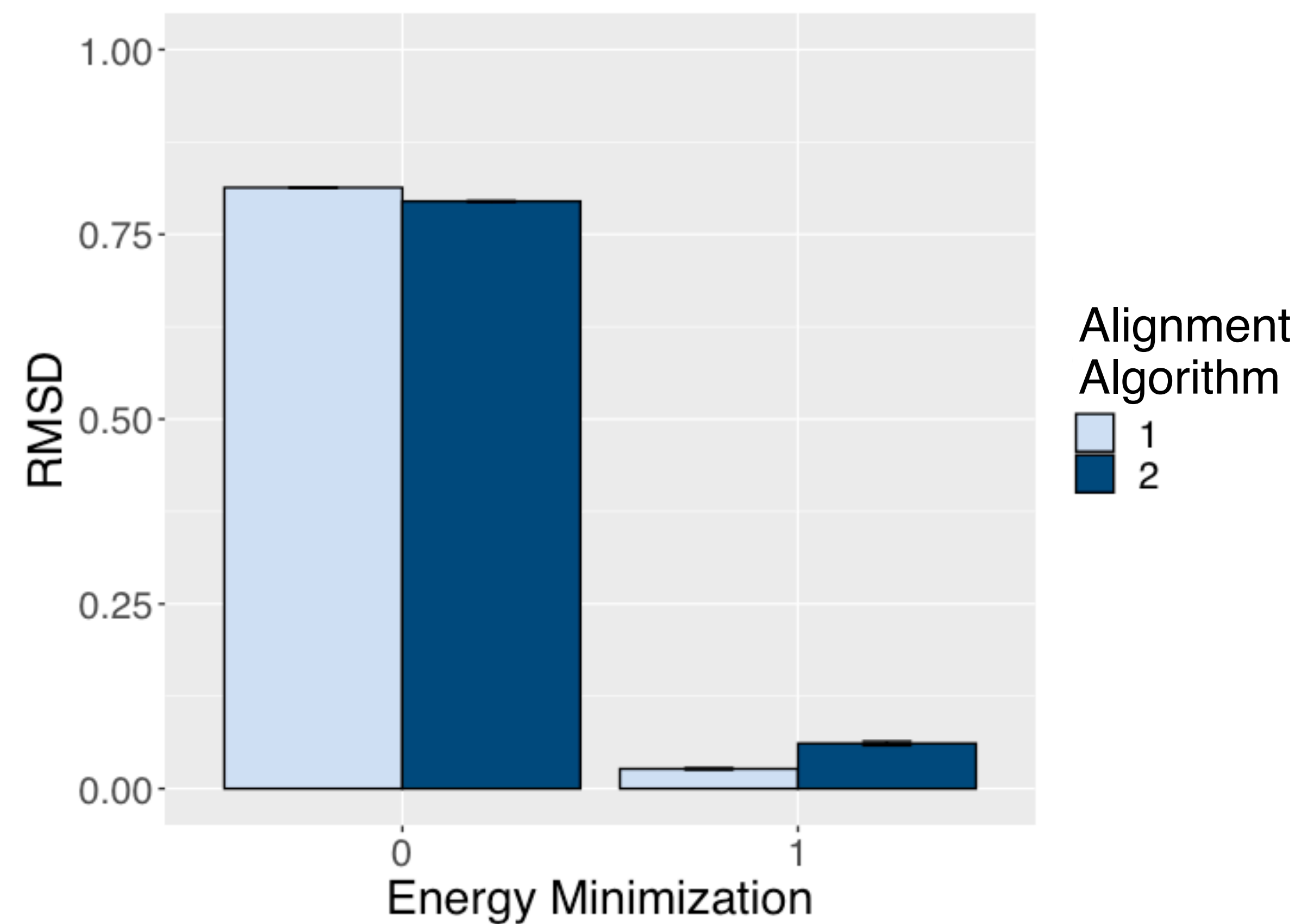
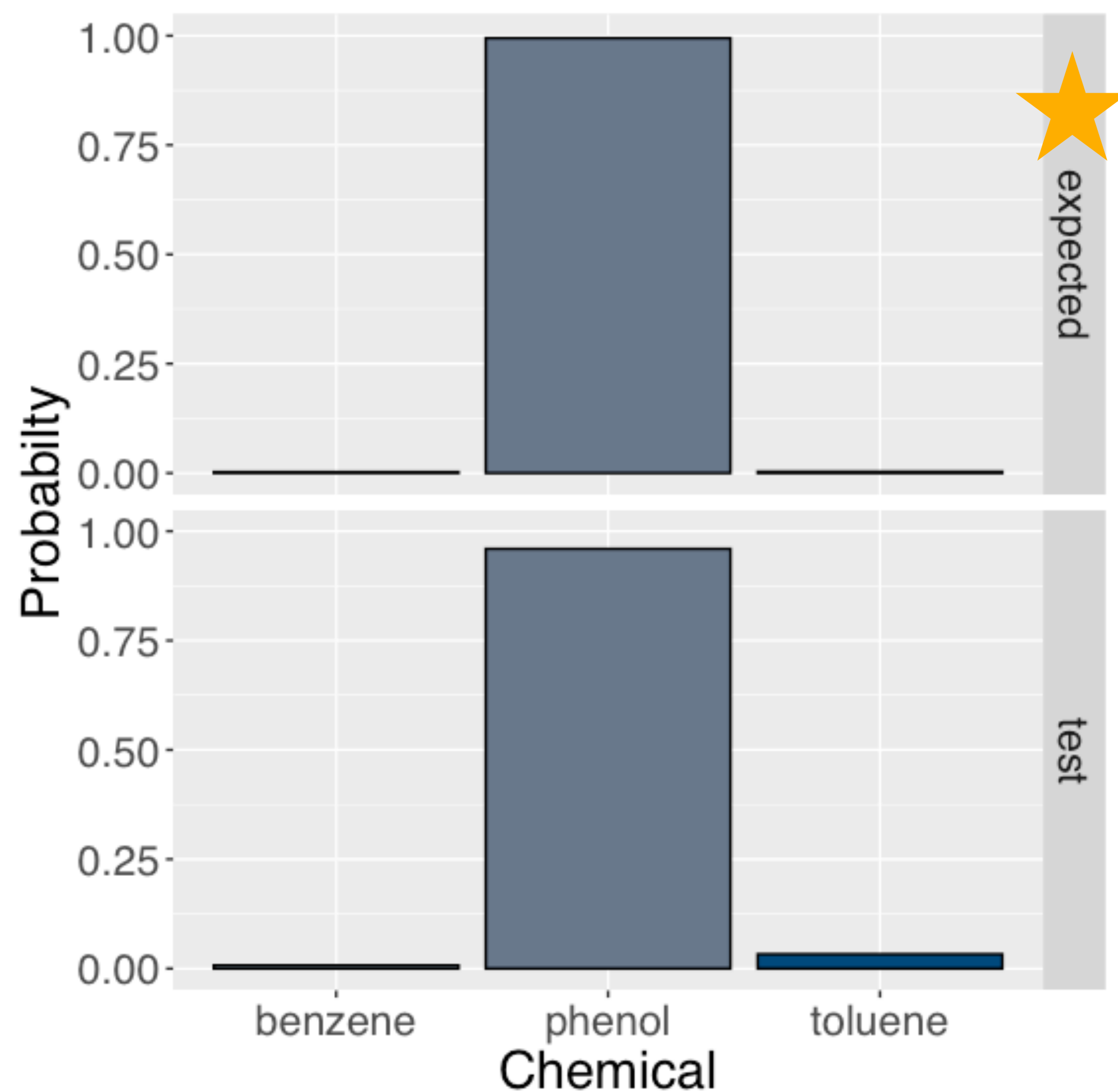


Dynamic Protein in Vacuum



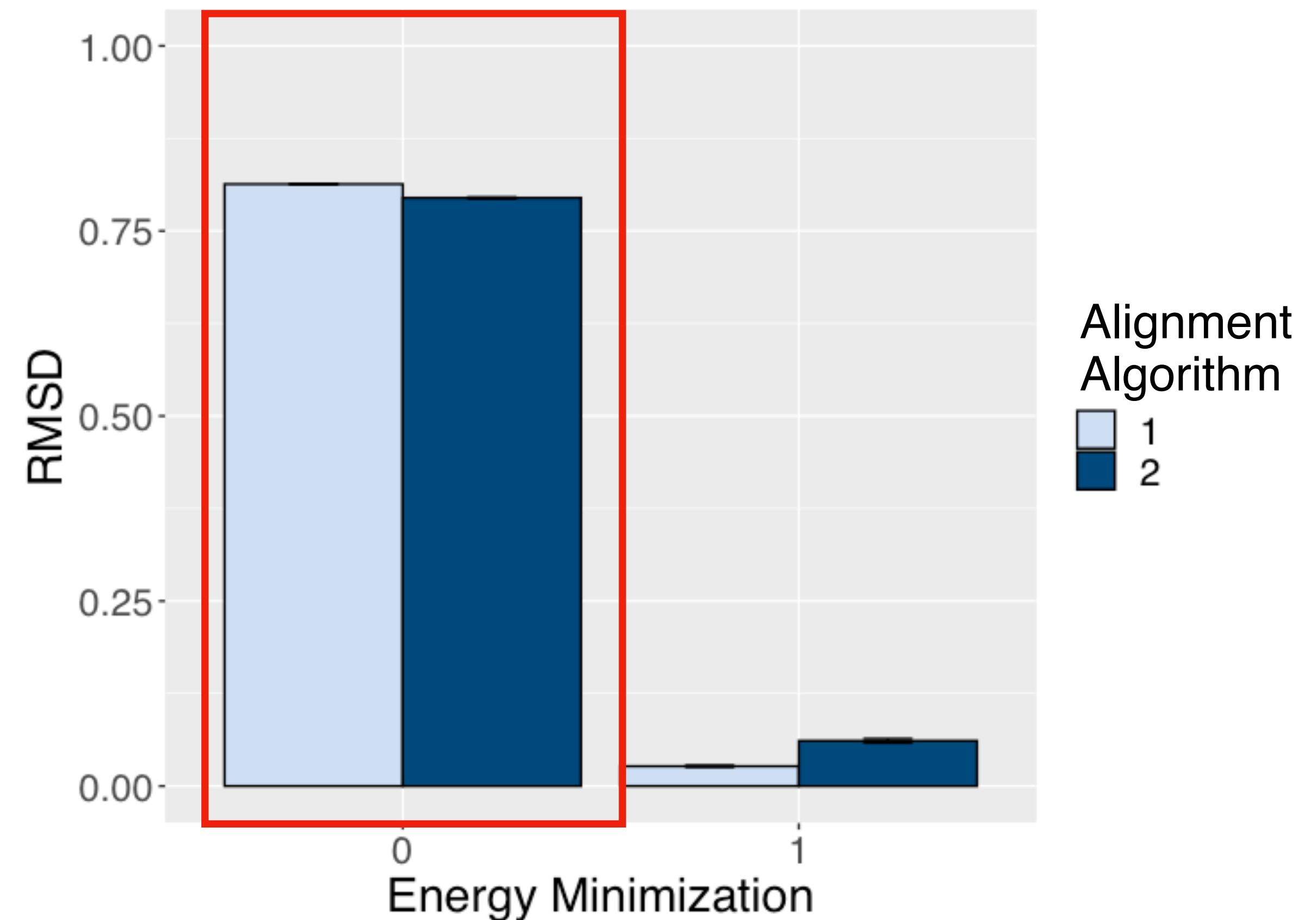
Parameter	Values
Environment	<b>Vacuum</b> , Implicit Solvent
Energy Minimization	ON, OFF
Alignment Algorithm	Version 1, Version 2

## Dynamic Protein in Vacuum



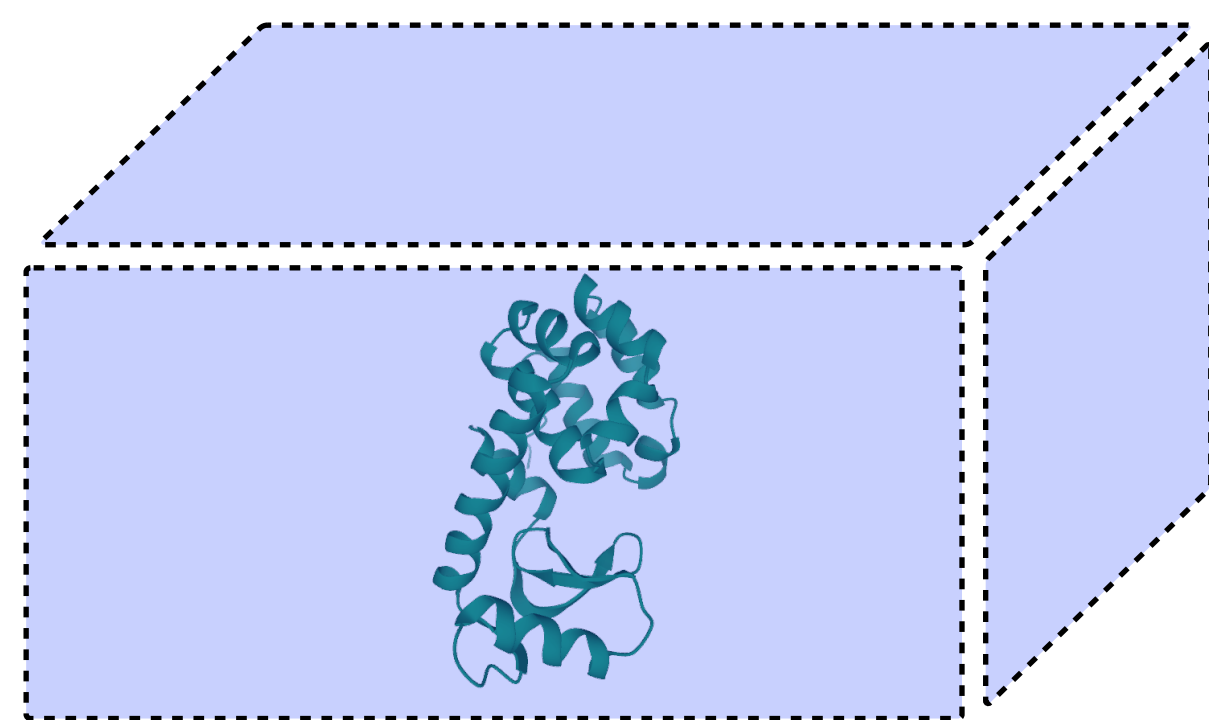
## Dynamic Protein in Vacuum

- Alignment algorithm 1 selected benzene 100% of the time
  - Didn't explore any other chemical space
- Alignment algorithm 2 did slightly better because it was able to move away from benzene

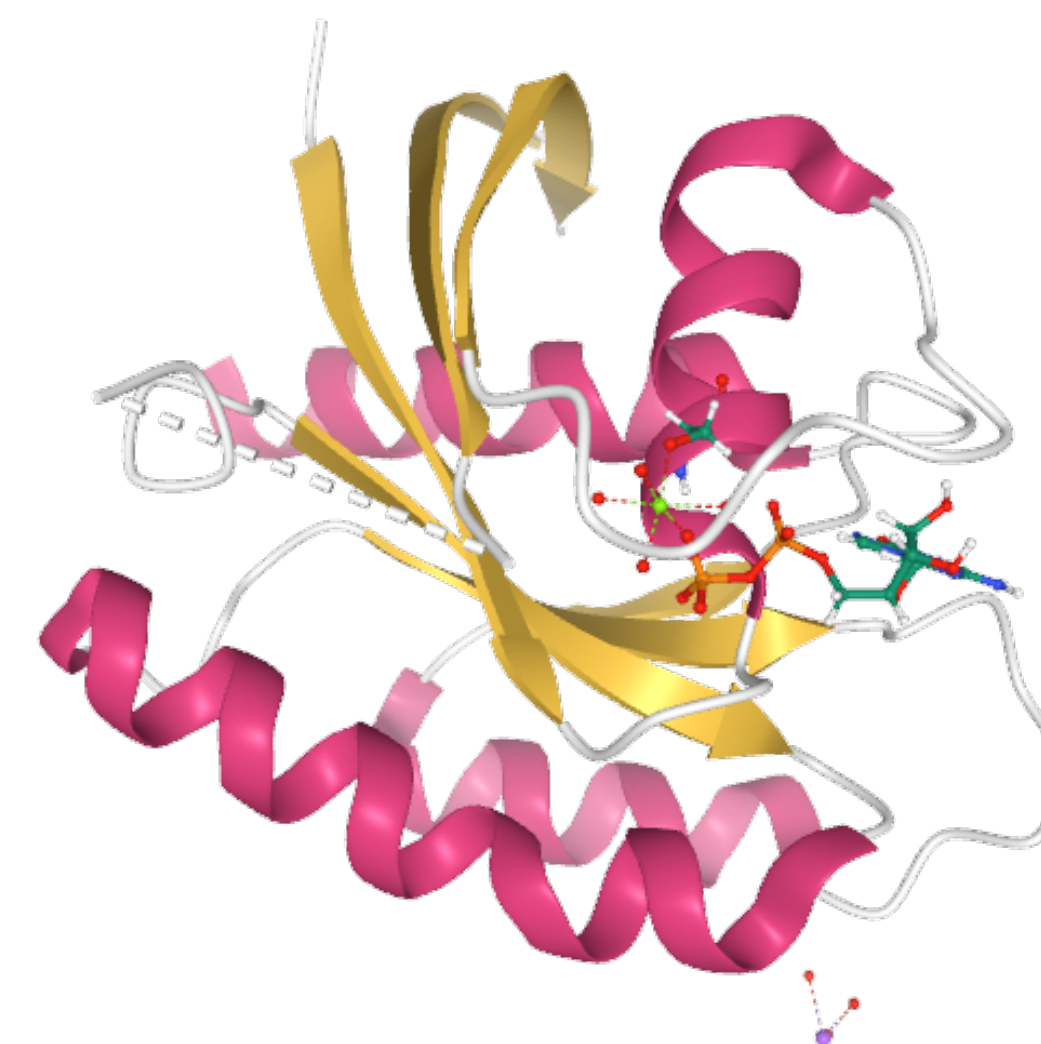
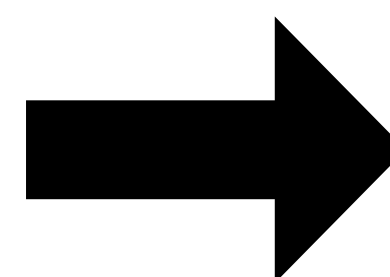


- Algorithm 1 and 2 showed very similar results under the same conditions
- However, algorithm 2 allows for more chemical space to be explored

## Next Steps:



T4 Lysozyme  
Implicit Solvent  
Model

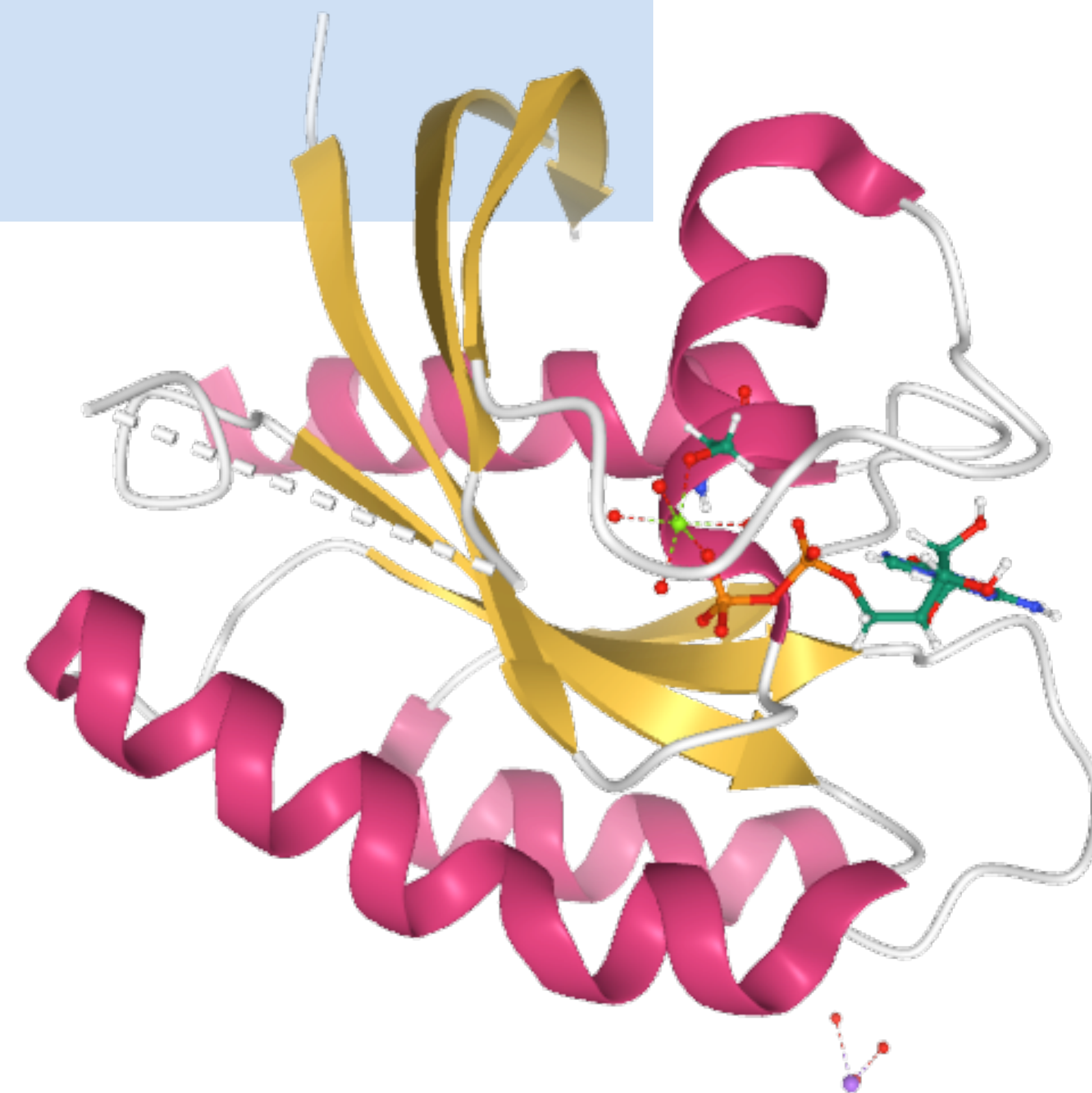


Apply software to  
find compatible drug  
candidates for RAS



# Thank You!

**Leah Darwin**  
**Mentor: Chris Neale**



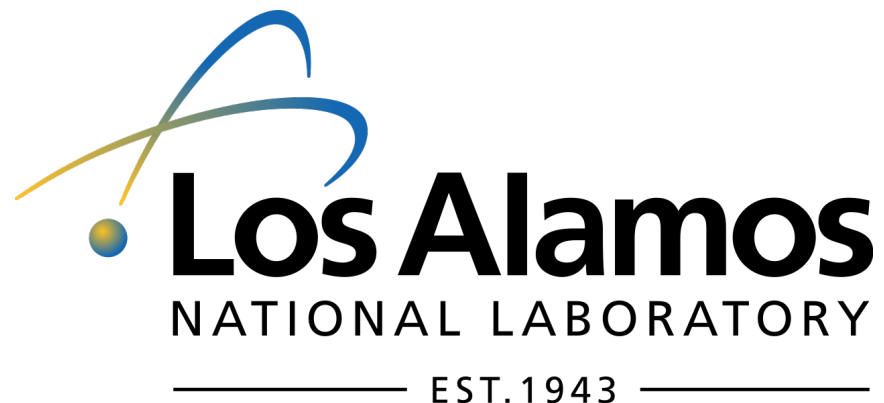
# Machine learning parametrization of empirical and semiempirical methods for improving transition metal chemistry

**Nikita Fedik**

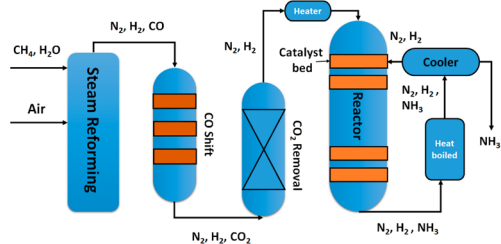
3<sup>rd</sup> year PhD student at Utah State University

GRA at T-1 division, LANL

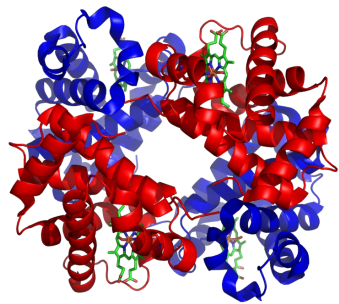
Mentors: Benjamin Nebgen, Sergei Tretiak



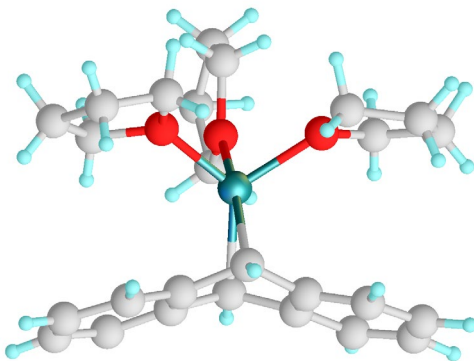
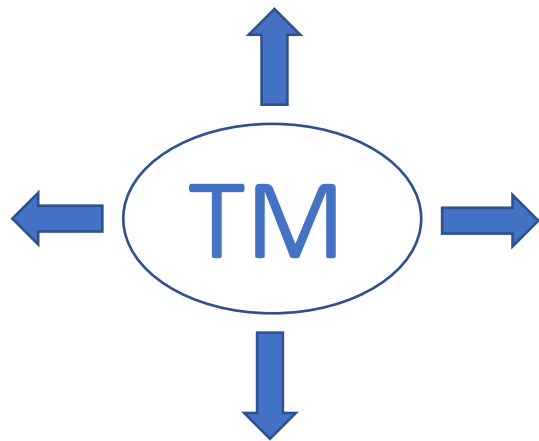
# Transition metal chemistry



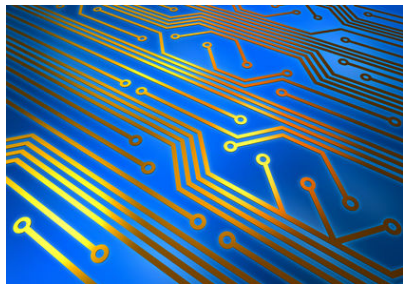
catalysis and industry



biochemistry



organometallics



semiconductors

## Simulation difficulties of TMs:

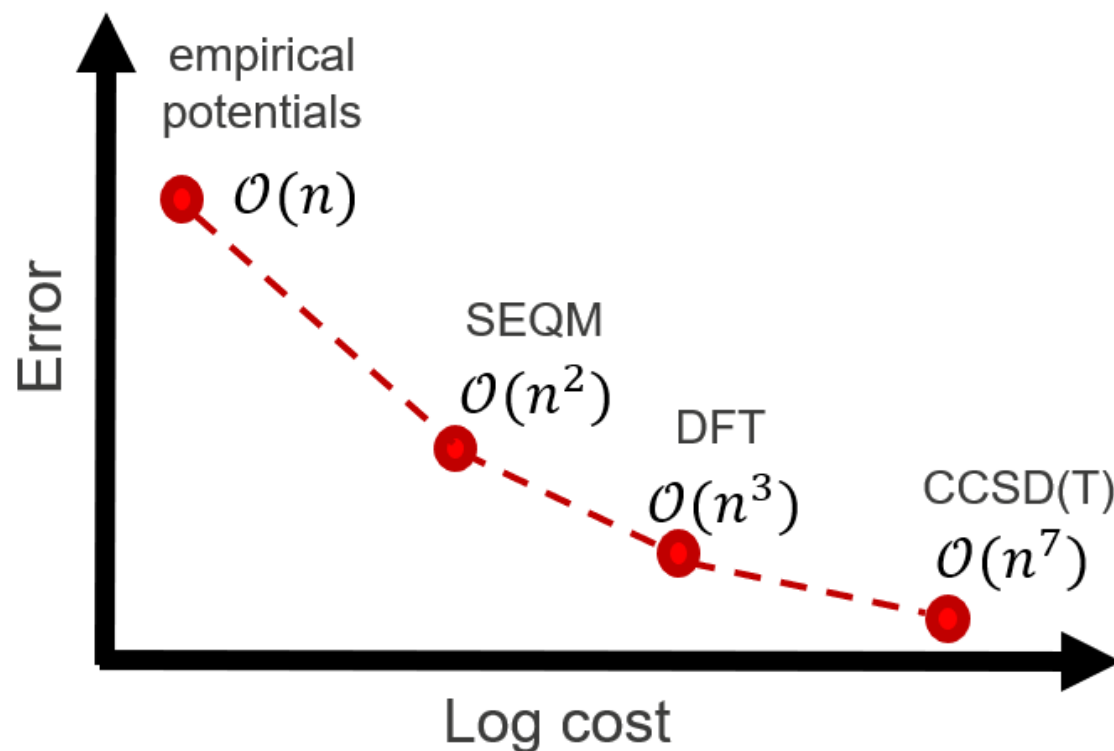
- large number of valence  $\bar{e}$
- degeneracy of states
- complex wave functions (multireference)
- possible relativistic effects



# Available methods

Empirical potentials usually do not have parameters for TMs at all.

Semiempirical QM are very inaccurate.



	MAD	RMSD
PM6	21.6	28.9
PM7	106.4	146.1
PM7-TS	68.2	141.9

THE JOURNAL OF  
PHYSICAL CHEMISTRY A

Cite This: J. Phys. Chem. A 2019, 123, 3761–3781

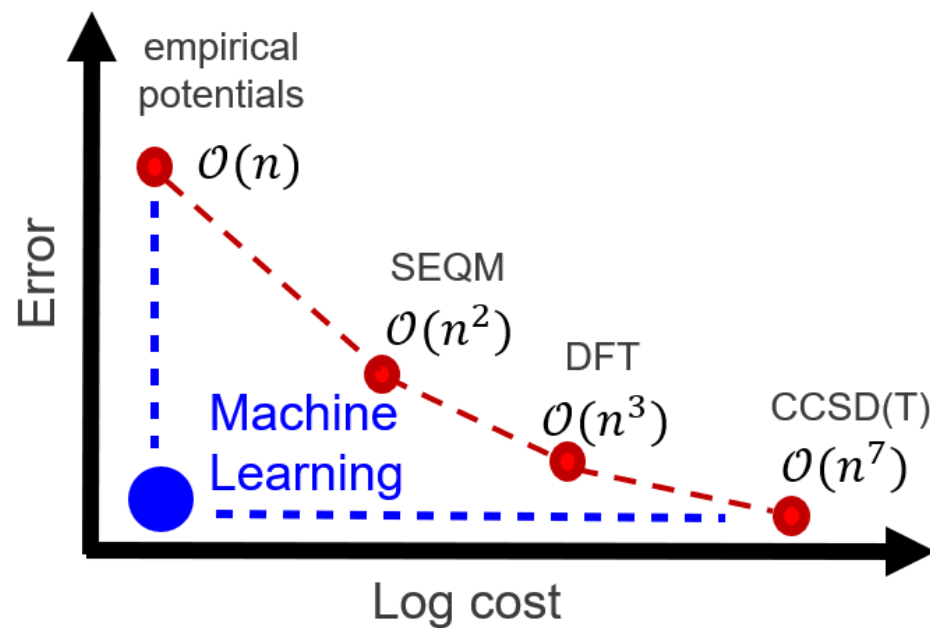
Article  
pubs.acs.org/JPCA

**Evaluating Transition Metal Barrier Heights with the Latest Density Functional Theory Exchange–Correlation Functionals: The MOBH35 Benchmark Database**

Published as part of The Journal of Physical Chemistry virtual special issue “Leo Radom Festschrift”.

Mark A. Iron<sup>\*,†</sup> and Trevor Janes<sup>‡</sup>

# Machine learning for dynamic parametrization



Neural Network potentials

- HIPNN architecture

THE JOURNAL OF CHEMICAL PHYSICS 148, 241715 (2018)



## Hierarchical modeling of molecular energies using a deep neural network

Nicholas Lubbers,<sup>1,a)</sup> Justin S. Smith,<sup>1,2</sup> and Kipton Barros<sup>1</sup>

<sup>1</sup>Theoretical Division and CNLS, Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA

<sup>2</sup>Department of Chemistry, University of Florida, Gainesville, Florida 32611, USA

ML parametrized SEQM

- PYSEQM

JCTC  
Journal of Chemical Theory and Computation

pubs.acs.org/JCTC

Article

## Graphics Processing Unit-Accelerated Semiempirical Born Oppenheimer Molecular Dynamics Using PyTorch

Guoqing Zhou, Ben Nebgen,<sup>\*</sup> Nicholas Lubbers, Walter Malone, Anders M. N. Niklasson, and Sergei Tretiak

# HIP-NN: Hierarchically Interacting Particle Neural Network

## 1. Local atomic environment

$$E \approx \hat{E} = \sum_{i=1}^{N_{\text{atom}}} \hat{E}_i.$$

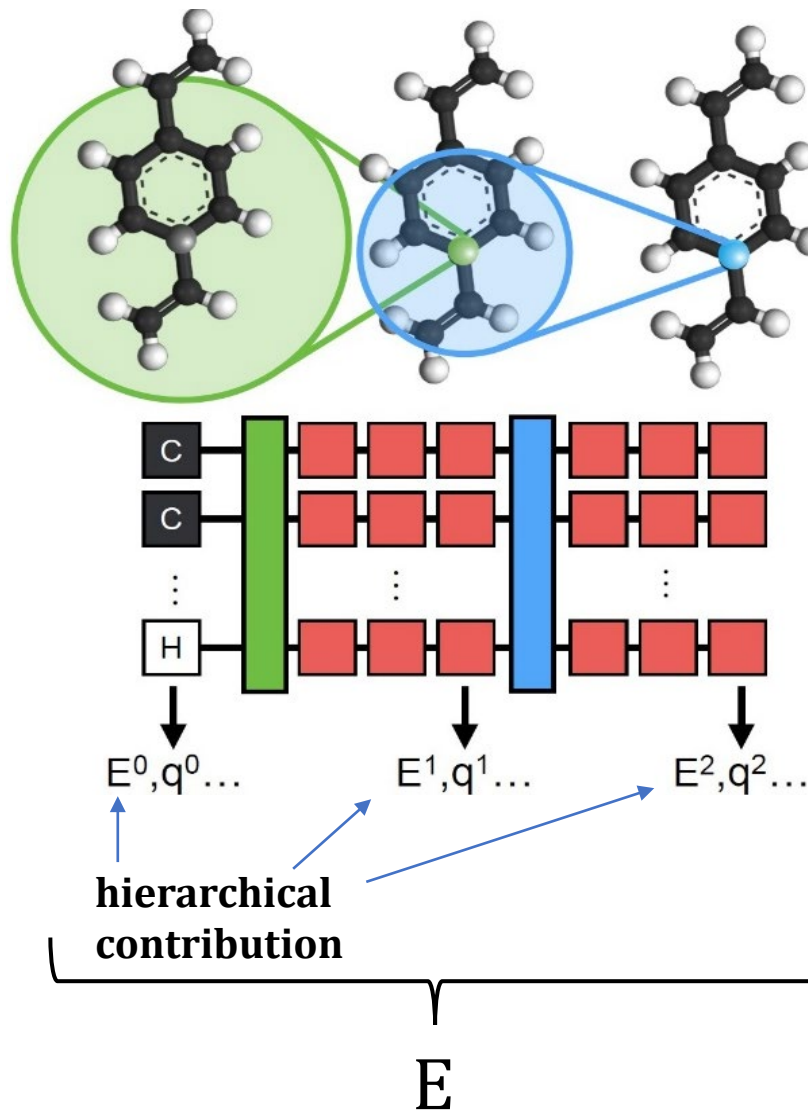
## 2. Hierarchical decomposition ansatz

$$\hat{E}_i = \sum_{n=0}^{N_{\text{interaction}}} \hat{E}_i^n.$$

## 3. Molecular representation

$$\{R\} \longrightarrow \{|r_i - r_j|\} \text{ if } r_{ij} < R_{\text{cutoff}}$$

*invariant to rotations and translations*



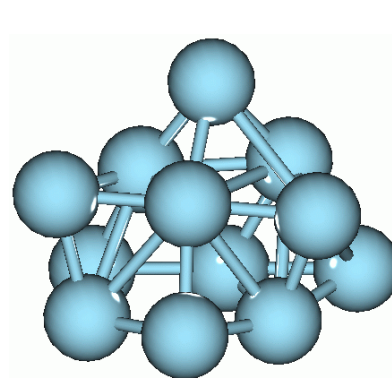
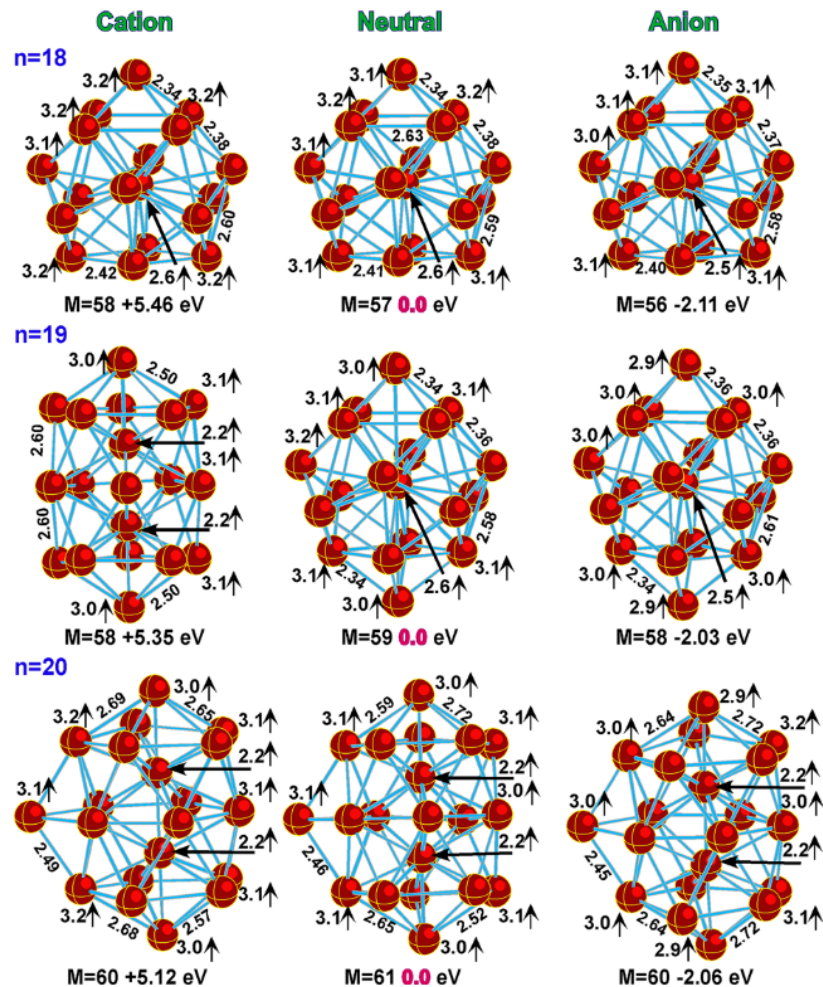
$$\hat{E}_i^n = \sum_{a=1}^{N_{\text{feature}}} w_a^n z_{i,a}^{\ell_n} + b^n,$$

**learned parameters**

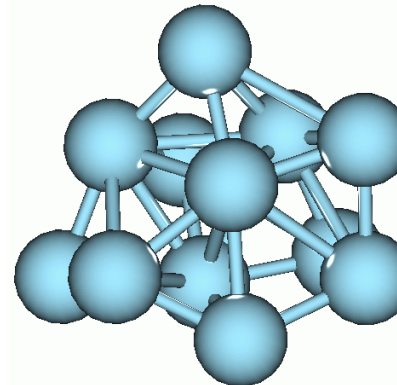
GPU training > 100 times faster than CPU

We need a big and consistent dataset ( $10^3 - 10^6$  entries)

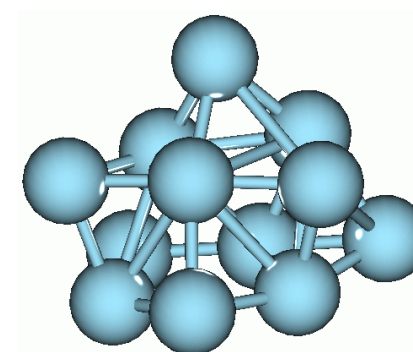
# Why pure TM clusters is a very a bad idea?



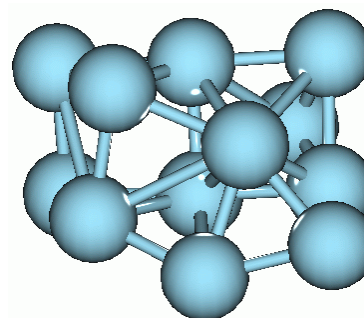
Pt<sub>12</sub>, <sup>7</sup>A “cube”  
-1430.881773 A.U.



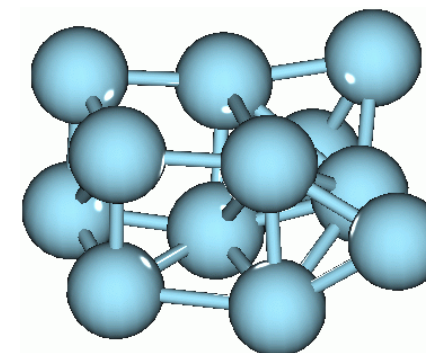
Pt<sub>12</sub>, <sup>5</sup>A “cube”  
-1430.875799 A.U.



Pt<sub>12</sub>, <sup>1</sup>A “cube” (from <sup>7</sup>A)  
-1430.871506 A.U.



Pt<sub>12</sub>, <sup>11</sup>A “cube”  
-1430.864530 A.U.



Pt<sub>12</sub>, <sup>9</sup>A “cube”  
-1430.866841 A.U.

$M = 2S+1$  are in range **22 – 61!**

G. L. Gutsev et al. J. Phys. Chem. A  
2012, 116, 10218–10228

# New organometallic database

## 1. Organometallic (HCNO) complexes of Ti

- closed shell singlets
- only 4 *ve* ( $3d^24s^2$ )

## 2. Reference method – what to choose?

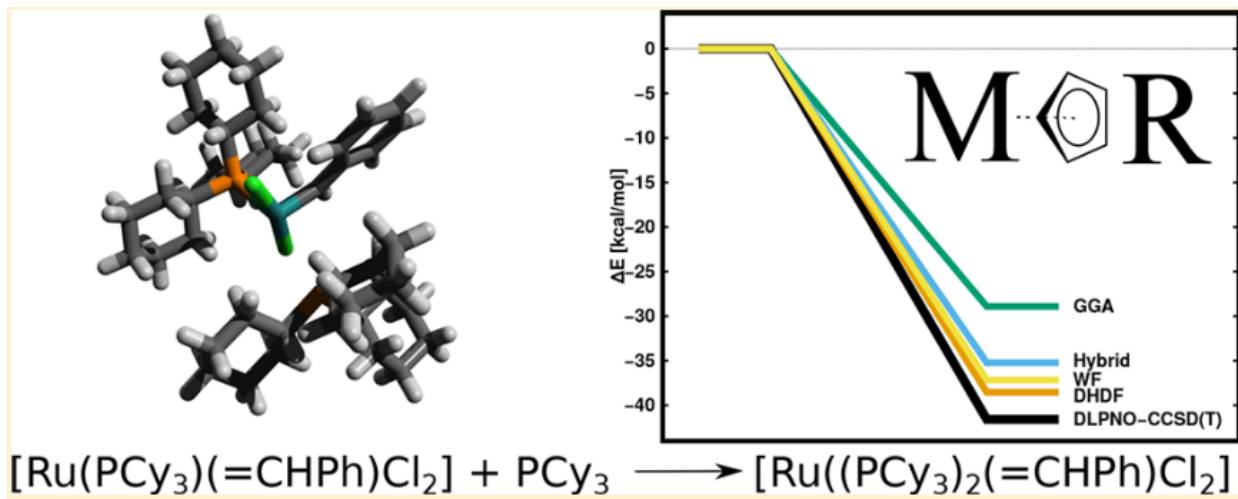
### Comprehensive Thermochemical Benchmark Set of Realistic Closed-Shell Metal Organic Reactions

Sebastian Dohm,<sup>†</sup> Andreas Hansen,<sup>\*,†,§</sup> Marc Steinmetz,<sup>†,§</sup> Stefan Grimme,<sup>†,§</sup> and Marek P. Checinski<sup>\*,‡</sup>

### Evaluating Transition Metal Barrier Heights with the Latest Density Functional Theory Exchange–Correlation Functionals: The MOBH35 Benchmark Database

Published as part of *The Journal of Physical Chemistry virtual special issue “Leo Radom Festschrift”*.

Mark A. Iron<sup>\*,†,§</sup> and Trevor Janes<sup>‡,§</sup>



1. Hybrids are probably the best choice:
  - $\omega$ B97-D3BJ (MAD = 1.9 kcal/mol)
2. Dispersion correction is essential
  - Reduces error by 2-10 kcal/mol
3. Def2-TZVPP is very accurate

# PES sampling

1. **Coalescence Kick (CK)** <https://github.com/averkiev75/Coalescence-Kick>

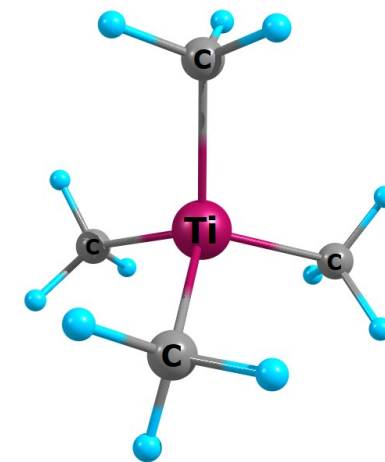
- only  $C_1$  structures
- allows orientation of predefined fragments

2. **AFFCK** <https://github.com/hczhai/AFFCK>

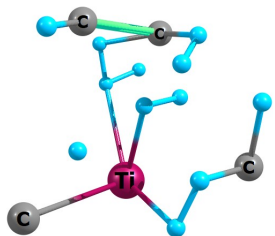
- $C_1$  and symmetrical structures

3. **PyXtal** <https://github.com/qzhu2017/PyXtal>

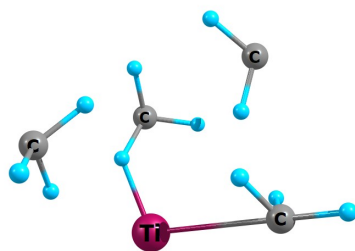
- Python library
- $C_1$  and symmetrical structures



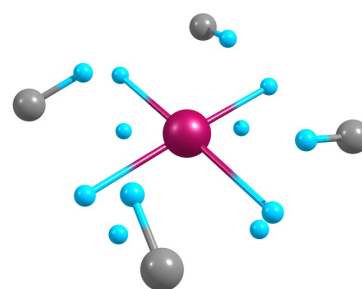
Ti(Me)<sub>4</sub> optimized structure



random  $C_1$



random  $C_1$  (fixed Me fragments)



symmetrical guesses

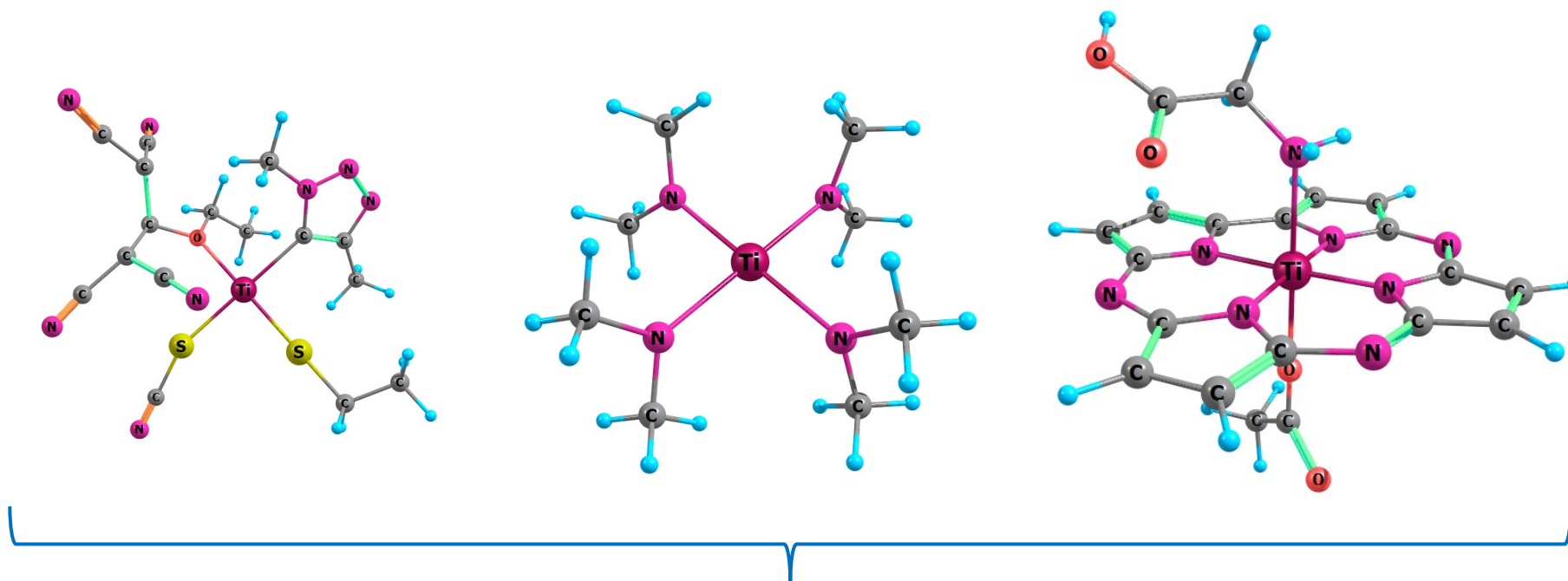
**unphysical structures: broken geometries and wave functions**



# PES sampling

## 4. MolSimplify <http://hjkgrp.mit.edu/molSimplify-tutorials> - designed for organometallics

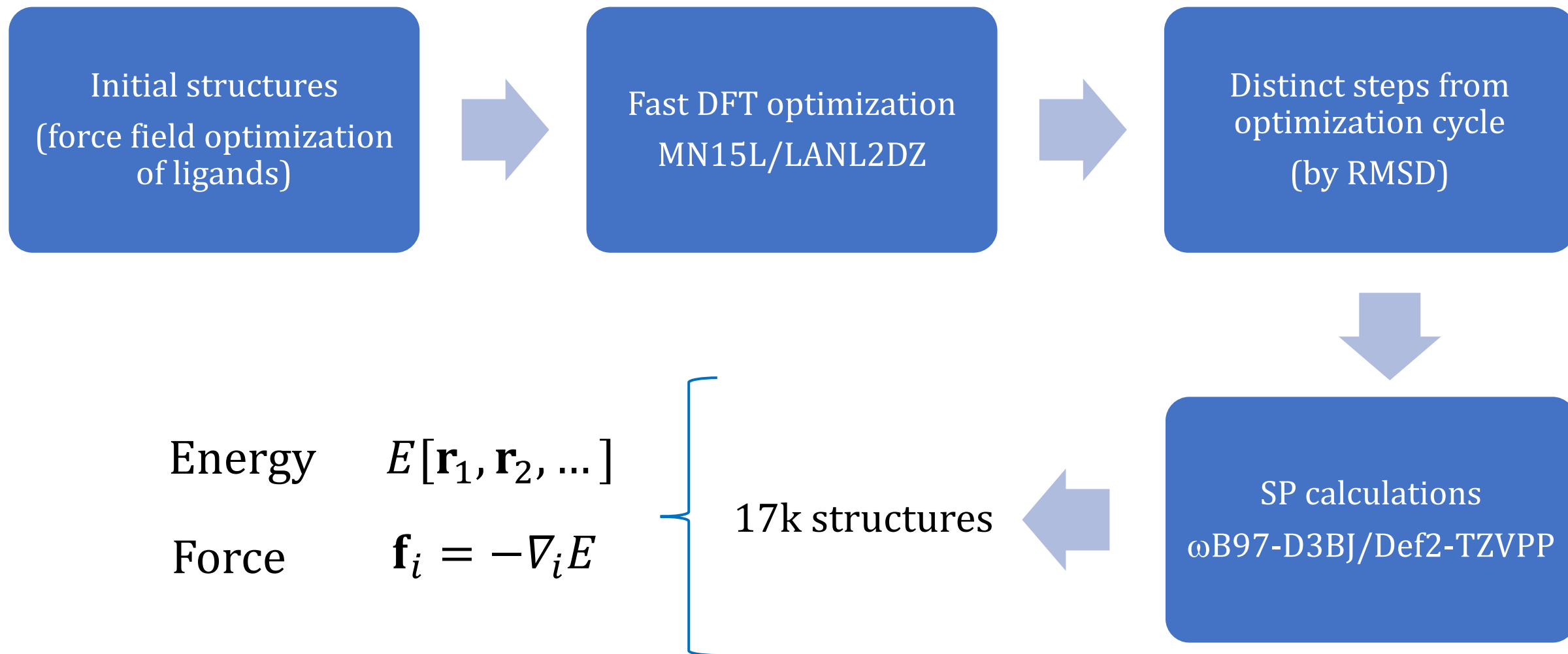
- Different coordination numbers
- Different geometries of complex
- Specific spin state (e. g. closed-shell)
- Supports SMILES



correct connectivity and valence angles

incorrect TM-L distance

# Dataset generation





# HIP-NN potential

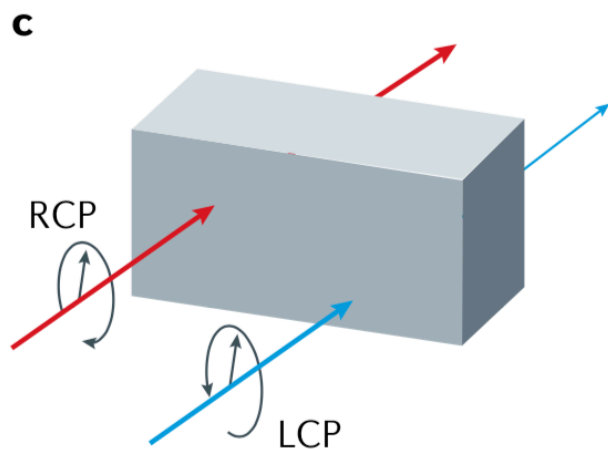
	E, kcal/mol				MAE	RMSE
	train	validation	test			
				B97-3c	<i>a</i>	<i>a</i>
				HF-3c	<i>a</i>	<i>a</i>
				PBEh-3c	<i>a</i>	<i>a</i>
				PM6	21.6	28.9
RMSE	2.23	13.62	14.59			
MAE	1.36	4.60	4.50			
MAE/atom	0.03	0.10	0.94			
				PM7	106.4	146.1
				PM7-TS	68.2	141.9

**Empirical potential with dynamic ML parametrization could outperform statically parametrized SEQM**

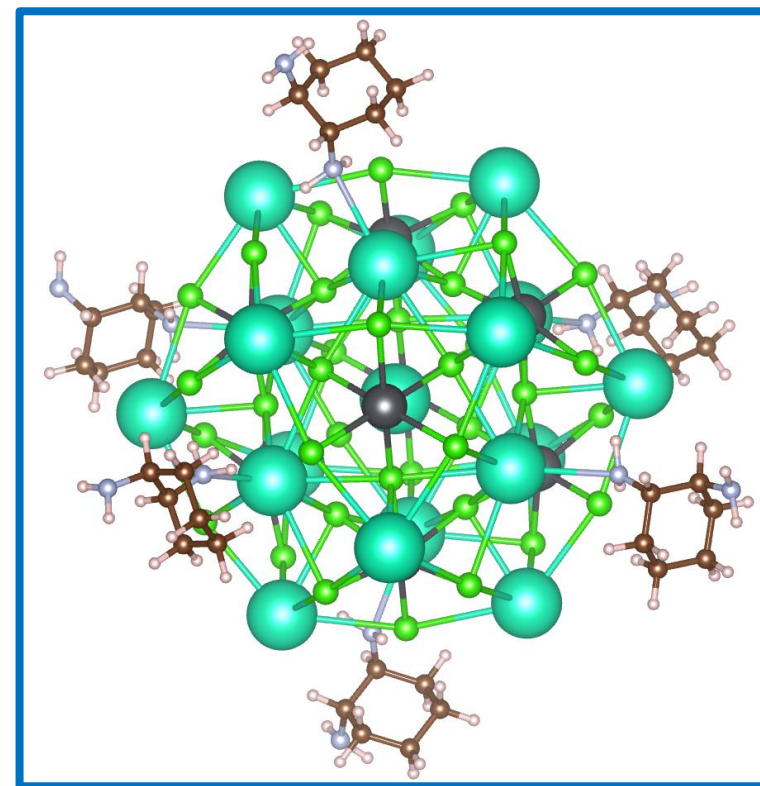
Further work:

- Increase transferability
- Add forces to the model
- Parametrize SEQM on the same dataset

# *Chiroptical Properties Induced into a Lead-Halide Perovskite Cluster Through Surface Chemistry*



Aaron Forde (GRA)  
North Dakota State University  
LANL Lightning Talk – 8-27-2020

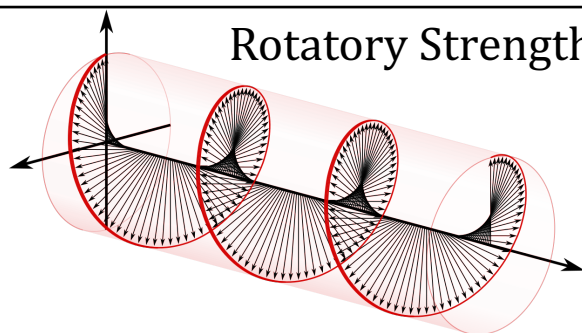


# Light-Matter Interactions: Circular Dichroism

## Classical Electromagnetic Field – Matter Interaction

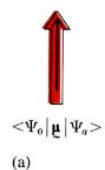
$$\vec{\epsilon} \cdot \vec{d}_{ij} = \left\langle i \left| \exp \left( \frac{i\omega}{c} \vec{n} \cdot \vec{x} \right) \vec{\epsilon} \cdot \vec{p} \right| f \right\rangle \approx \underbrace{\langle i | \vec{\epsilon} \cdot \vec{p} | j \rangle}_{\text{Electric Dipole}} + \underbrace{\langle i | (\vec{n} \cdot \vec{x}) (\vec{\epsilon} \cdot \vec{p}) | j \rangle}_{\text{Magnetic Dipole and Electric Quadrupole}} + \dots$$

## Rotatory Strength

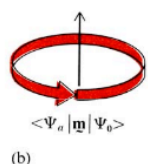


$$R_{0a} = \text{imag}(\langle \Psi_0 | \vec{\mu} | \Psi_a \rangle \cdot \langle \Psi_a | \vec{m} | \Psi_0 \rangle)$$

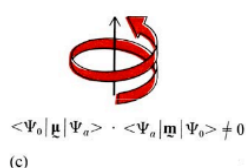
Electric Dipole      Magnetic Dipole



Pure  
electronic  
absorption



Pure  
magnetic  
absorption



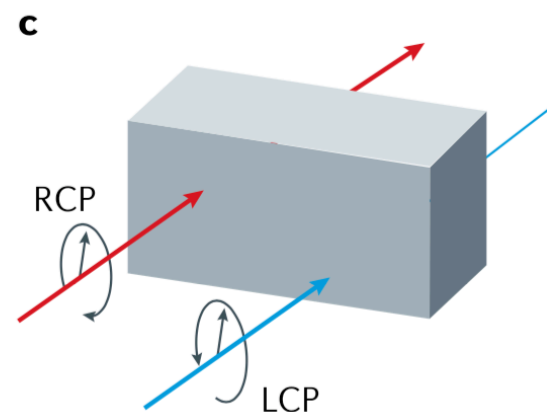
**Optical  
activity!**

Chiroptical  
Activity

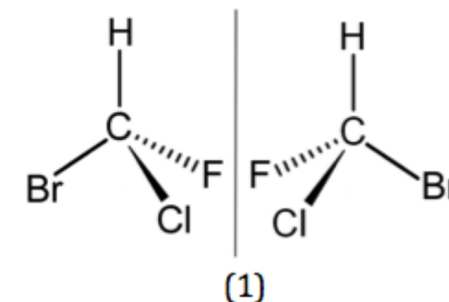
## Circular Dichroism

$$A_{RCP} \neq A_{LCP}$$

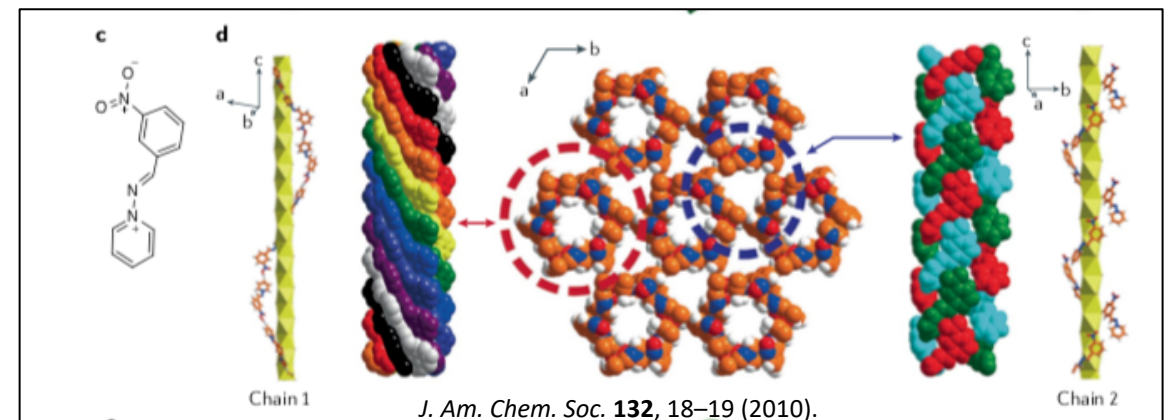
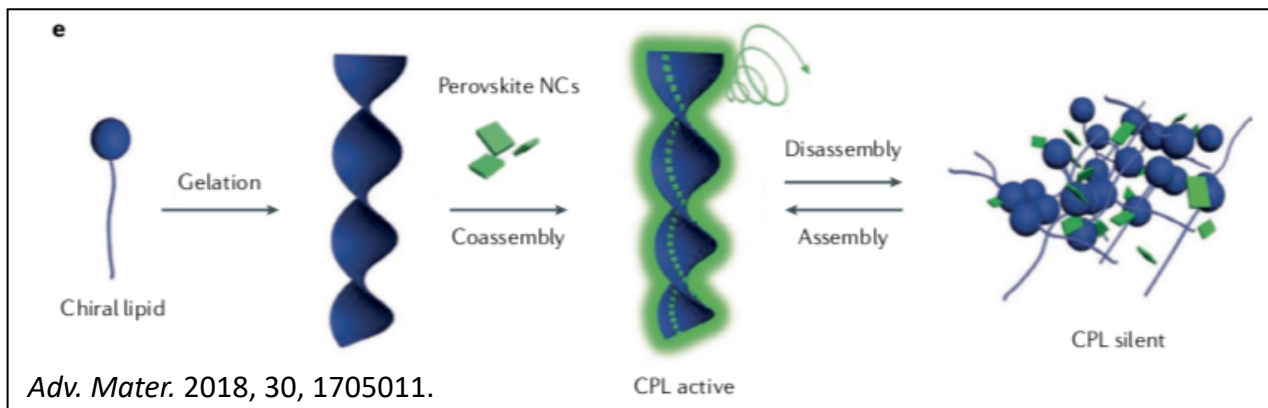
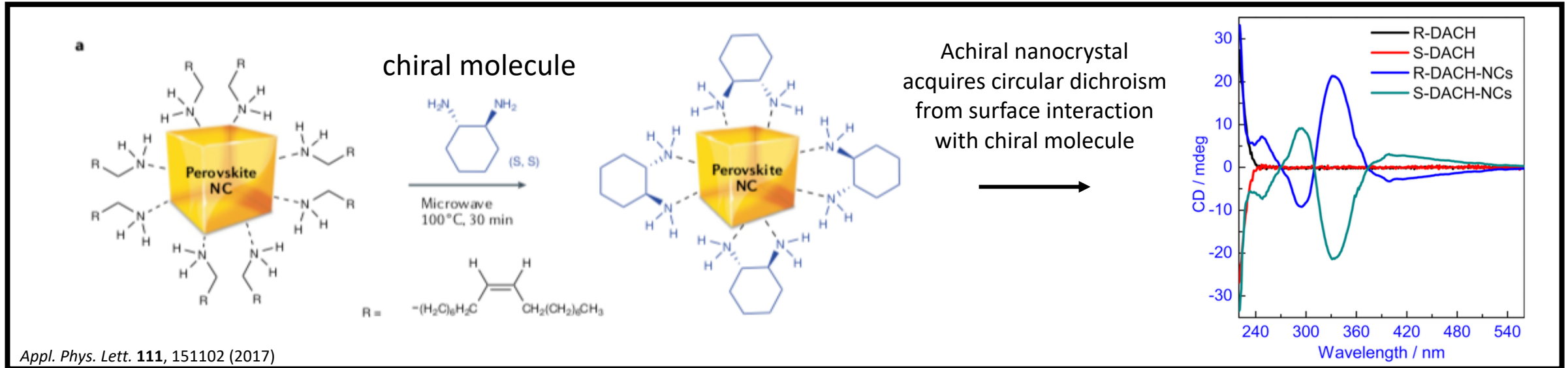
Difference in absorption of  
right-handed and left-handed  
circularly polarized light



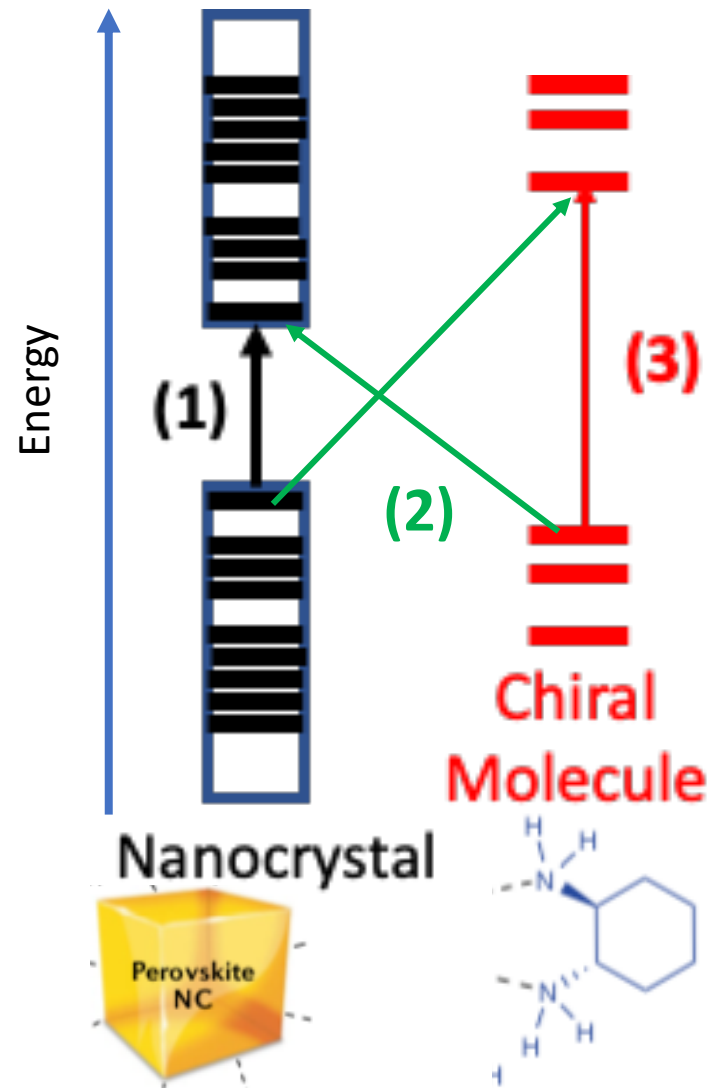
Occurs in materials  
that do not have  
mirror symmetry



# Chirality Induced into Lead-Halide Perovskite Semiconductors Through Surface Interactions



# What is the Photo-Physical Mechanism for Induced Chirality in Lead Halide Perovskite Nanocrystal?



## Possible Radiative Pathways

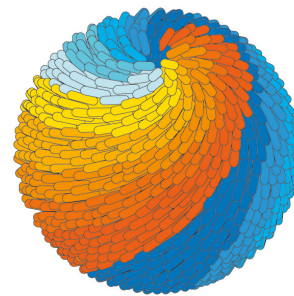
(1) Perovskite-to-Perovskite

(2) Charge-Transfer Excitations

(3) Molecule-to-Molecule

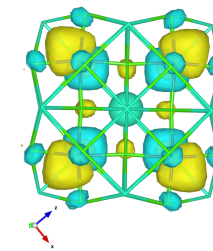
## (1) Chiral Templating

Chiral templating of the nanocrystal surface gives circular dichroism for perovskite-to-perovskite transitions

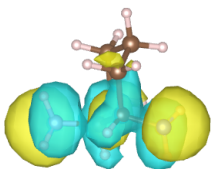


## (2) Charge-Transfer Excitations

Electronic coupling between the nanocrystal and the chiral molecule give charge-transfer excitations that have chiral signatures.



(2)





# TD-DFT Investigation of Circular Dichroism in Chiral Capped Lead Halide Perovskite Cluster

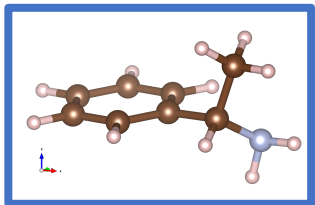
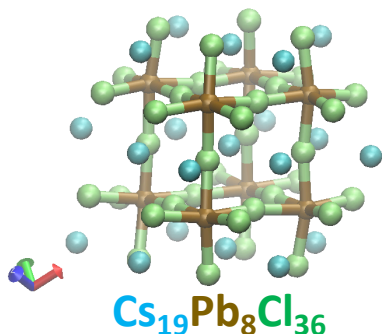
Functional: CAM-B3LYP

Basis Set for Cluster: LANL2DZ

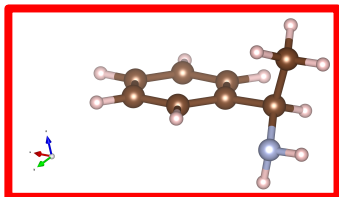
Basis Sets for Molecule: 6-311++g\*\*

Solvent: Toluene

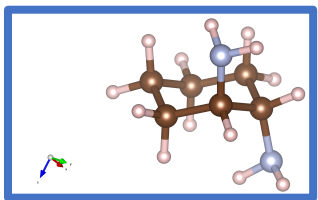
## (i) Individual Components



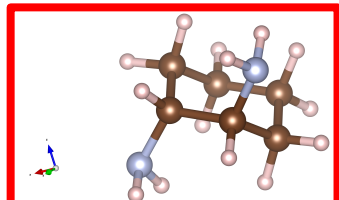
1R Methylbenzylamine



1S Methylbenzylamine

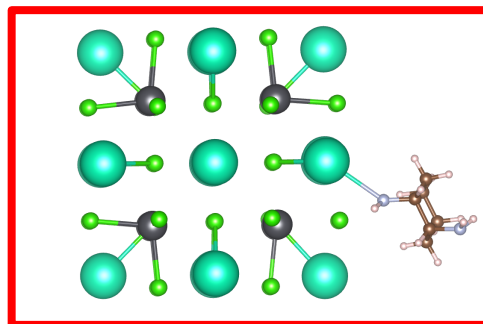


1R,2R Diaminocyclohexane

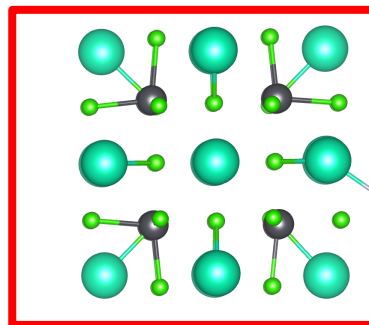


1S,2S Diaminocyclohexane

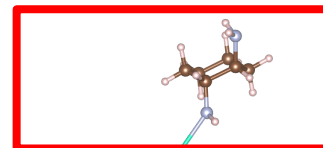
## (ii) Single Molecule on Cluster Surface



$\text{Cs}_{19}\text{Pb}_8\text{Cl}_{36}$  + 1S,2S Diaminocyclohexane

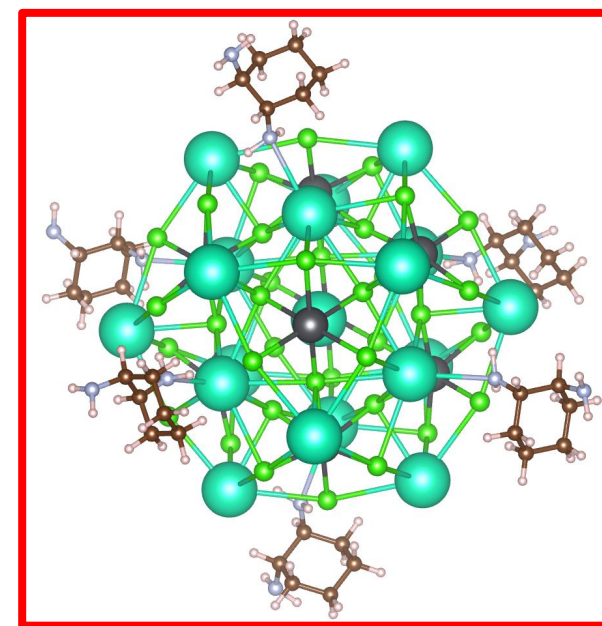


$\text{Cs}_{19}\text{Pb}_8\text{Cl}_{36}$



1S,2S Diaminocyclohexane

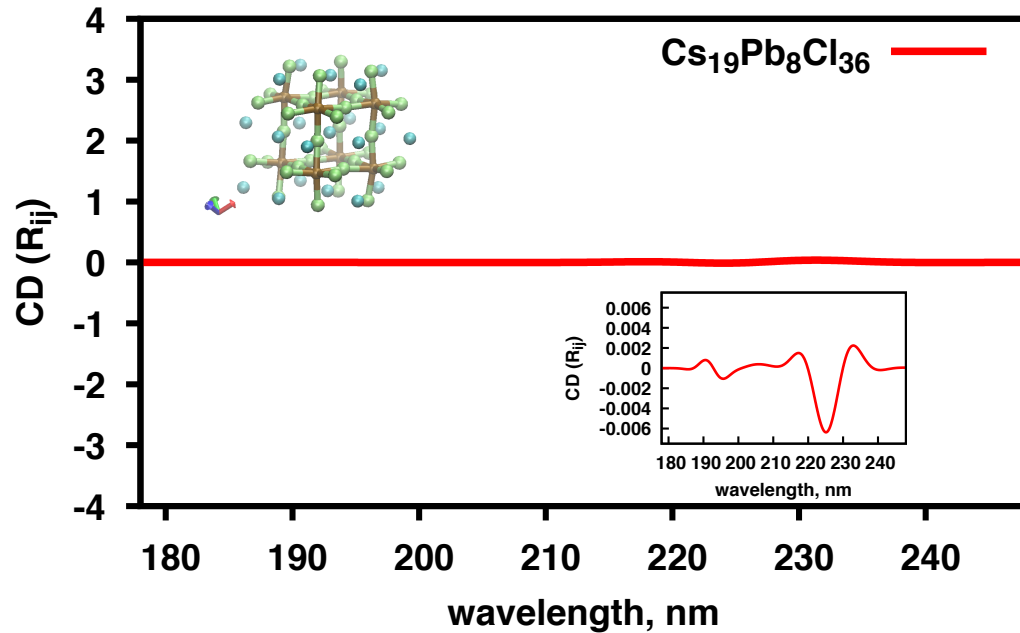
## (iii) Multiple Molecules on Cluster Surface



$\text{Cs}_{19}\text{Pb}_8\text{Cl}_{36}$  1S,2S Diaminocyclohexane

# (i) Circular Dichroism of Individual Components

Circular Dichroism of Perovskite Cluster



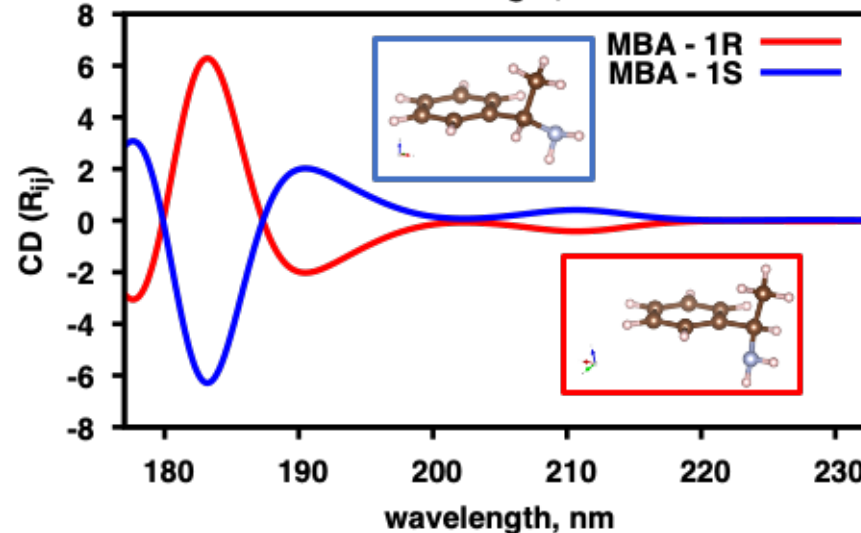
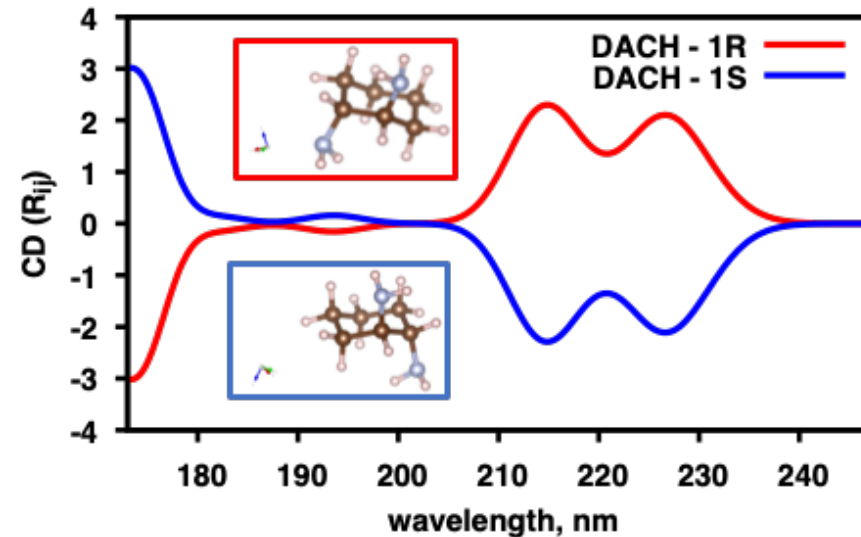
Cluster shows negligible chiroptical activity

$$CD(\omega_{ij}) = \sum_{ij} R_{ij} \delta(\hbar\omega - \hbar\omega_{ij})$$

$$R_{ij} \propto \text{Im} \left( \left| \vec{\mu}_{ij}^{elec} \cdot \vec{\mu}_{ij}^{mag} \right| \right)$$

$\delta$  is broadened as a Gaussian with broadening parameter 0.01 eV

Circular Dichroism of Chiral Molecules



Enantiomers show expected mirrored CD spectra

# (ii) Circular Dichroism with a Single Chiral Molecule on the Cluster Surface

**Cs<sub>19</sub>Pb<sub>8</sub>Cl<sub>36</sub> 1R,2R Diaminocyclohexane**

	binding [eV]
Configuration 1	-0.46
Configuration 2	-0.67
Configuration 3	-0.57
Configuration 4	-0.58

**Cs<sub>19</sub>Pb<sub>8</sub>Cl<sub>36</sub> 1R Methylbenzlamine**

	binding [eV]
Configuration 1	-0.45
Configuration 2	-0.53
Configuration 3	-0.44
Configuration 4	-0.48

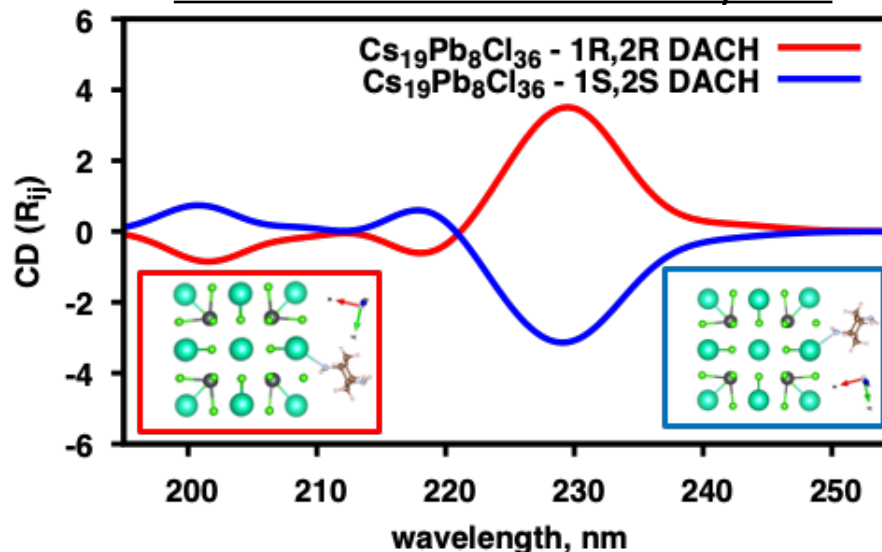
$$E_{\text{binding}} = E(\text{cluster} + \text{molecule}) - [E(\text{molecule}) + E(\text{Cluster})]$$

$$CD(\omega_{ij}) = \sum_{ij} R_{ij} \delta(\hbar\omega - \hbar\omega_{ij})$$

$$R_{ij} \propto \text{Im} \left( \left| \vec{\mu}_{ij}^{\text{elec}} \cdot \vec{\mu}_{ij}^{\text{mag}} \right| \right)$$

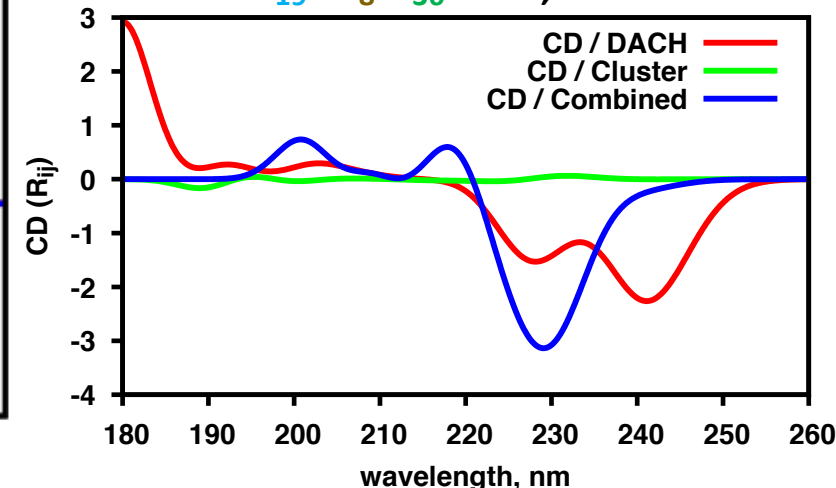
$\delta$  is broadened as a Gaussian with broadening parameter 0.01 eV

Circular Dichroism of Joint System

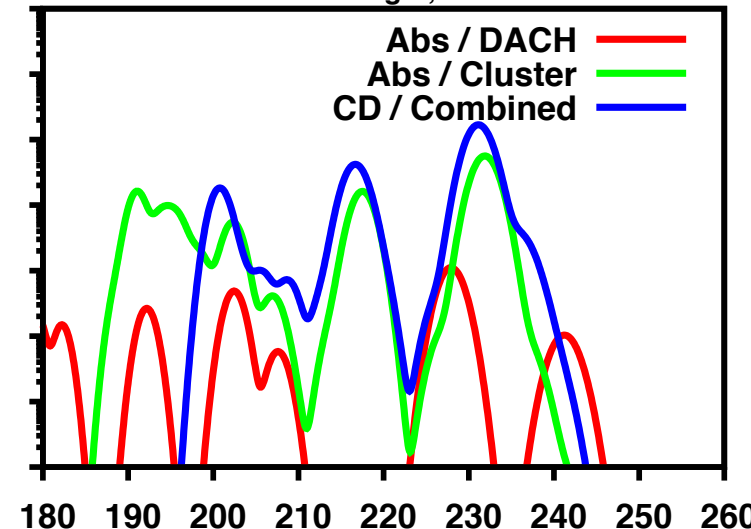


'Decomposed' System:

**Cs<sub>19</sub>Pb<sub>8</sub>Cl<sub>36</sub> 1S,2S DACH**

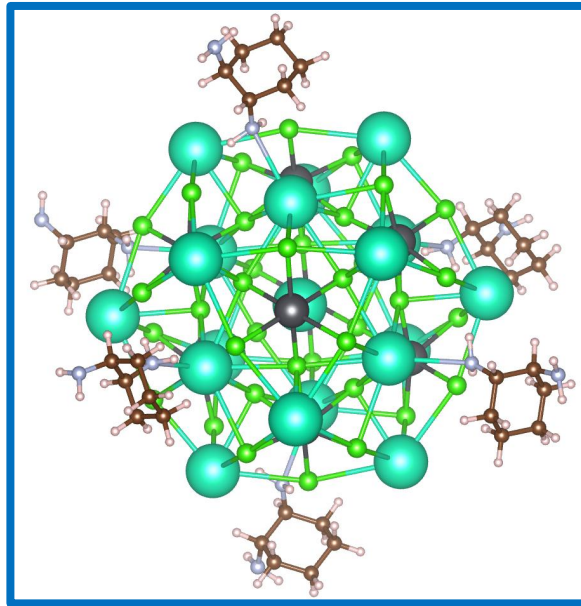
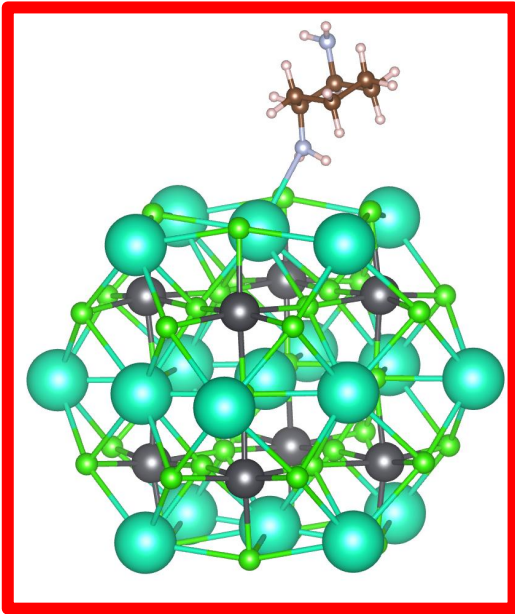


CD spectra for the joint system is distinct from individual contributions

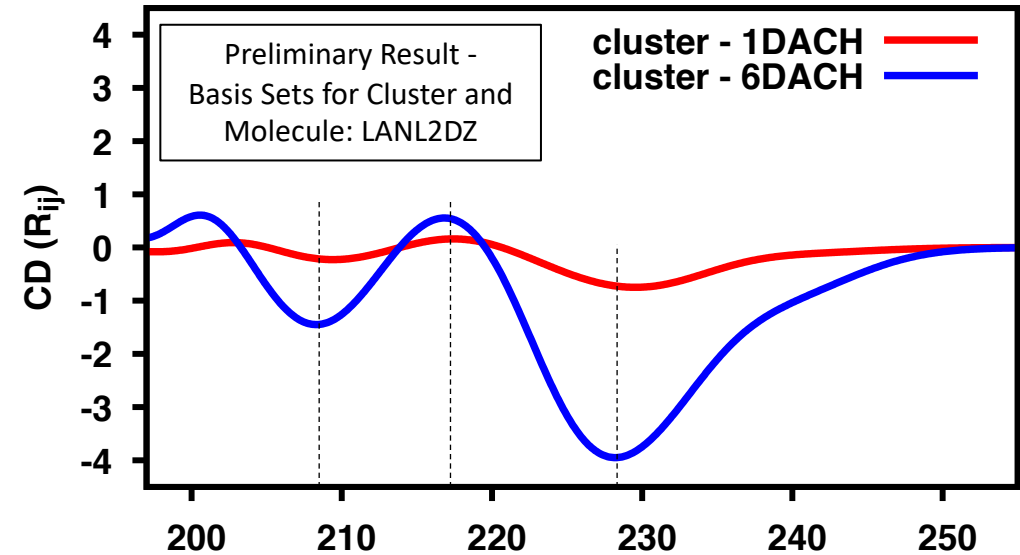




# (iii) Circular Dichroism with Multiple Chiral Molecules on the Cluster Surface



Circular Dichroism of With Chiral Molecules on the Surface



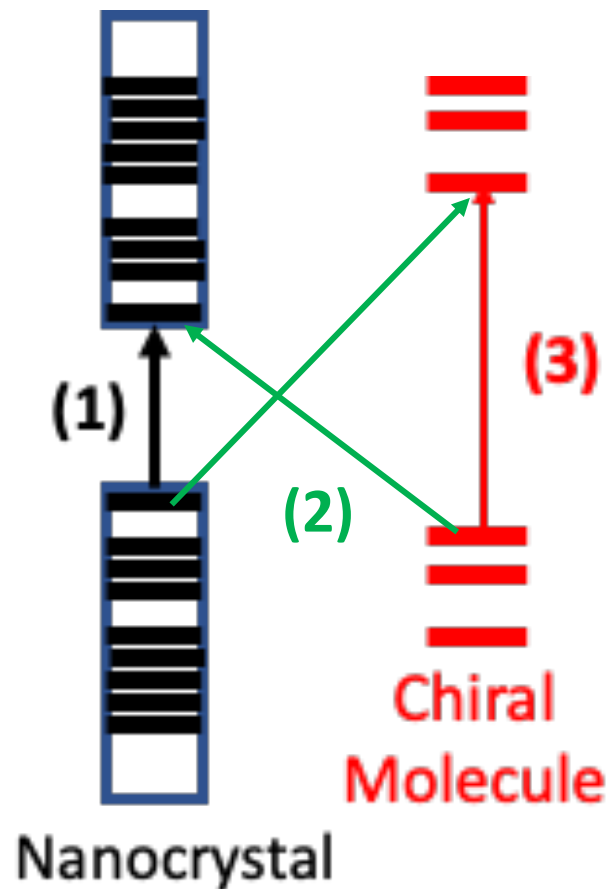
$$CD(\omega_{ij}) = \sum_{ij} R_{ij} \delta(\hbar\omega - \hbar\omega_{ij})$$

$$R_{ij} \propto \text{Im} \left( \left| \vec{\mu}_{ij}^{elec} \cdot \vec{\mu}_{ij}^{mag} \right| \right)$$

$\delta$  is broadened as a Gaussian with broadening parameter 0.01 eV

Effect of having multiple chiral ligands on the surface is additive to the CD spectra

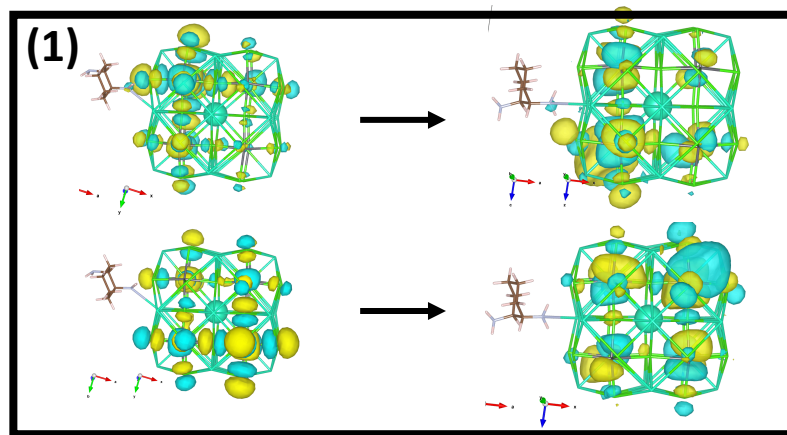
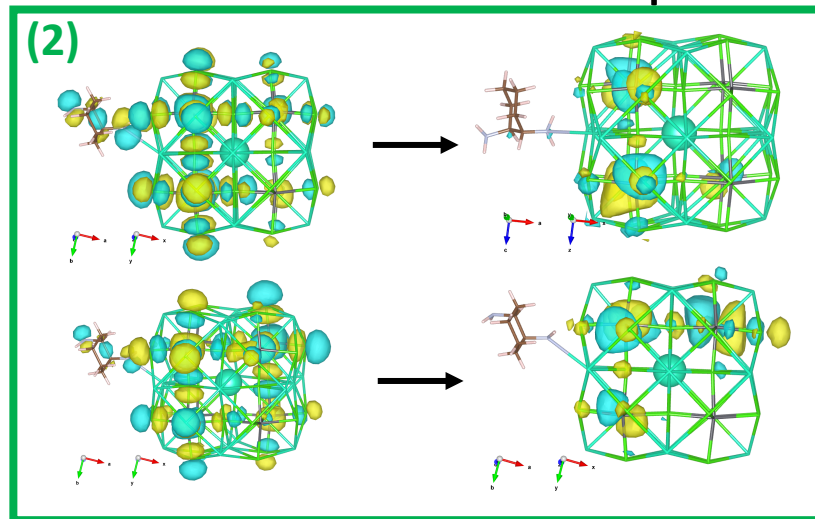
# Contribution of Charge-Transfer Excitations Towards Circular Dichroism Spectra



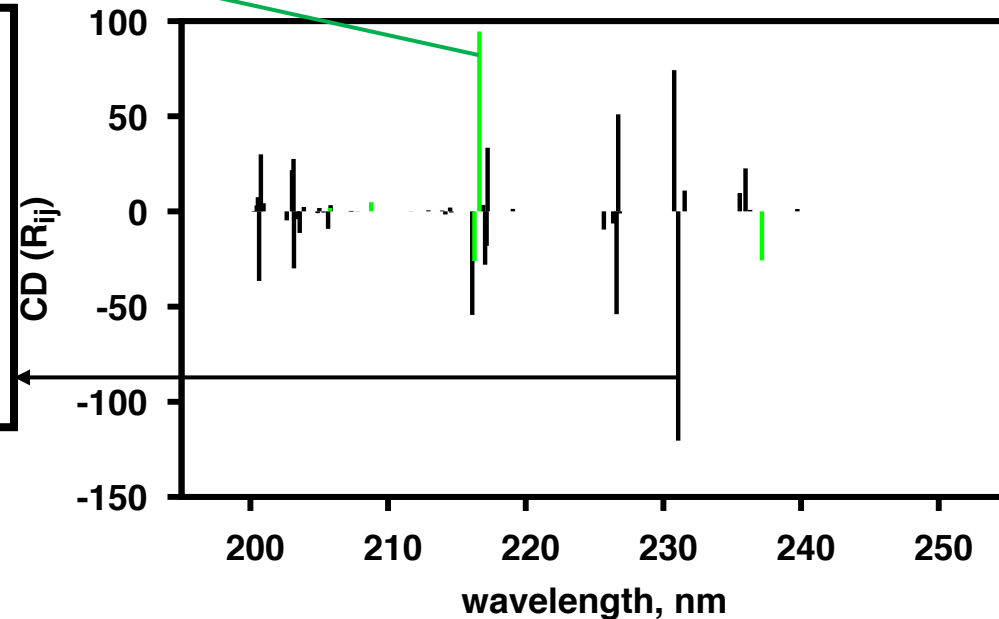
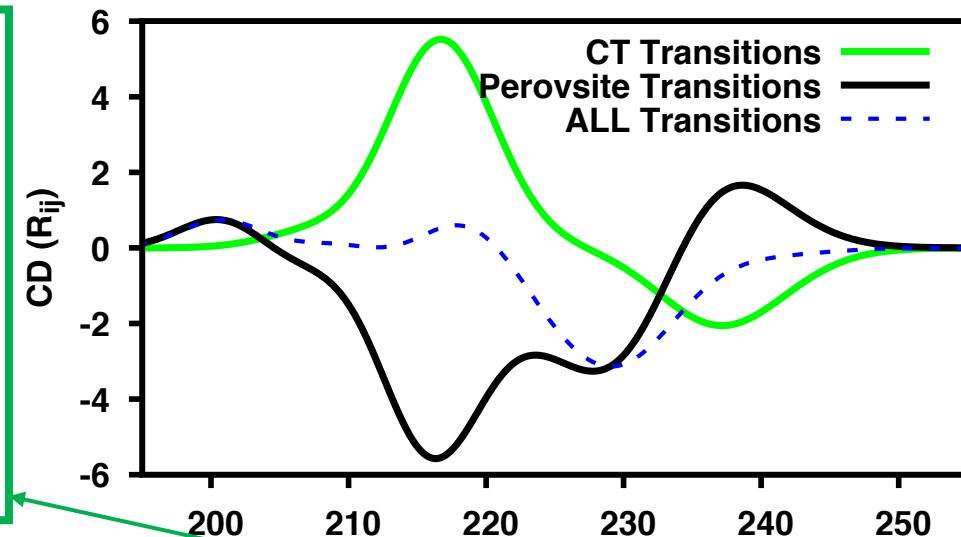
$$CD(\omega_{ij}) = \sum_{ij} R_{ij} \delta(\hbar\omega - \hbar\omega_{ij})$$

$$R_{ij} \propto \text{Im} \left( \left| \vec{\mu}_{ij}^{elec} \cdot \vec{\mu}_{ij}^{mag} \right| \right)$$

$\delta$  is broadened as a Gaussian with broadening parameter 0.01 eV



Additional features in CD spectra are not only attributable to (2) Charge-Transfer Excitations



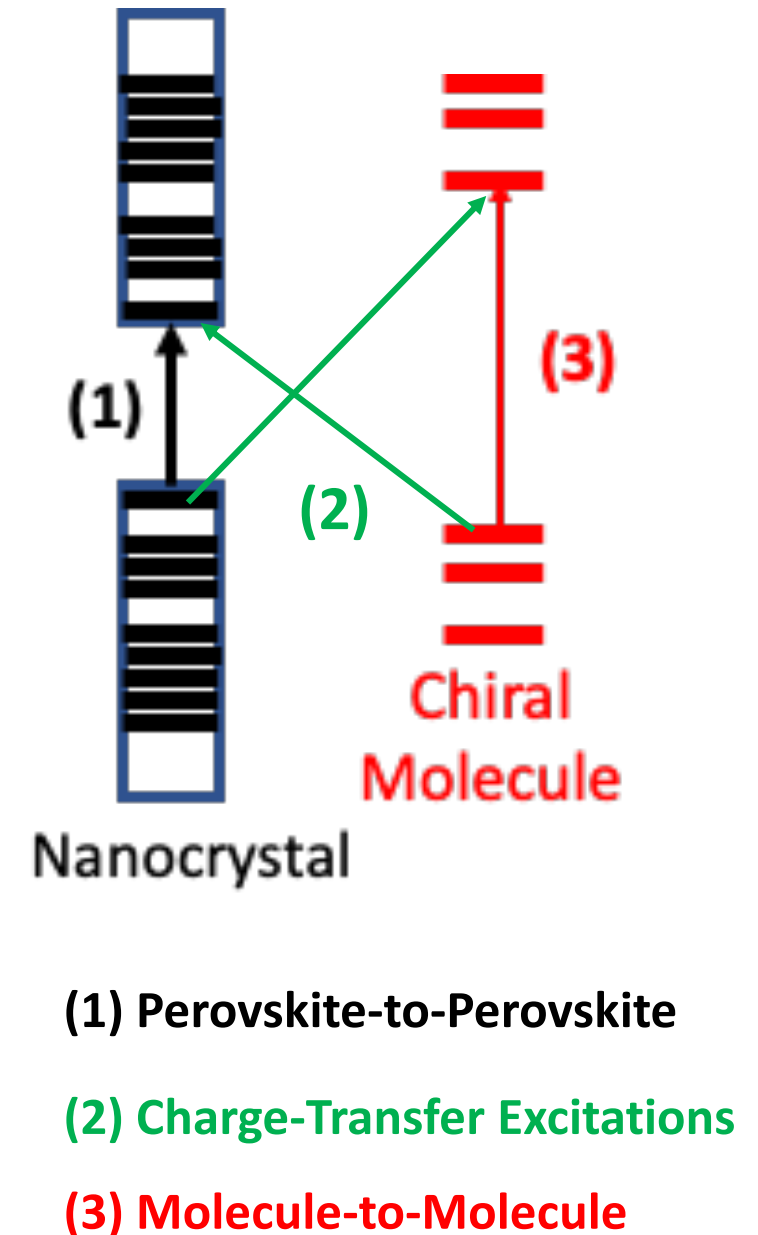
# Key Findings:

(i) Circular Dichroism spectra for the joint system is distinct from individual contributions

- Total CD spectra is not sum of radiative pathways  
**(1) + (2) + (3)**

(ii) Effect of having multiple chiral ligands on the surface is additive to the CD spectra

(iii) Chemical interaction between the chiral molecule and achiral perovskite imprints chirality into **(1)** radiative pathways



# Acknowledgements

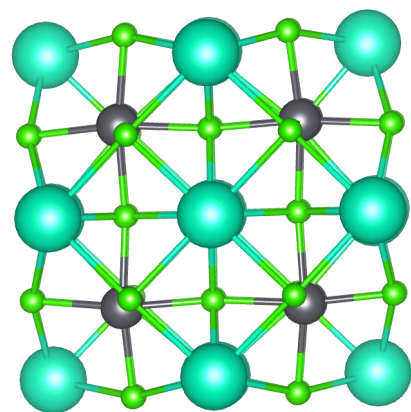
- LANL Mentors

- Amanda Neukirch
- Dibya Ghosh
- Sergei Tretiak

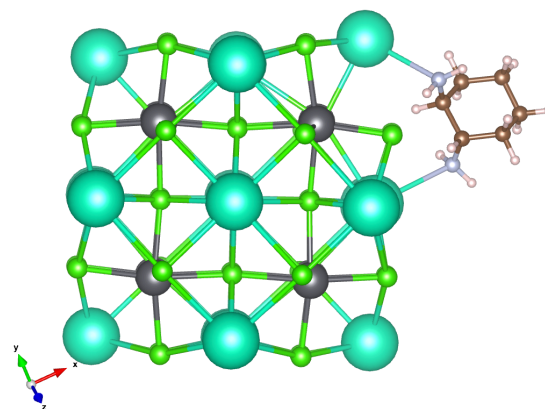
- NDSU

- Dmitri Kilin
- Svetlana Kilina
- Erik Hobbie

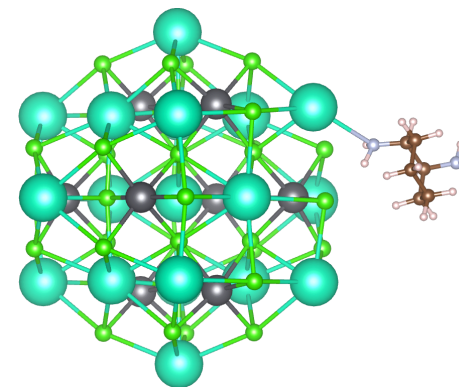
Supplemental

**(a)**

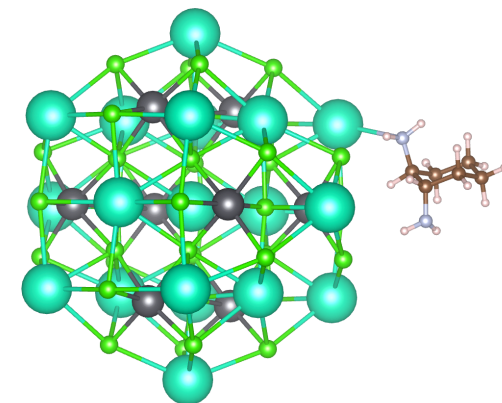
$\langle 100 \rangle$   
Configuration 1

**(b)**

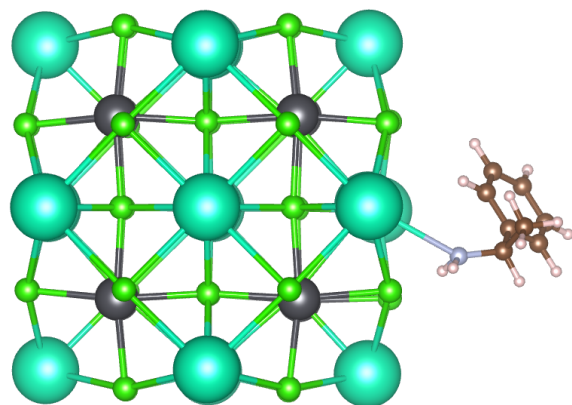
$\langle 100 \rangle$   
Configuration 2

**(c)**

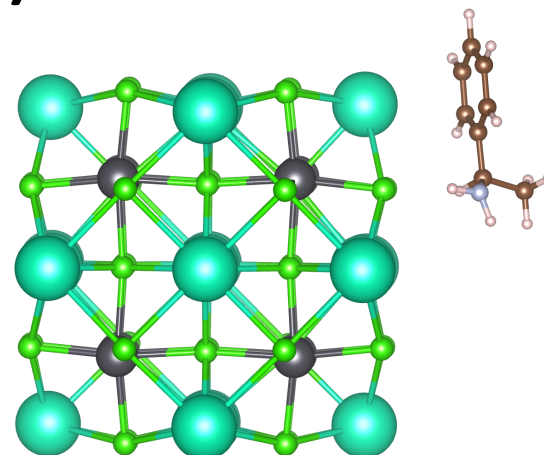
$\langle 111 \rangle$   
Configuration 3

**(d)**

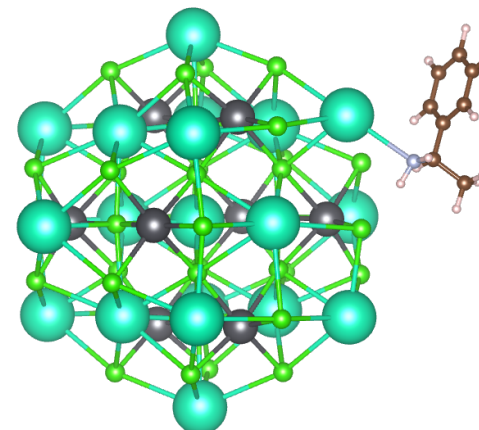
$\langle 111 \rangle$   
Configuration 4

**(e)**

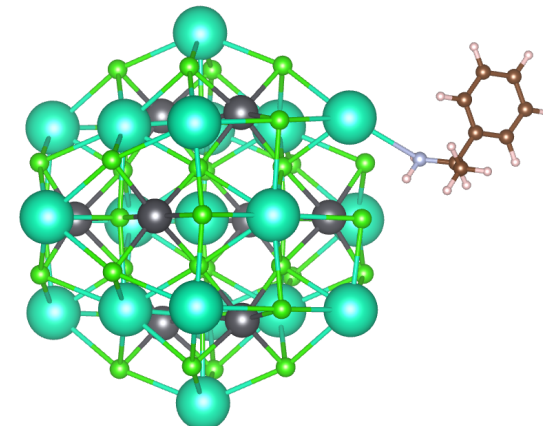
$\langle 100 \rangle$   
Configuration 1

**(f)**

$\langle 100 \rangle$   
Configuration 2

**(g)**

$\langle 111 \rangle$   
Configuration 3

**(h)**

$\langle 111 \rangle$   
Configuration 4

# Uncertainty quantification of atomic data for collisional-radiative modeling of fusion plasmas

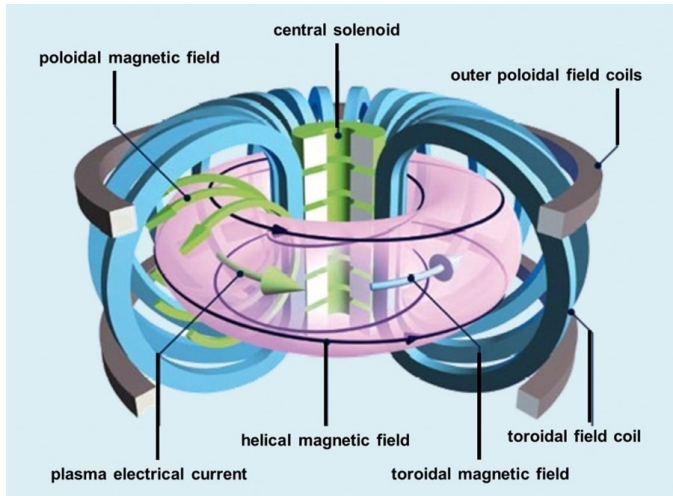
Andrew George – Student intern at T-5 with Willem Kupets at T-1

Mentors: Nathan Garland (T-5) and Mark Zammit (T-1)

LA-UR-20-25815, Funding through ER Grant  
20200356ER

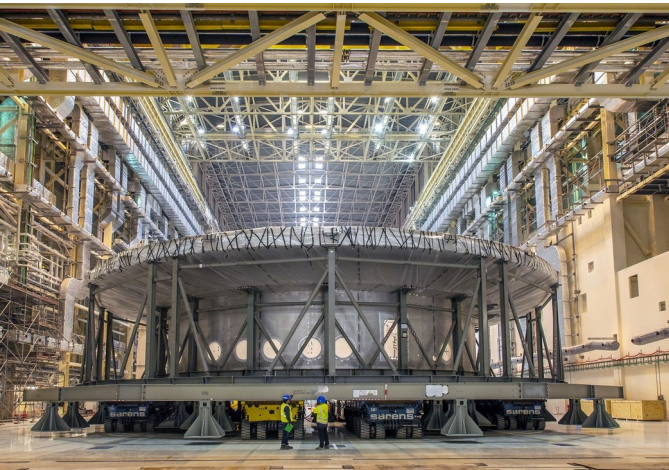


# The Tokamak



A tokamak is a magnetic confinement device that can confine plasma using counter propagating currents. It is a solenoid that uses the magnetic field that is generated by the electric field to pinch the plasma and confine it.

[1]



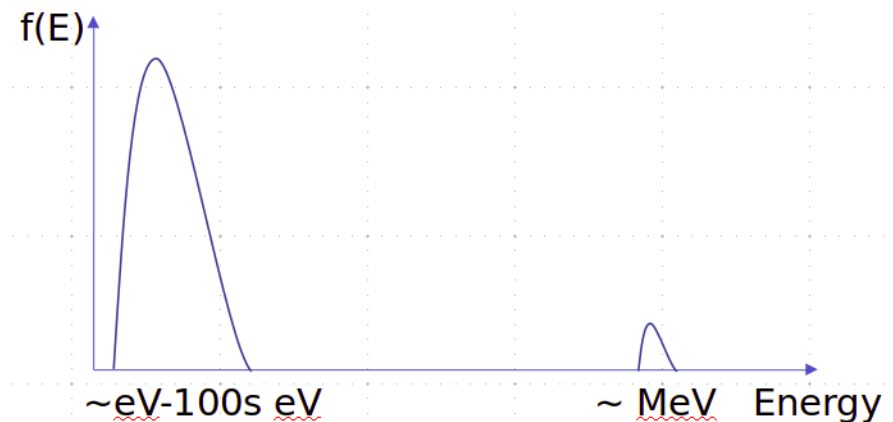
The tokamak that is related to our work is ITER. ITER is the international collaboration that is attempting to make self heating fusion.



# Tokamak Disruption Problem

In a tokamak a factor that needs to be accounted for is the fact that a disruption can happen. A disruption is a plasma instability, and a possible scenario of a disruption are runaway electrons. The runaway electrons cause a beam that hits a part of the tokamak.

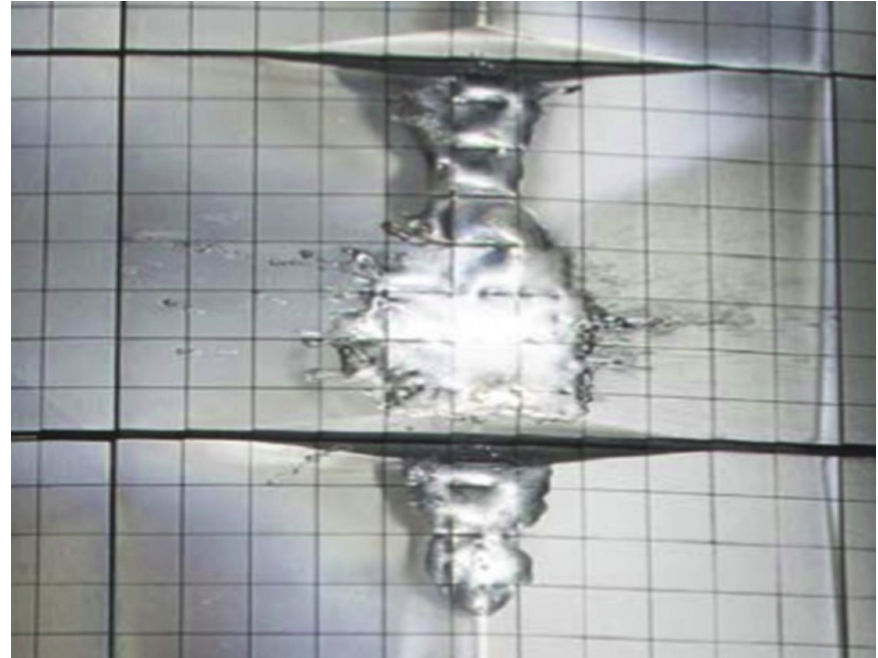
Our team at LANL is studying the underlying physics of disruptions in order to understand how they occur, in order to develop prevention or mitigation options.



# A Potential Solution?

In ITER the walls are mostly beryllium with a small area that is tungsten. Since electrons are highly relativistic and energetic, they need to be mitigated, in order to reduce potential destruction of the vessel.

A high  $z$  material such as argon or neon can be inserted to act as an energy sink for the runaway electrons. Therefore we need to understand the composition of ions in the plasma.



Damaged wall

# Modeling the Plasma- CR Model

A collisional radiative model is a model that describes the collisional events and radiative processes (I.E. the up and down processes).

The collisional radiative model needs two things – state completeness and the appropriate rate matrix.

We are using a superconfiguration refinement which uses the principle quantum  $n$  for our states.

Quantum number	Name	What it labels	Possible values	Notes
$n$	principal	electron energy level or shell number	1, 2, 3, ...	Except for d-orbitals, the shell number matches the row of the periodic table.
$\ell$	azimuthal	orbital type: s, p, d, f	0, 1, 2, ..., $n-1$	0 = s orbital 1 = p orbital 2 = d orbital 3 = f orbital
$m_\ell$	magnetic	orbital sub-type	integers between and including $-\ell$ and $+\ell$ : $-\ell, -\ell+1, \dots, \ell-1, \ell$	$\ell = 0$ (s): 2 $e^-$ in one orbital $\ell = 1$ (p): 2 $e^-$ in each of three sub orbitals ( $p_x, p_y, p_z$ ) $\ell = 2$ (d): 2 $e^-$ in each of 5 sub orbitals ( $d_{xy}, d_{xz}, d_{yz}, d_{x^2-y^2}, d_{z^2}$ )
$m_s$	spin	electron spin	$\pm \frac{1}{2}$	Spins in any single sub-orbital must be paired.

$$\vec{J} = \vec{L} + \vec{S}$$

# My Project - The Rate Matrix

The Rate matrix is the matrix that contains the collision and radiative transitions.

The rates coefficients are first found through the cross sections. The coefficients can then be used to create the rate matrix.

Using a code called FLYCHK and then an uncertainty quantification package called Dakota I'll be able to take an iterative approach to analyze the effect of uncertainty of the cross sections.

$$\frac{dn_i}{dt} = -n_i \sum_{j \neq i}^{N_L} W_{ij} + \sum_{j \neq i}^{N_L} n_j W_{ji} \quad 1 \leq i \leq N_L$$

For upward transitions ( $i < j$ ),

$$W_{ij} = B_{ij} \bar{J}_{ij} + n_e C_{ij} + \beta_{ij} + n_e \gamma_{ij} + \sigma_{ij} + I_{ij}$$

For downward transitions ( $j > i$ )

$$W_{ij} = A_{ji} + B_{ji} \bar{J}_{ji} + n_e D_{ji} + n_e \alpha_{ji}^{RR} + n_e \kappa_{ji}^{EC} + n_e^2 \delta_{ji}$$

where  $n_e$  is an electron density and  $\bar{J}$  is a frequency-averaged mean intensity which is relevant to the transition. The subscript  $ji$  refers to a transition from level  $j$  to level  $i$ . The rate coefficients correspond to the following atomic processes:

$A_{ji}$ spontaneous emission	$\alpha_{ji}$ radiative recombination
$B_{ij}$ stimulated absorption ( $i < j$ ) or emission ( $i > j$ )	$\beta_{ij}$ photoionization plus stimulated recombination
$C_{ij}$ collisional excitation	$\gamma_{ij}$ collisional ionization
$D_{ij}$ collisional deexcitation	$\delta_{ji}$ collisional recombination
$I_{ij}$ beam and non-thermal electron collisions	$\kappa_{ji}$ electron capture
	$\sigma_{ij}$ autoionization

$$dn/dt = Rn$$

# What is FLYCHK

- Upgrade from FLY which is a K-shell based code designed to study hot plasmas
- FLYCHK is a general purpose plasma code that applies to most laboratory plasmas (non-LTE, LTE, collisionally driven etc).
- Provides  $\langle Z \rangle$  for fixed densities: ion, electron, mass
- Uses the principle quantum number  $n$  (superconfiguration) to describe ion energy levels[2]

# Dakota UQ package

“At a high level, uncertainty quantification (UQ) or nondeterministic analysis is the process of (1) characterizing input uncertainties, (2) forward propagating these uncertainties through a computational model, and (3) performing statistical or interval assessments on the resulting responses.”[3]

Dakota also does sensitivity analysis. Sensitivity analysis goes hand in hand with the variance.

Uses Latin Hypercube Sampling(LHS) – Creates near random sample of parameter values from a multidimensional distribution. Dakota samples specifically from a Latin square which in arbitrary dimensions is a hypercube.

# Our UQ approach

Inherent uncertainty exists in approximate formulas for cross-sections and rate input data.

We want to probe this uncertainty and its impact on our outputs.

We vary FLYCHK fundamental atomic data by 25%.

$$R_{\text{emission}}^{UQ} = x_{\text{arat}} A_{ij}$$

$$\sigma_{\text{excitation}}^{UQ} = x_{\text{exc}} \sigma_{\text{exc}}$$

$$0.75 \leq x_i \leq 1.25$$

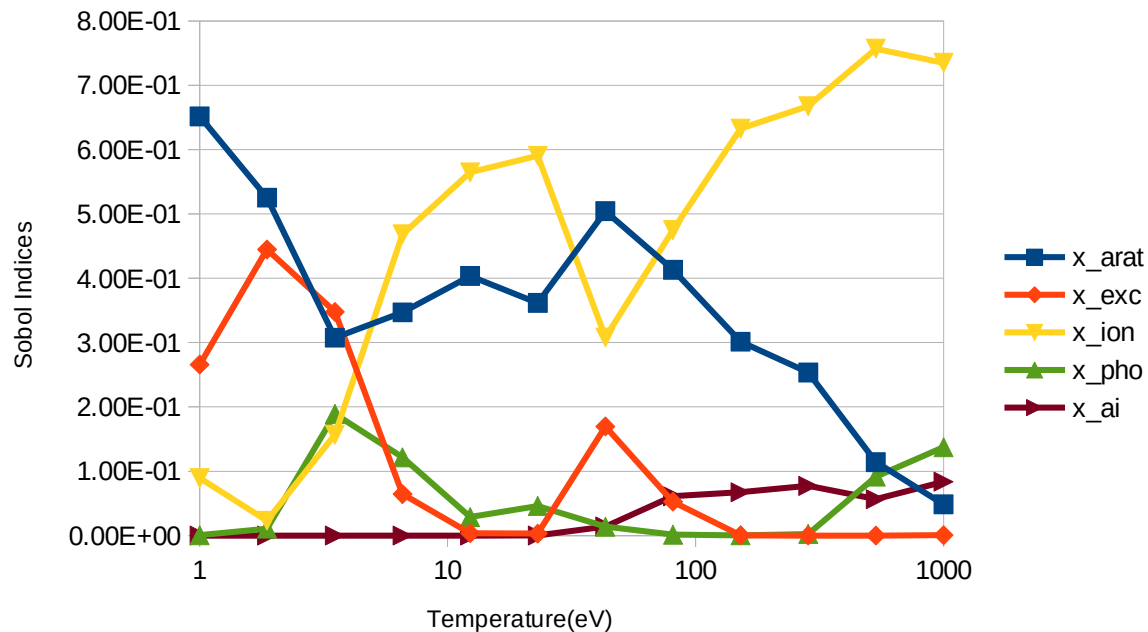
$$\sigma_{\text{ionization}}^{UQ} = x_{\text{ion}} \sigma_{\text{ion}}$$

$$\sigma_{\text{photon}}^{UQ} = x_{\text{pho}} \sigma_{\text{pho}}$$

$$R_{\text{autoion}}^{UQ} = x_{\text{ai}} R_{\text{ai}}$$



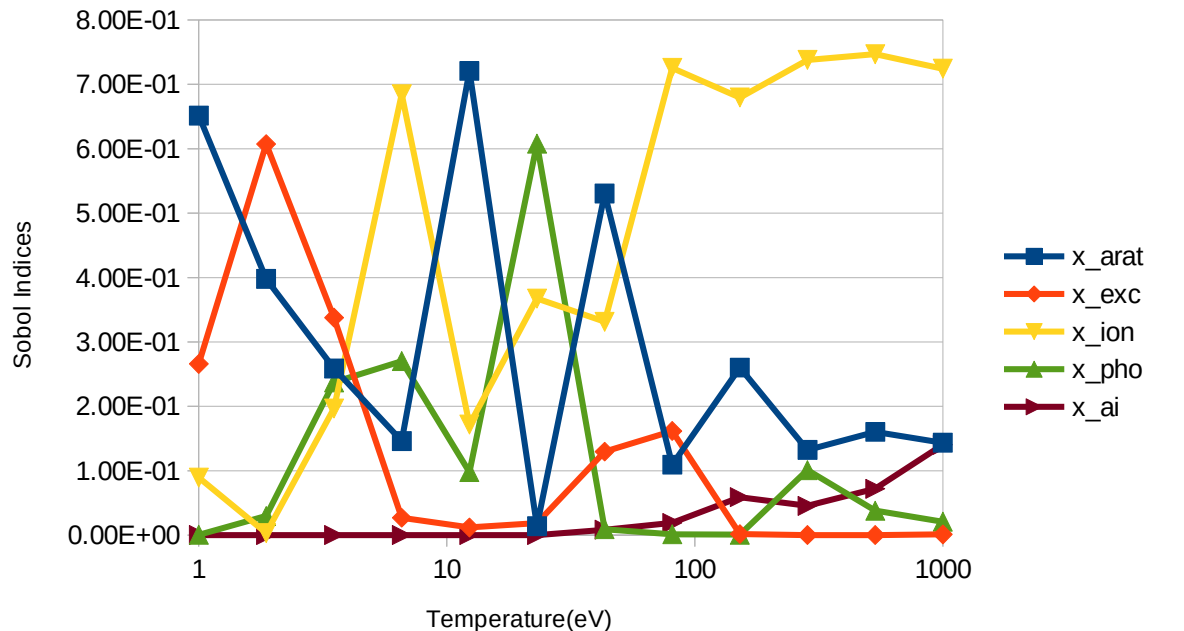
# Results: $\langle Z \rangle$



Surprisingly, excitation and spontaneous emission are the dominant processes at low temperatures,

At higher temperatures ionization becomes the dominant effect as one should expect,

# Results: $Z\text{var} < \sigma_z^2 >$

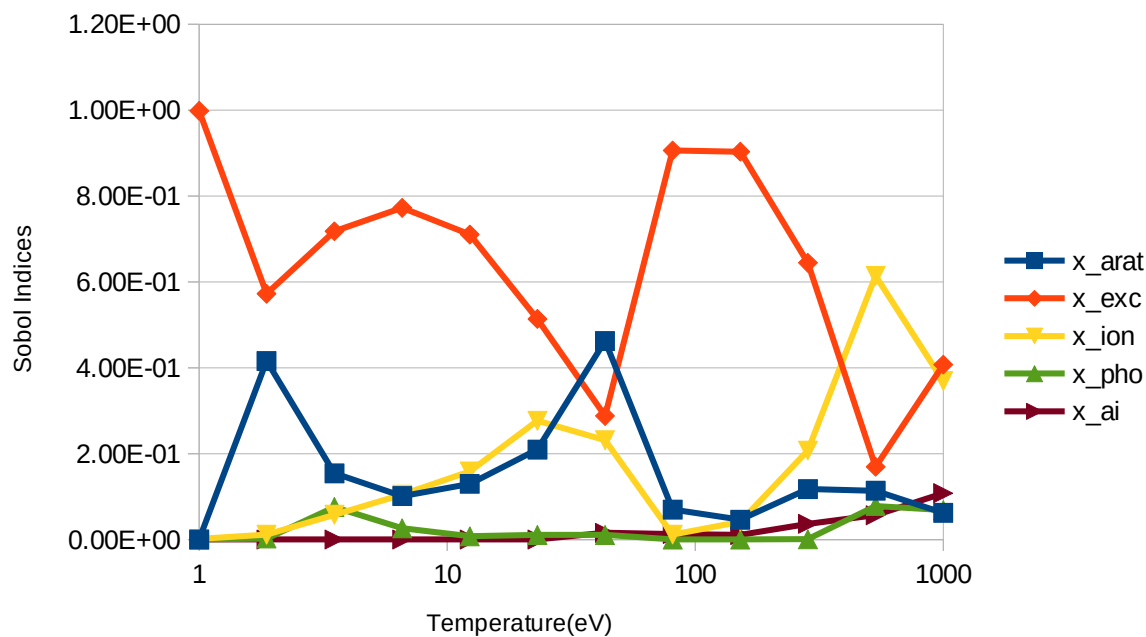


Below 10eV excitation, spontaneous emission and ionization have the biggest effect.

From 10eV to 100 eV, ionization has less of an effect. Spontaneous emission and photoionization are the dominant effects.

At high temperatures ionization is the dominant effect.

# Results: RPL(Radiative Power Loss)



Excitation is the dominant effect at low temperatures.

At 100eV excitation is the dominant effect, but loses its impact as ionization begins to have more of an effect.

# References

- [1]Haupt, K. (2018, June 11). Fusion machines: Searching for the perfect shape. Retrieved July 02, 2020, from <https://www.iter.org/newsline/-/3037>
- [2]Chung, H., Chen, M., Morgan, W., Ralchenko, Y., & Lee, R. (2005). FLYCHK: Generalized population kinetics and spectral model for rapid spectroscopic analysis for all elements. High Energy Density Physics, 1(1), 3-12. doi:10.1016/j.hedp.2005.07.001
- [3]Adams, B.M., Bohnhoff, W.J., Dalbey, K.R., Ebeida, M.S., Eddy, J.P., Eldred, M.S., Geraci, G., Hooper, R.W., Hough, P.D., Hu, K.T., Jakeman, J.D., Khalil, M., Maupin, K.A., Monschke, J.A., Ridgway, E.M., Rushdi, A.A., Stephens, J.A., Swiler, L.P., Vigil, D.M., Wildey, T.M., and Winokur, J.G., "Dakota, A Multilevel Parallel Object-Oriented Framework for Design Optimization, Parameter Estimation, Uncertainty Quantification, and Sensitivity Analysis: Version 6.11 User's Manual," Sandia Technical Report SAND2014-4633, July 2014; updated November 2019.

# Validating Experimental Results of toxins using Cheminformatics and Quantum Chemistry

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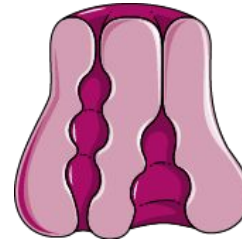
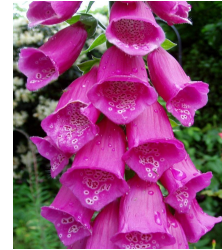
Beauty Kolade  
T-CNLS / B-11  
08/14/2020

Mentor: Jacob Miner (B-11)  
Co-Mentor / PI: Robert F. Williams (B-11)

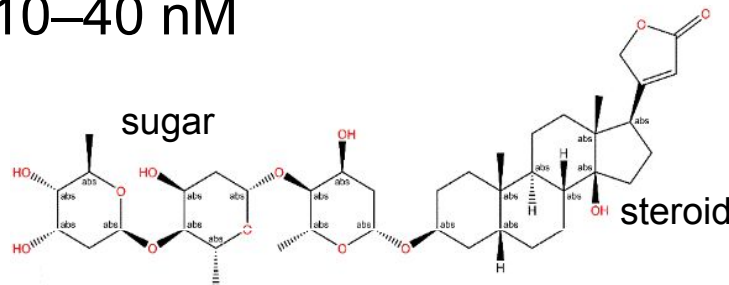


# Plant-based toxins

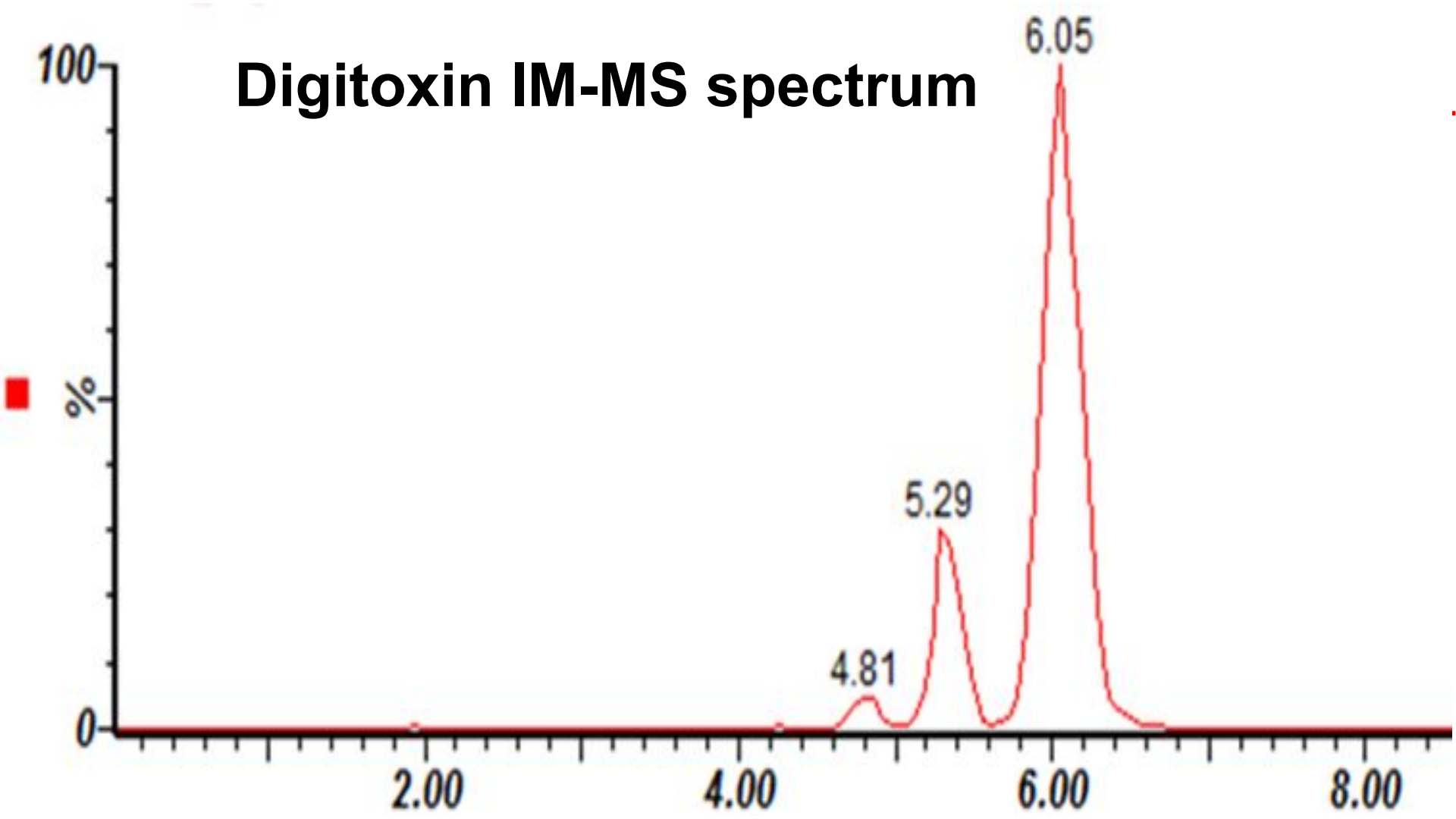
- Derived from natural sources (i.e Foxglove plant)
- Treat heart failure and irregular heartbeats
- Binds to protein receptors
- Therapeutic range=10–40 nM



Structure

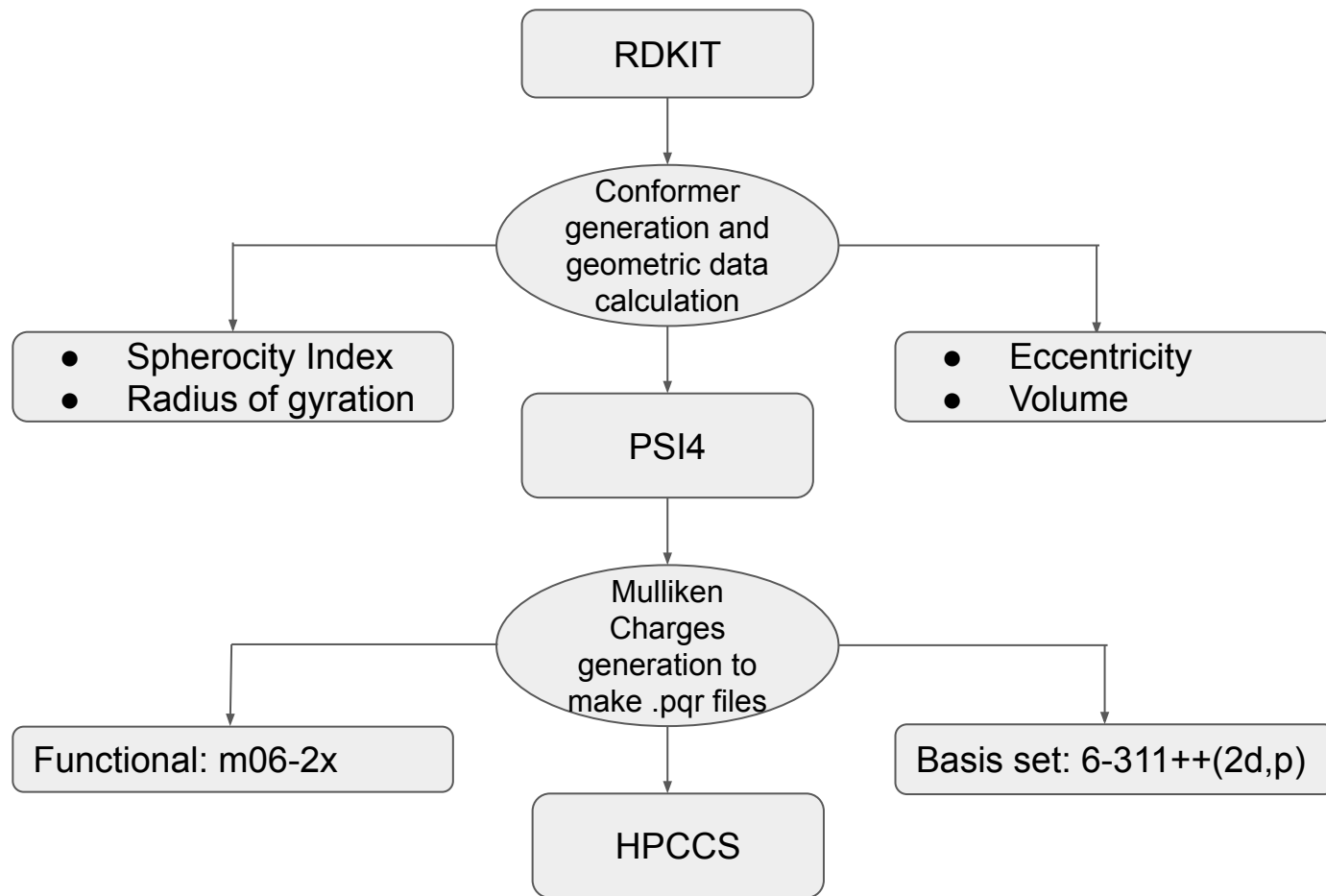


# Digitoxin IM-MS spectrum





# New procedure



# Goals

---

- Develop an approach for rapidly determining toxin molecule configurations to compare with experiment
- Create a library of ion-mobility-based data for cardiac glycosides
- Generate ensembles of molecular configurations without complex parameterization and intensive simulations through molecular dynamics
- Correlate geometric criteria of molecules with calculated CCS values

# Systems

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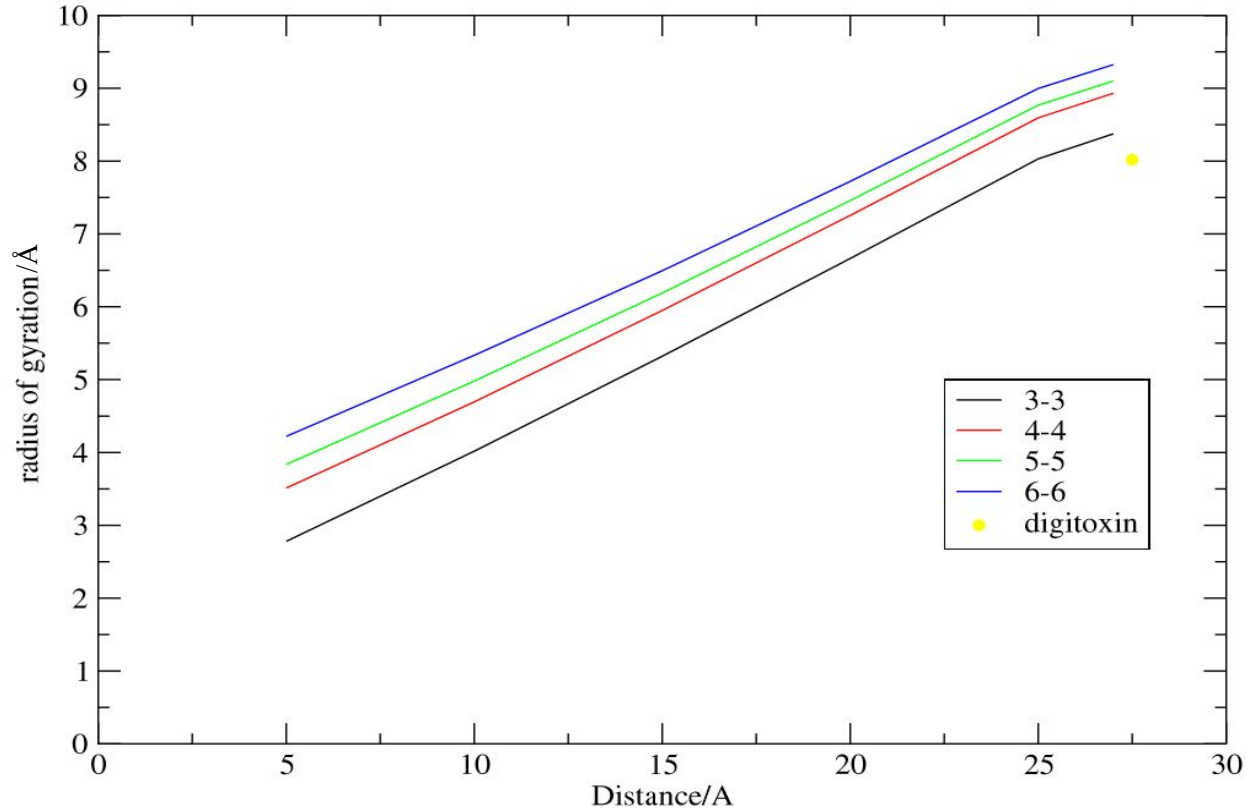
- Digitoxin- Complex structure
- Single-wall carbon nanotubes - oblong shape is similar to Digitoxin shape
- Buckyballs - used to calibrate the original mobility calculators ("like mobcal") and the IM-MS machines
- Single-wall carbon nanotubes and Buckyballs are simple structures
- Their geometric properties can be used as a standard to compare with the geometric properties of Digitoxin

# Geometric data Results-Digitoxin

Molecule	End-to-end distance (Å)	Volume(Å <sup>3</sup> )	Radius of gyration(Å)	Eccentricity	Sphericity Index
Digitoxin(RDKIT)	27.5	721±0.40	8.02±0.48	1.00± 0.009	0.11±0.03
Digitoxin(GROMACS)	23.6	708±3.00	6.85±1.21	0.98±0.02	0.19±0.08
3-3 Nanotubes	27.0	1121	8.38	1.00	0.09

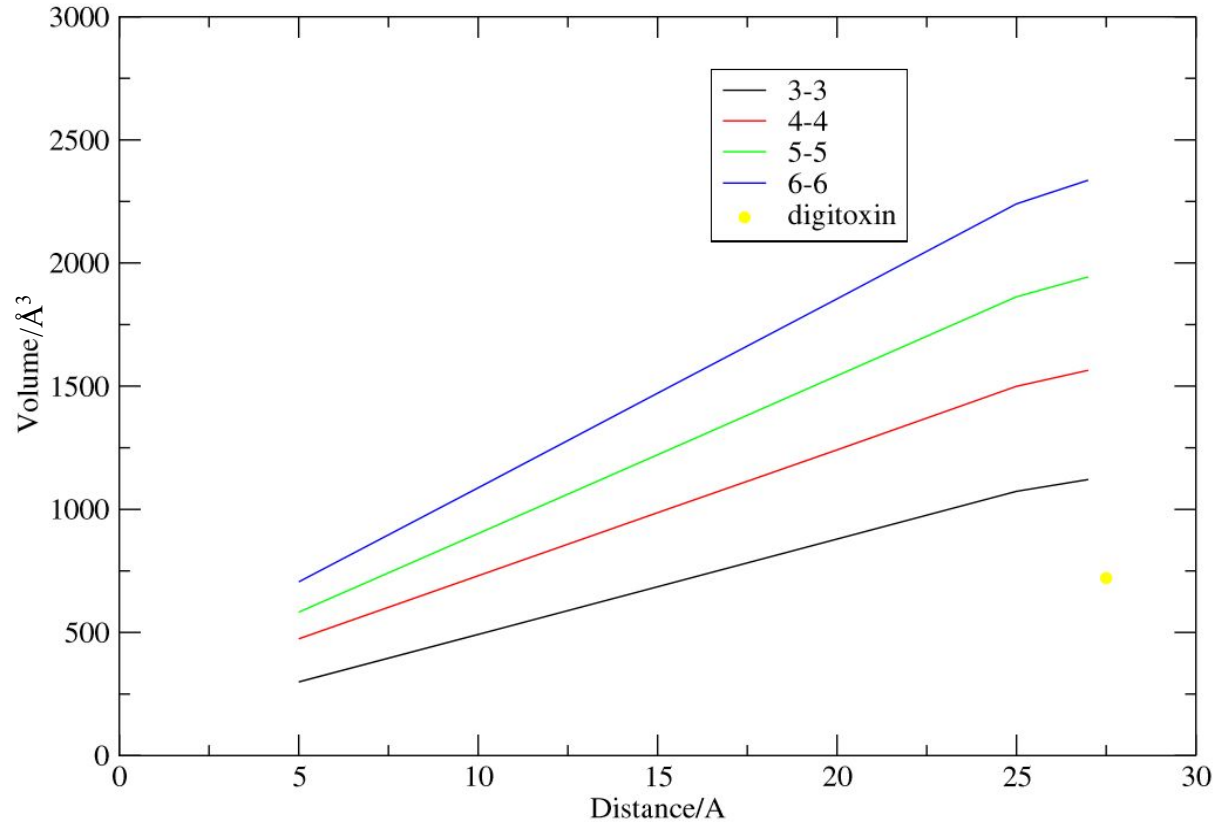
# Geometric data Results -Nanotubes

Rg of the nanotubes



# Geometric data Results -Nanotubes

Volume of the nanotubes



# Next steps

---

- Calculate the CCS values of the most representative Digitoxin Configurations
- Apply clustering methods (i.e. k-means, principal component analysis) to identify the most salient geometric features that correlate with CCS in buckyballs and nanotubes
- ➔ Improve the pre-screening of CCS for biological molecules
- Compare with results of MD simulations (Amber, Ligpargen) to see if RDKit comes closer to predicting CCS.



# Acknowledgement

---

CNLS:Angel Garcia



Mentors: Jacob Miner, Robert Williams

Funding: DHS funding

**Thank you!**

# **Diversification of ML Datasets via Bias Potential as Function of “Uncertainty”**

Maksim Kulichenko

3<sup>rd</sup> year PhD, Utah State University

T-1 Division

Mentor: Justin Smith

Co-mentor: Sergei Tretiak

Collaborator: Kipton Barros

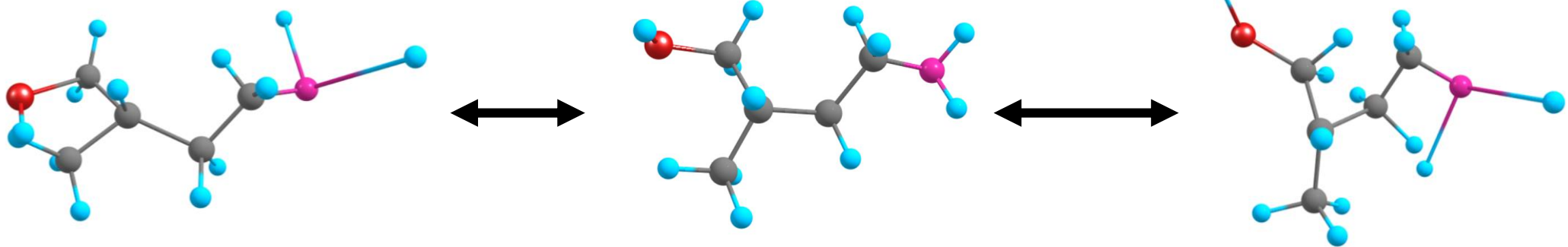
## Make Dataset Better

Purpose: the ideal  
ratio between  
*Time, Size, Diversity*

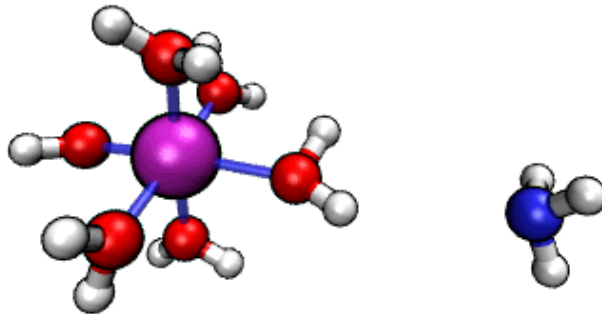
- Capture as much conformational and configurational space as possible
- Dataset diversity increase
- Dataset preparation time decrease

# Different Approaches to Space Sampling

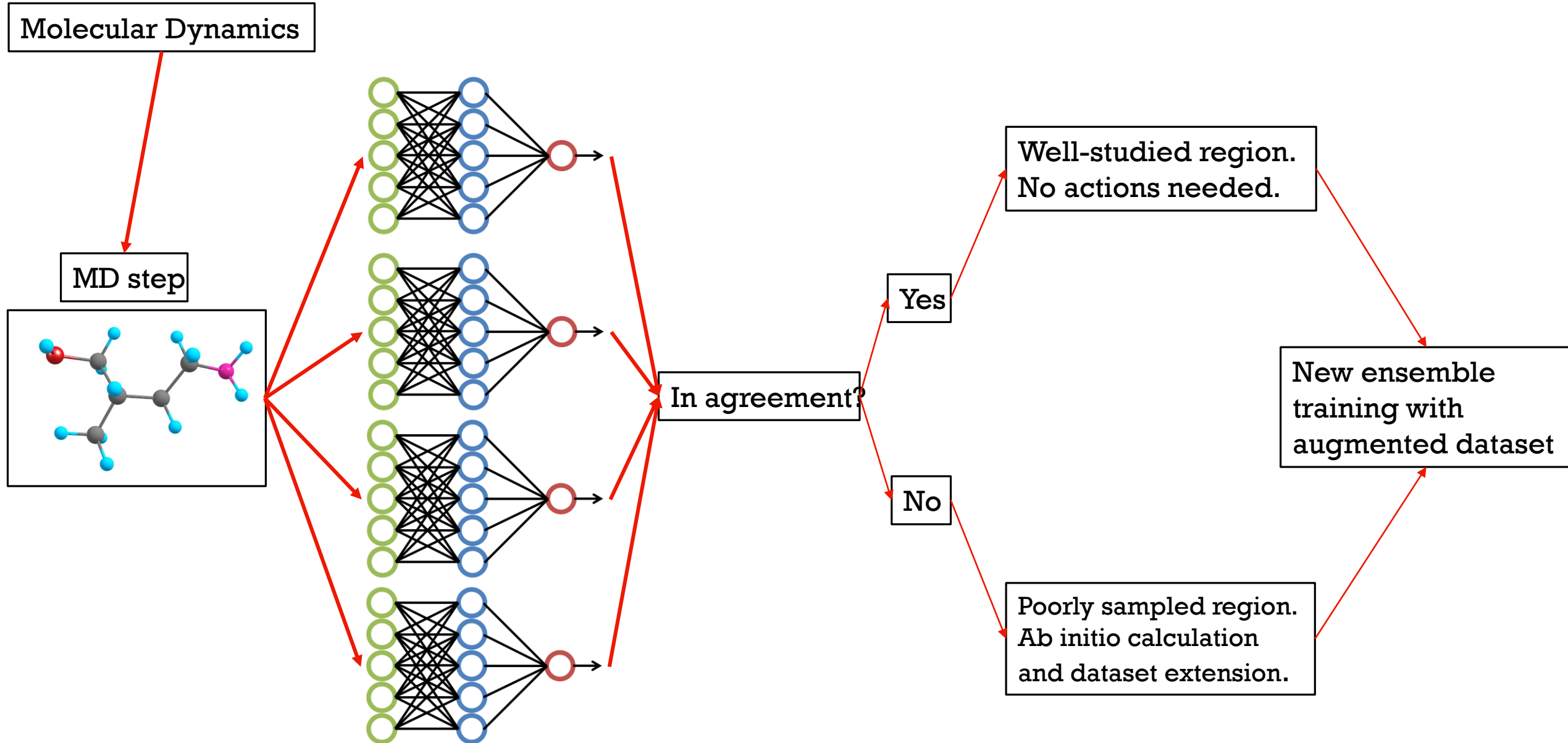
- Random Sampling
- Normal modes descent



- Molecular Dynamics

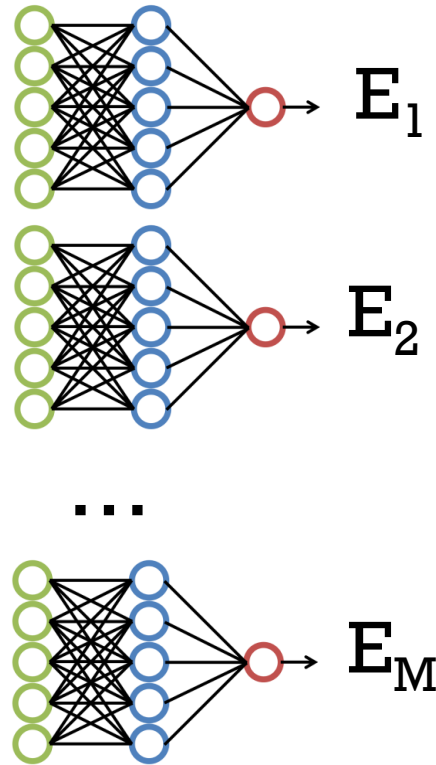


# Active Learning MD Workflow



Moving towards  
Low Sampled  
Region

■ Math of Energy uncertainty



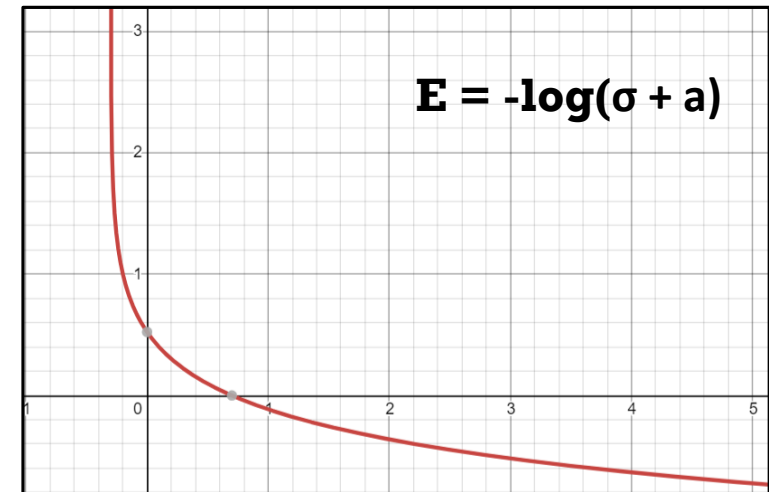
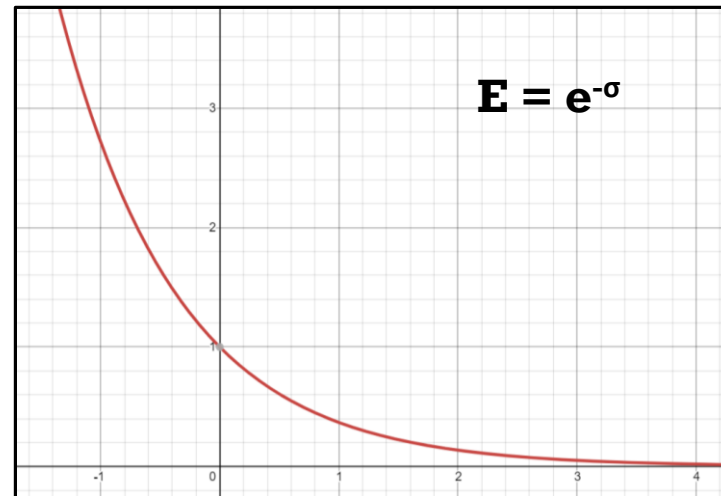
$$\bar{E} = \frac{1}{M} \sum_{m=1}^M \hat{E}_m$$

$$\sigma_E^2 = \frac{1}{M(M-1)} \sum_{m=1}^M \delta E_m^2,$$

$$\delta E_m = \hat{E}_m - \bar{E},$$



# Energy as Function of Uncertainty



## Driving Force of Uncertainty

$$-\nabla_i V(\sigma_{\mathbf{f}}^2) = -V'(\sigma_{\mathbf{f}}^2) \nabla_i \sigma_{\mathbf{f}}^2,$$

Energy uncertainty

$$\sigma_E^2 = \frac{1}{2} \sum_m (\hat{E}_m - \bar{E})^2,$$

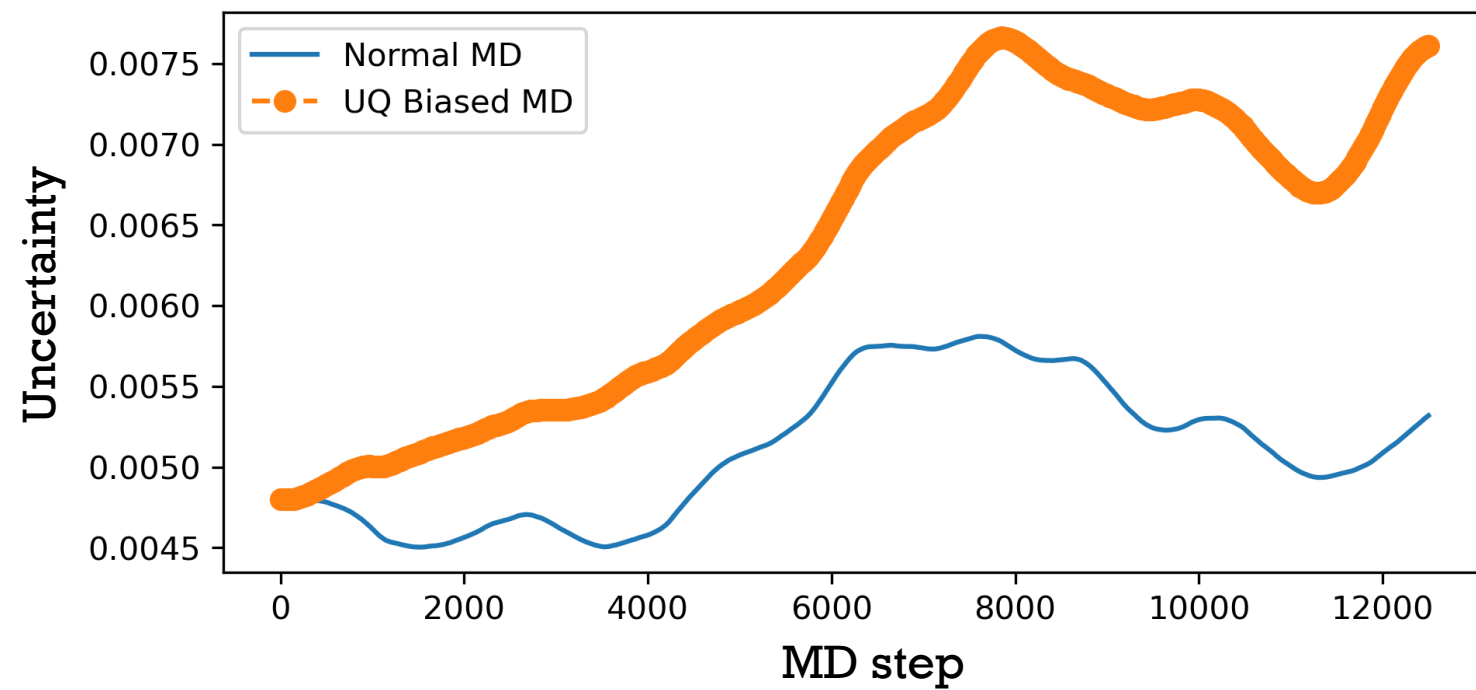
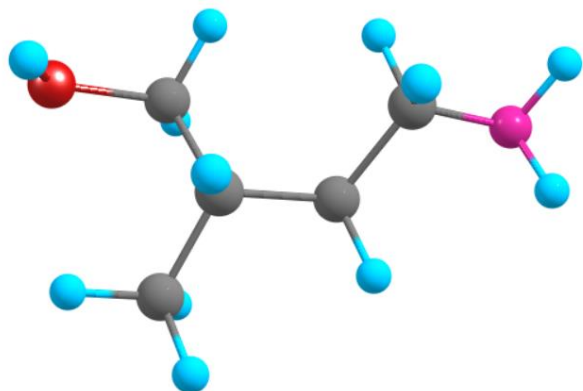
where

$$\bar{E} = \frac{1}{N_{models}} \sum_m \hat{E}_m$$

Bias forces

$$\begin{aligned} -\frac{\partial}{\partial \mathbf{r}_i} \sigma_E^2 &= -\sum_m (\hat{E}_m - \bar{E}) \frac{\partial}{\partial \mathbf{r}_i} (\hat{E}_m - \bar{E}) \\ &= \sum_m (\hat{E}_m - \bar{E}) (\hat{\mathbf{f}}_m - \bar{\mathbf{f}}) \end{aligned}$$

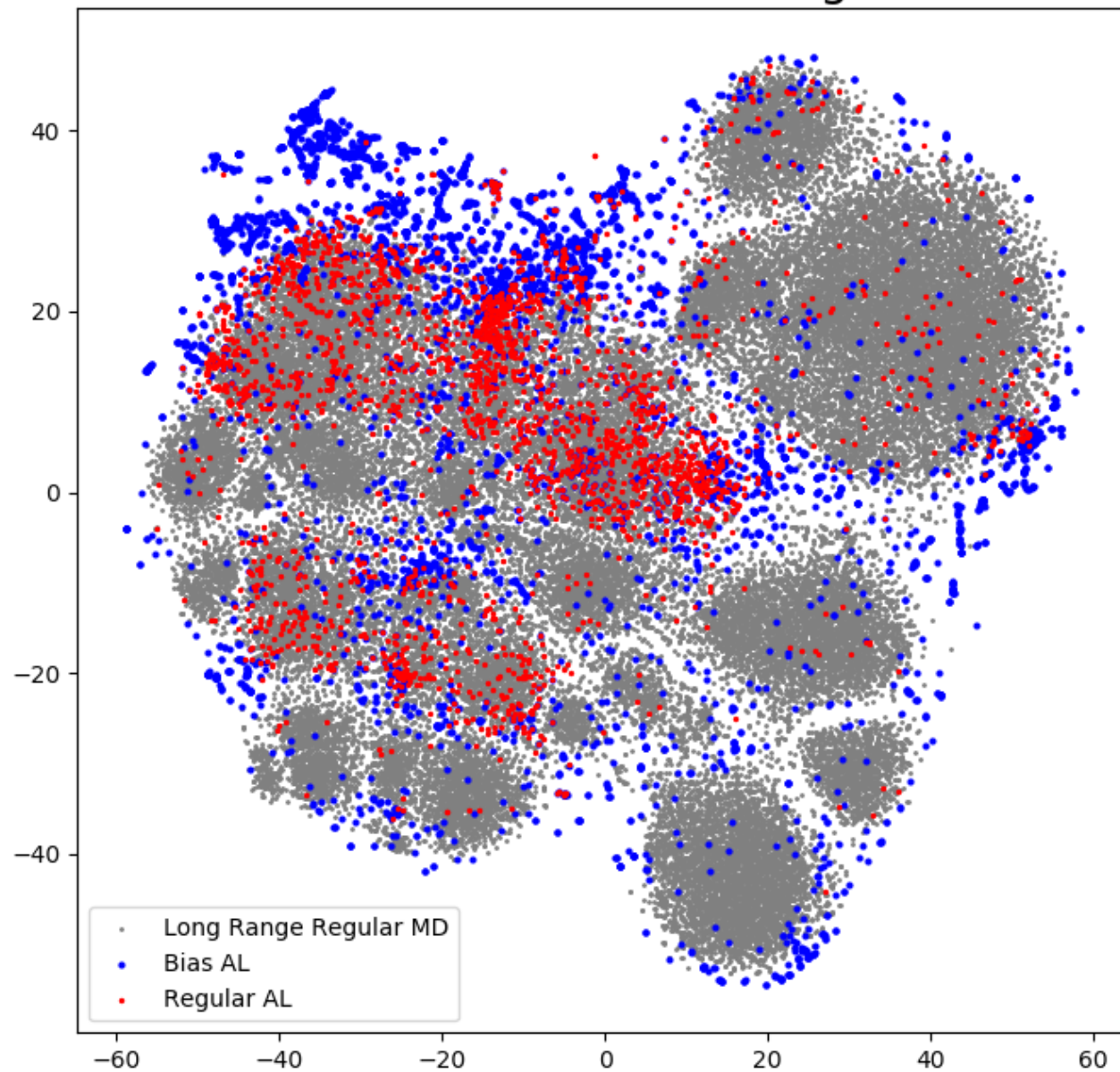
# UQ Test



# t-SNE Latent Space

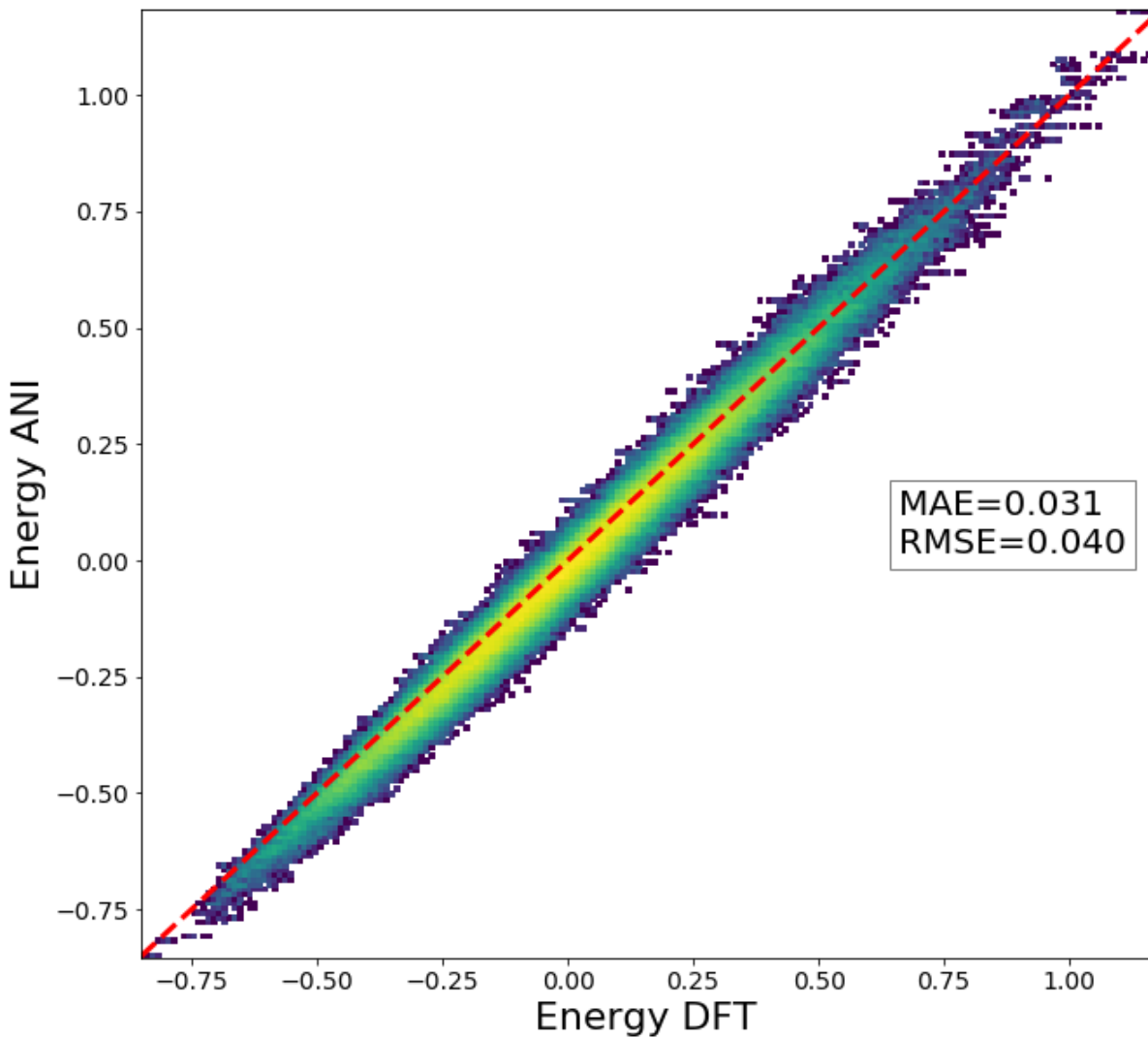
Plot biased and non-biased  
Active learning over long scale  
MD (25 ns)

Molecule Embedding

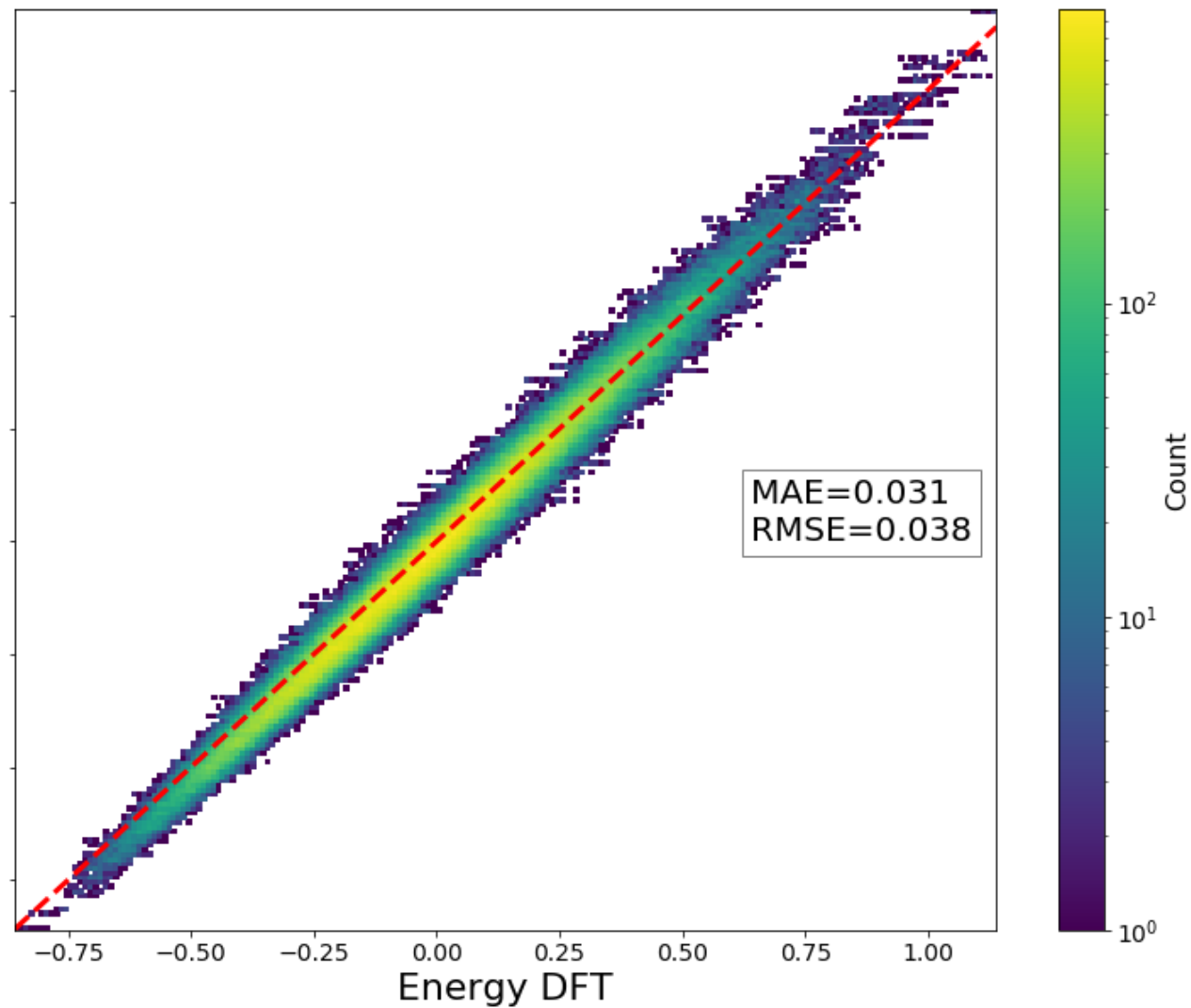


# Feeding Long Range MD data to...

Model trained on non-biased AL data



Model trained on biased AL data



x1.5-2 faster

# Further Development

- Application to larger systems and combinations of different systems
- Implementation of Hessian matrix
- Nano-reactors

# *Bayesian Uncertainty Quantification for COVID-19 Situational Awareness*

Abhishek Mallela

Department of Mathematics  
University of California, Davis

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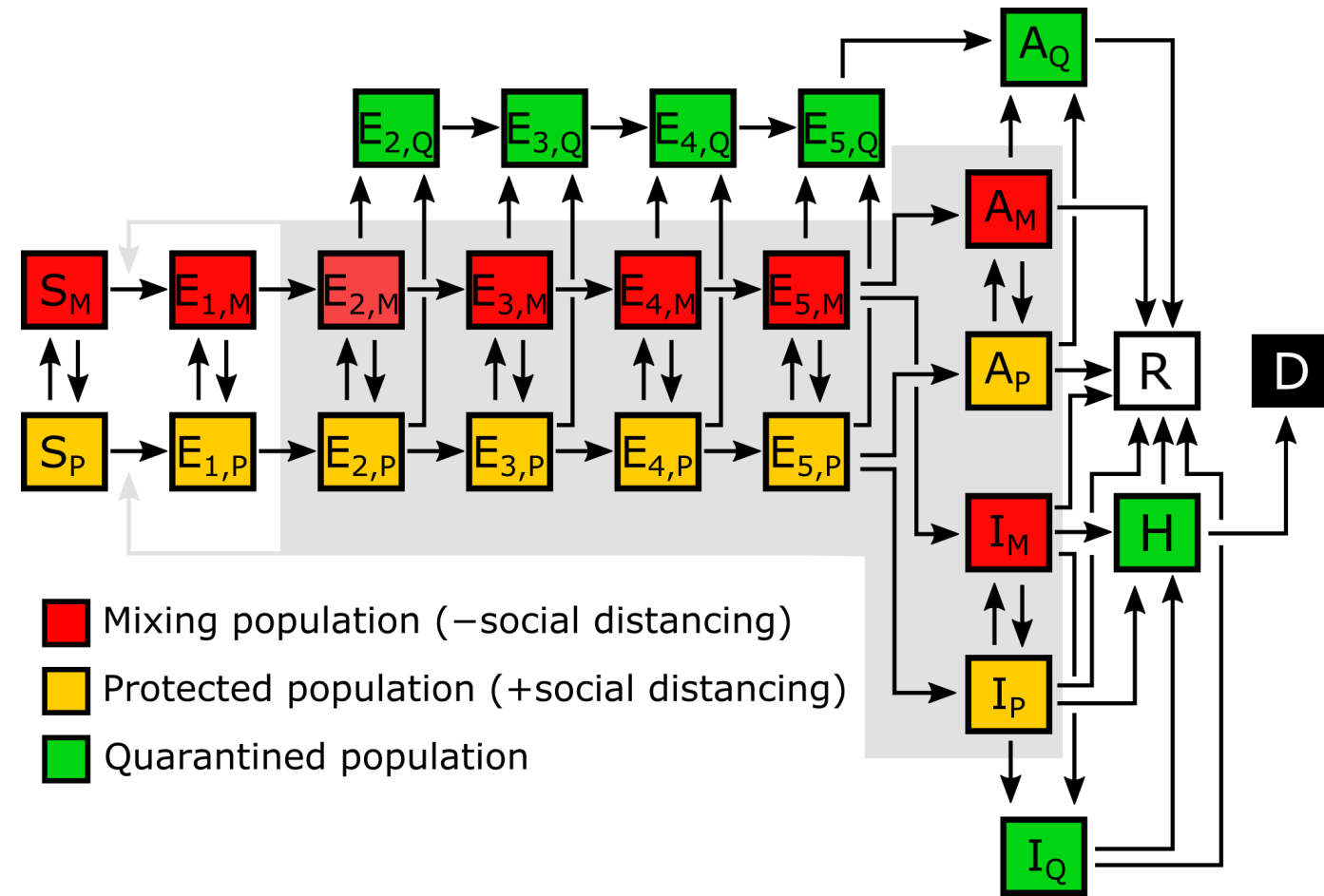
Mentor: William S. Hlavacek

Group: T-6 (Theoretical Biology and Biophysics)



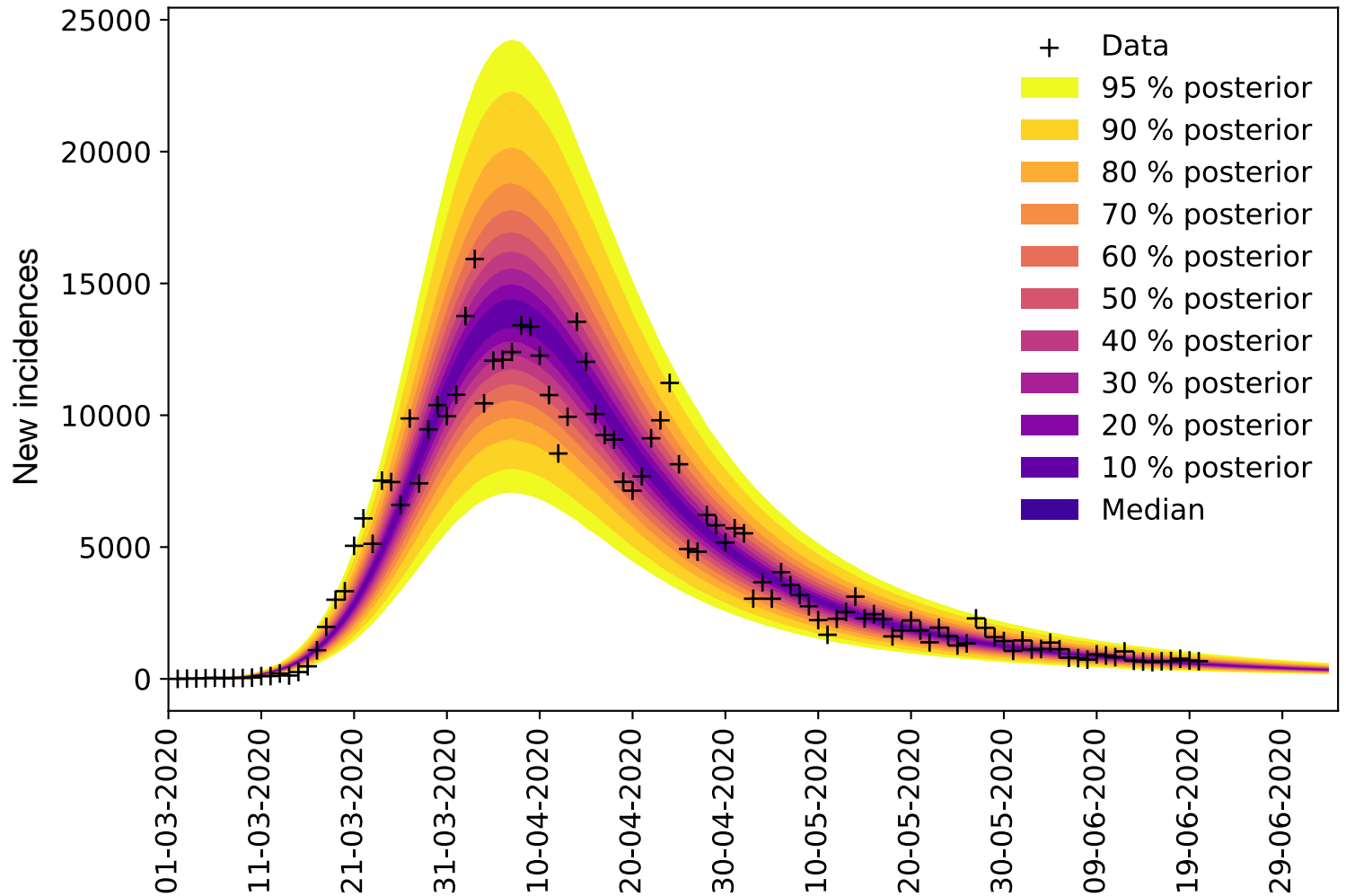
## Illustration of the populations and processes considered in a mechanistic compartmental model for the dynamics of COVID-19 transmission.

The model accounts for susceptible (S), exposed (E), asymptomatic (A), symptomatic (I), hospitalized (H), recovered (R), and deceased (D) populations. It also accounts for social distancing, which establishes mixing and protected subpopulations, quarantine driven by testing and contact tracing, and self-isolation spurred by symptom awareness. The incubation period is divided into 5 stages, which allows the model to reproduce an empirically determined Erlang distribution of waiting times for the onset of symptoms after infection. The exposed population (consisting of individuals incubating the virus) includes presymptomatic and asymptomatic individuals. The *A*-populations consist of true asymptomatic individuals in the immune clearance phase. The gray background indicates the populations that contribute to disease transmission. An auxiliary measurement model accounts for imperfect detection and reporting of new cases.

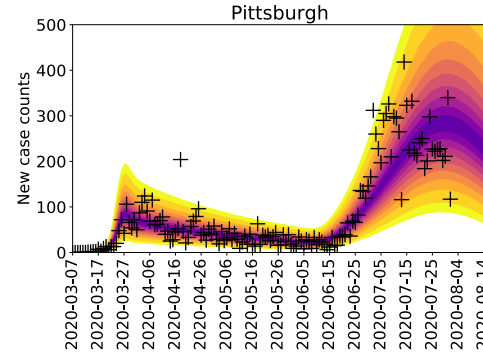
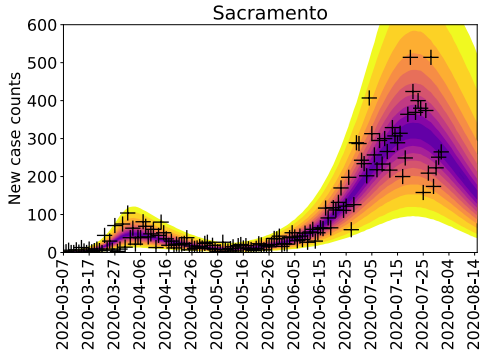
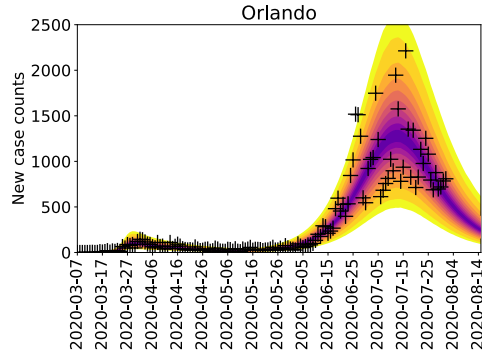
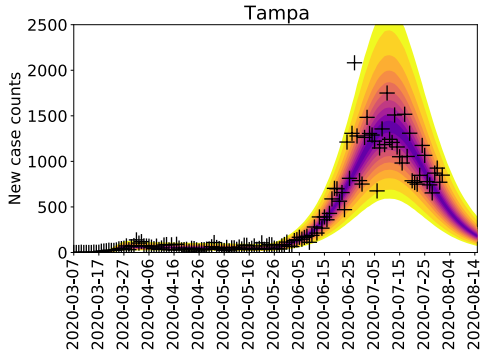
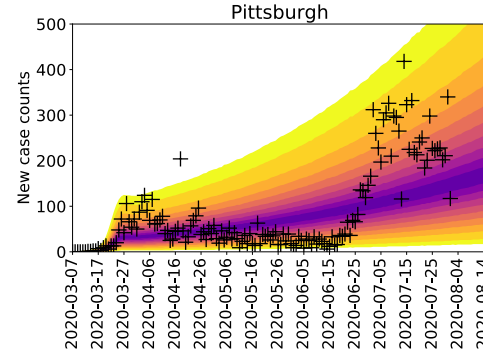
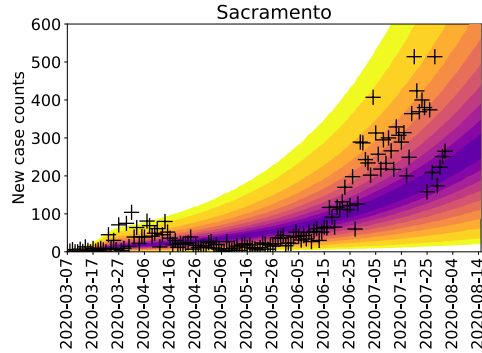
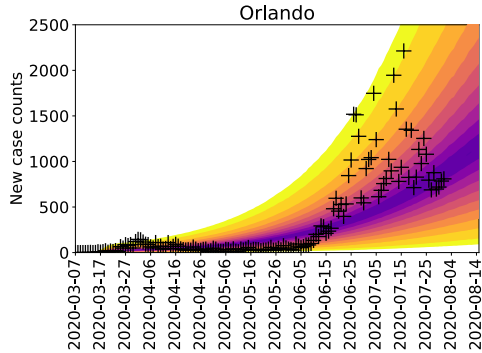
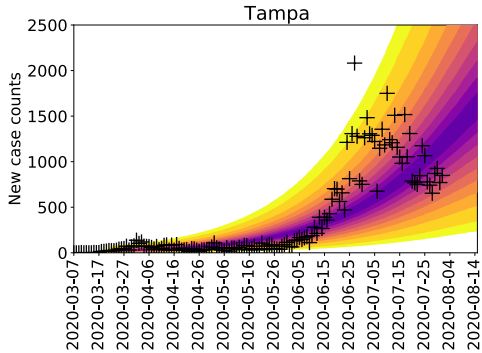


## Illustration of Bayesian predictive inference.

We forecast future daily reports of new COVID-19 cases with rigorous uncertainty quantification (UQ) through online Bayesian learning of model parameters. Each day, using all daily case-reporting data available up to that point, we perform Markov chain Monte Carlo (MCMC) sampling of the posterior distribution for a set of adjustable parameters. Subsampling of the posterior samples then allows us to use the relevant model to generate trajectories of the epidemic curve that account for both parametric and observation uncertainty. The entire shaded region indicates the 95% credible interval for predictions of daily case reports. In other words, the central 95% of all predictions lie within the shaded region. The color-coded bands within the shaded region indicate other credible intervals, as indicated in the legend.



Top row: One phase; Bottom row: Two phases



# Summary

- In this work, we report the computational/statistical/mathematical modeling methodology we are using to make daily forecasts of new-case detection with rigorous quantification of uncertainty for each of the regional epidemics in the 15 most populous metropolitan areas in the United States.
- Our results
  - justify a modeling focus on cities
  - illustrate early detection of upward trends in new-case detection
  - explain rising-falling-rising epidemic curves because of distinct periods of adherence to social-distancing practices that reduce disease transmission
- Building on this work, I have shown that local COVID-19 epidemics are well-explained by a compartmental model with time-varying social-distancing parameters. The model can predict second-wave dynamics for any of the 384 metropolitan statistical areas in the United States.

# References

Lin, Yen Ting, Jacob Neumann, Ely Miller, Richard G. Posner, **Abhishek Mallela**, Cosmin Safta, Jaideep Ray, Gautam Thakur, Supriya Chinthavali, and William S. Hlavacek. "Daily Forecasting of New Cases for Regional Epidemics of Coronavirus Disease 2019 with Bayesian Uncertainty Quantification." In submission, *aRxiv preprint: arXiv:2007.12523* (2020).



USC University of  
Southern California

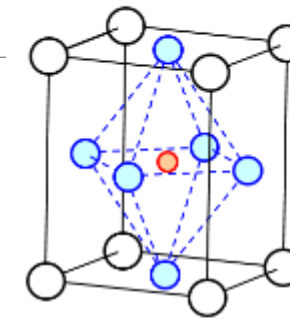
# TDDFT Characterization of $\text{CsPbBr}_3$ Clusters Optical Properties

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CARLOS MORA PEREZ, AMANDA NEUKIRCH, DIBYAJYOTI GHOSH,  
SERGEI TRETIK

# Motivation

## Metal halide perovskite

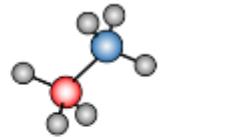


○ A (organic cation,  $\text{Cs}^+$ )

● B ( $\text{Pb}^{2+}$ ,  $\text{Sn}^{2+}$ , ...)

● X ( $\text{I}^-$ ,  $\text{Cl}^-$ ,  $\text{Br}^-$ )

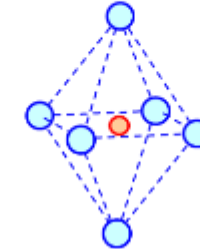
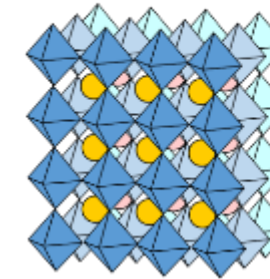
$\text{CH}_3\text{NH}_3\text{PbI}_3, \dots$



methylammonium  
( $\text{CH}_3\text{NH}_3^+$ )



formamidinium  
( $\text{NH}_2\text{CHNH}_2^+$ )



$\text{BX}_6$  octahedron

Perovskites most abundant mineral on earth!!

Halide perovskites ( $\text{ABX}_3$ ) for optoelectronic devices

❖ Hybrid

❖ All-inorganic

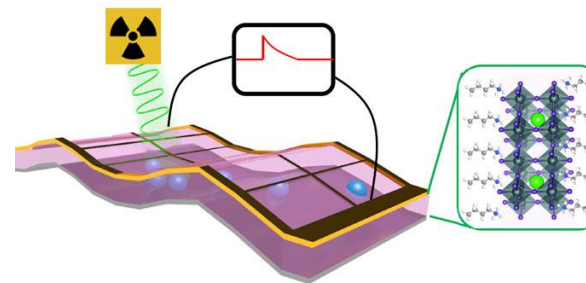
Optoelectronic applications

❖ Electroluminescence (EL)

❖ Photoluminescence (PL)

❖ X-ray

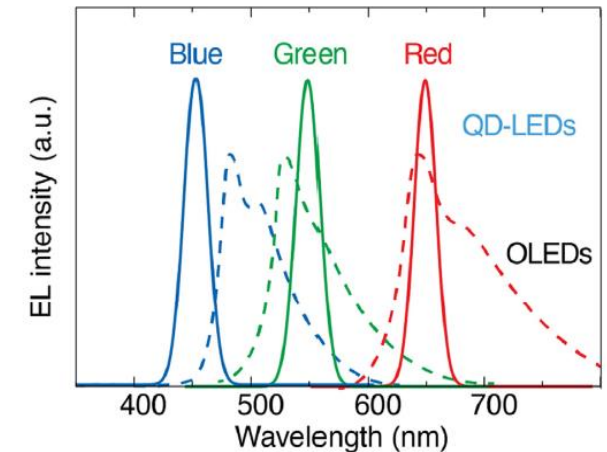
❖ LEDs



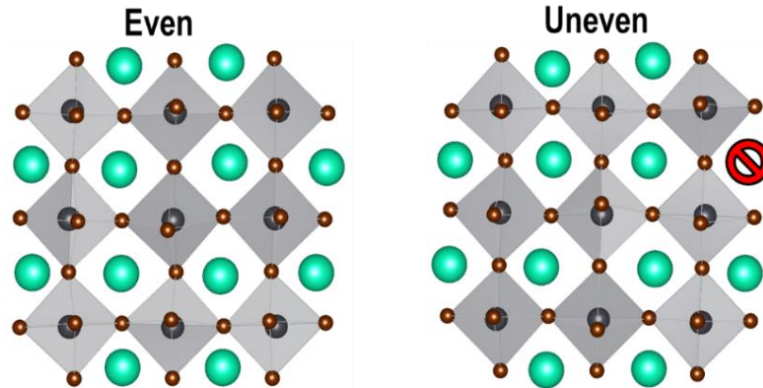
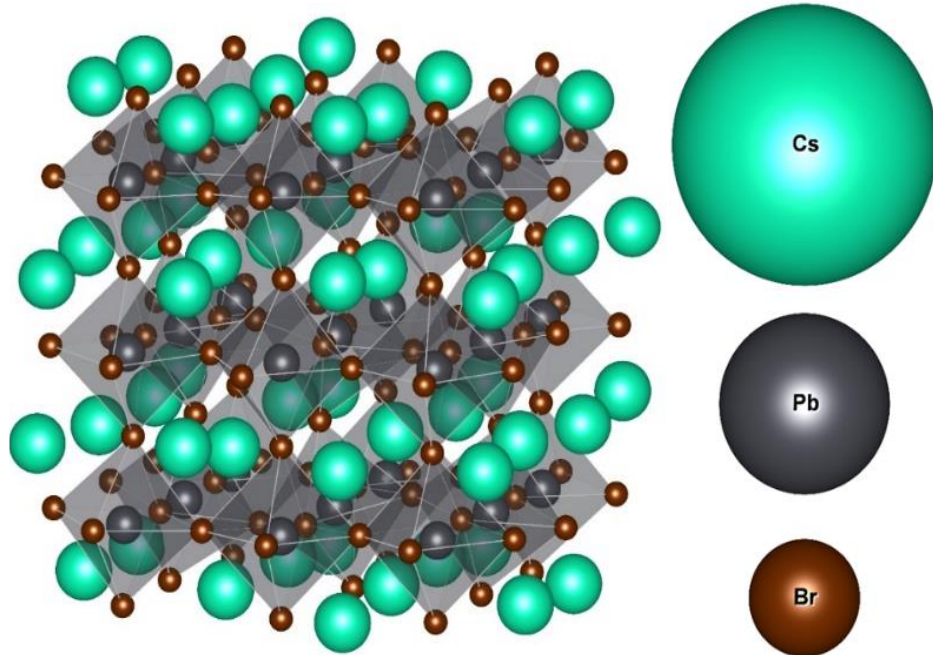


# Why perovskites?

1. Direct gap semiconductor
  1. Easily tunable
  2. Cover Whole visible light range
  3. High PL quantum yield
2. High defect tolerance
3. Cheaper raw materials and equipment
4. Ambipolar Charge Transfer
  1. Can transport both holes and electrons
  2. Better than OLEDs
5. Facile Synthesis
  1. Solution-based methods at low temperature.
  2. VS. all-vacuum thermal evaporation method in OLEDs



# CsPbBr<sub>3</sub> Cluster



## Pristine System

❖ 189 Atoms

○ Cs<sub>54</sub>Pb<sub>27</sub>Br<sub>108</sub>

❖ Dimension

○ 1.8x1.8 nm

❖ Unique Surfaces

○ Even

○ Uneven

Point Defects (low formation energy)

❖ Br<sup>-1</sup>

❖ Pb<sup>+2</sup>

❖ Cs<sup>+1</sup>

❖ Positive: [CsPbBr<sub>3</sub>]<sup>+1</sup>

❖ Negative: [CsPbBr<sub>3</sub>]<sup>-1</sup>

## Computational Methods

❖ Software: Gaussian16

○ Optimization

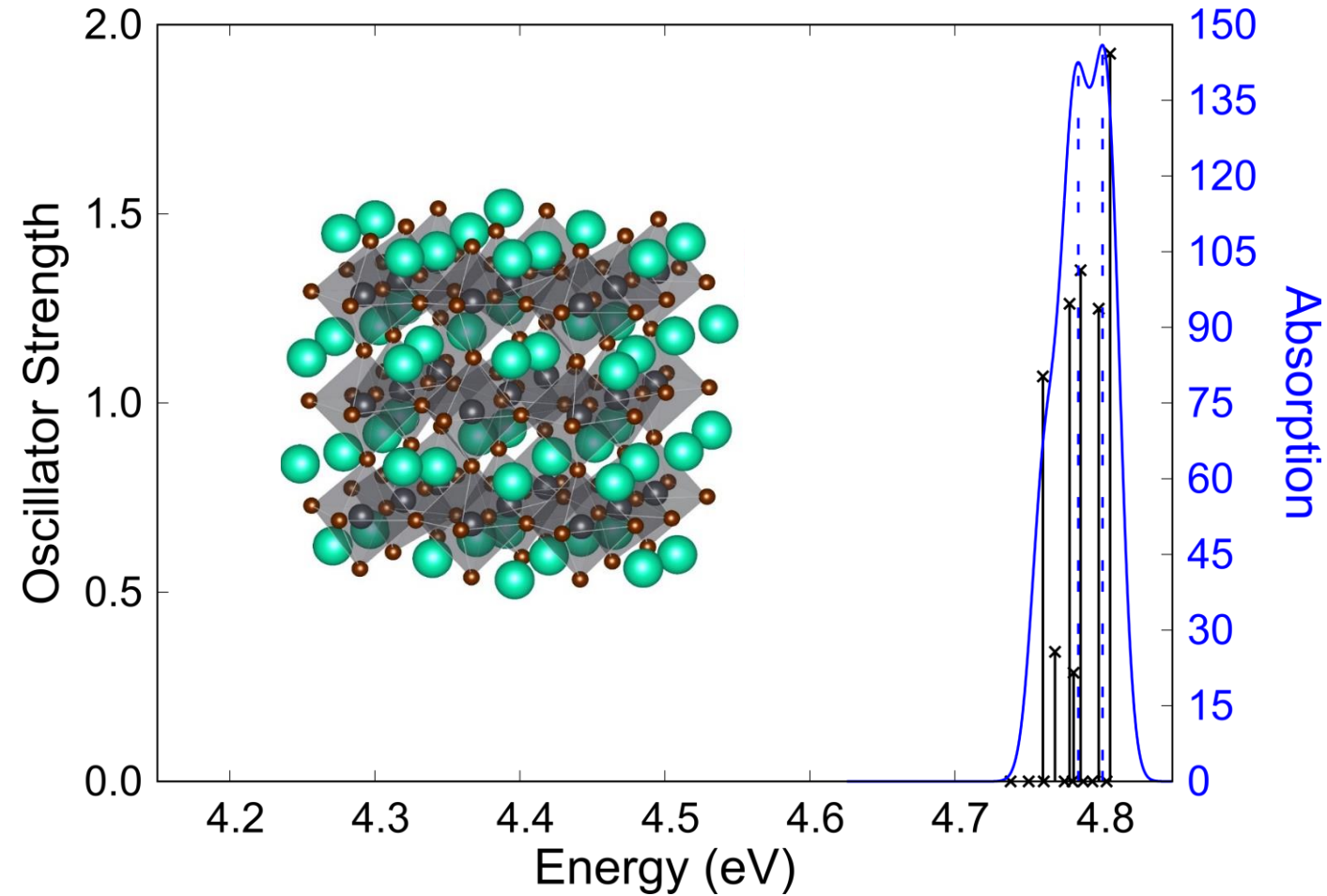
○ TDDFT Calculations

❖ Functional: CAM-B3LYP

❖ Basis Set: LANL2DZ

❖ Solvent: Water

# UV-Vis Profile



High energy due to small size of cluster

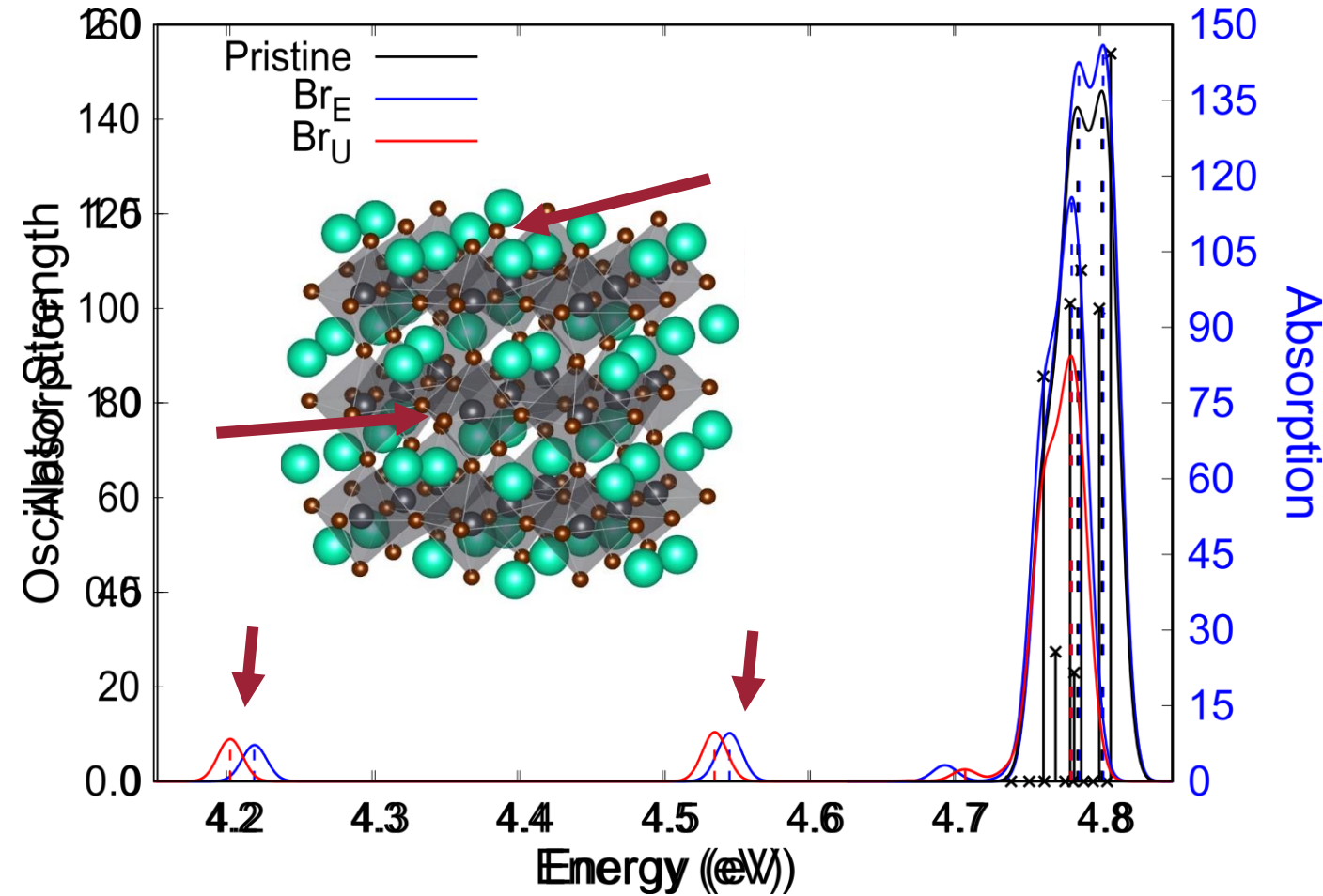
Band peaks at 4.78 and 4.80 eV

Vacancies introduces defects states which lowers energy gap

✦ ~4.2 eV band peak

✦ ~0.5 eV red shift

# UV-Vis Profile



High energy due to small size of cluster

Band peaks at 4.78 and 4.80 eV

Vacancies introduces defects states which lowers energy gap

✦ ~4.2 eV band peak

✦ ~0.5 eV red shift

# Summary

---

Vacancies introduce defect states which lowers bandgap

- ❖ Leads to lower absorption
- ❖ Lower performance devices

Mitigate defects to gain performance

- ❖ Halide rich/Cs poor environment
- ❖ Surface passivation
  - Ligands
  - Oxygen

The background is a solid blue color. Overlaid on it are several faint, light-blue circular patterns. On the left side, there is a large circular scale with tick marks and numbers ranging from 140 to 260. Other smaller circular patterns with arrows indicating direction are scattered across the page.

# Physics based machine learning report card

CONJUGATED HYDROCARBONS

Yamilée Morency  
Mentor: Justin Smith

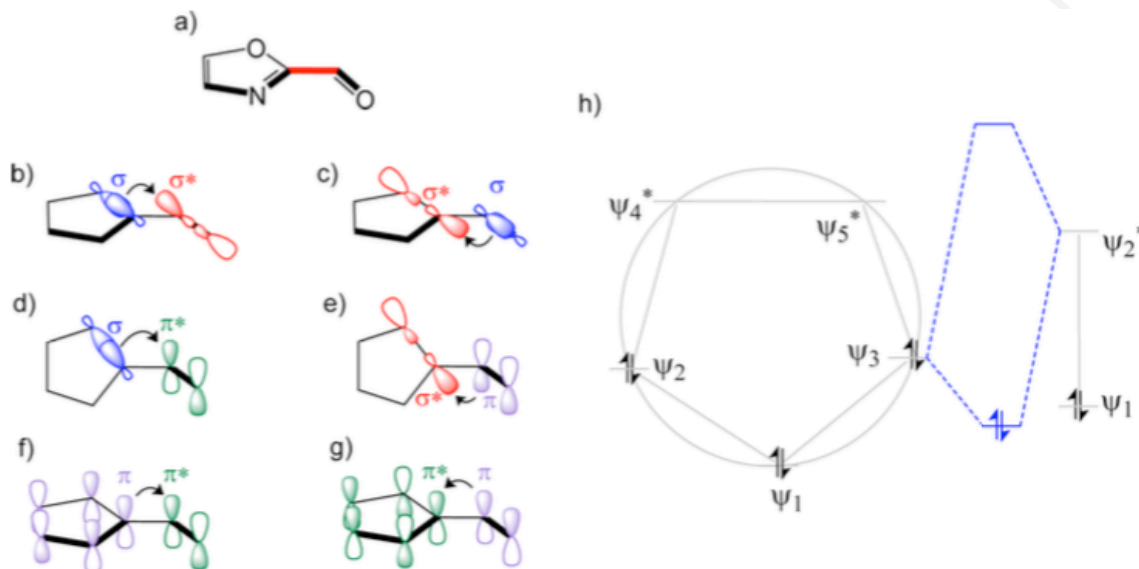


# BACKGROUND

- Molecular potential energy surfaces are essential to solving problems in physics, chemistry and biology
- Theoretical ab initio methods are not tractable for very large systems
- Machine Learning Potentials (MLPs) provide a faster alternative with little to no loss in accuracy
- MLPs depend on chemical feature representation, quality of the data & ML algorithms
- MLPs have no physical basis
- How well do these models reproduce complex physical phenomena?



# Conjugated Hydrocarbons



Wei, W., Champion, C., Barigye, S. J., Liu, Z., Labute, P., & Moitessier, N. (2020). Use of Extended-Hückel Descriptors for Rapid and Accurate Predictions of Conjugated Torsional Energy Barriers. *Journal of Chemical Information and Modeling*, 60(7), 3534-3545. doi:10.1021/acs.jcim.0c00440

# Models & Training data

## Data

- All hydrocarbons from GDB13 & GDB17
- All PubChem hydrocarbons w/ up to 20 carbons
- Total 122k hydrocarbons
- Coordinates for molecules in training generated from rdkit & MMFF94 optimized

## Model 1 (no Active Learning)

- ANI-1 type model
- Trained with 3.5 k conjugated molecules selected at random from the 122k

## Model 2 (Active Learning):

- Active learning sampled from 122k molecules

# ANI Functional Form

$$E_T = \sum_{i=1}^{\text{Natom}} E_i$$

Heavily modified Behler Parrinello symmetry functions (BPSFs)

BPSFs compute atomic environment vector (AEV)

Each element in an AEV

$$G_m^{\text{Amod}} = 2^{1-\xi} \sum_{i \neq j, k}^{\text{Natom}} (1 + \cos(\theta_{ijk} - \theta_s))^\xi e^{-v(R_{ij} + R_{ik} + R_s)} f_c(R_{ij}) * f(R_{ik})$$

The local AE is achieved with this piecewise cutoff function:

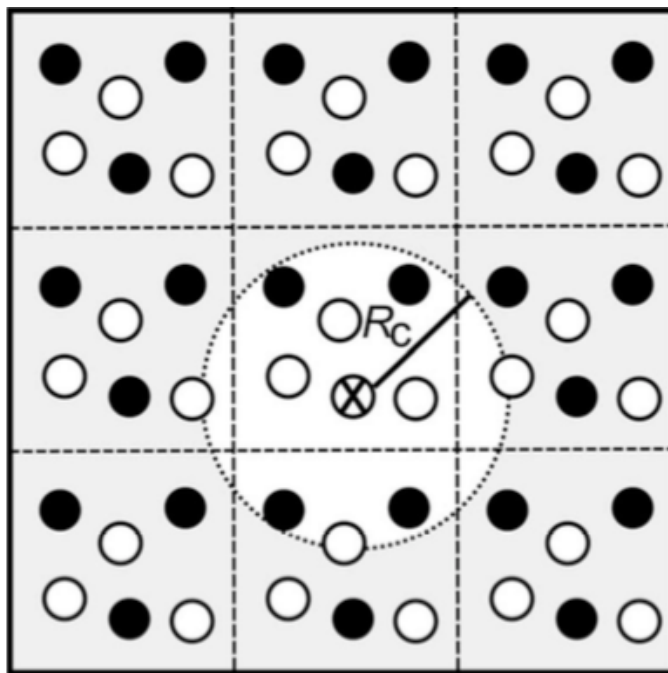
$$f_c(R) = \begin{cases} 0.5 \cos\left(\frac{\pi R_{ij}}{R_c}\right) + 0.5 & \text{for } R_{ij} \leq R_c \\ 0 & \text{for } R_{ij} > R_c \end{cases}$$

The AE cuts off at :

4.6 Å for radial symmetry functions

3.1 Å for angular symmetry functions

# Conjugated Hydrocarbons



Behler, J. (2015). Constructing high-dimensional neural network potentials: A tutorial review. *International Journal of Quantum Chemistry*, 115(16), 1032-1050. doi:10.1002/qua.24890

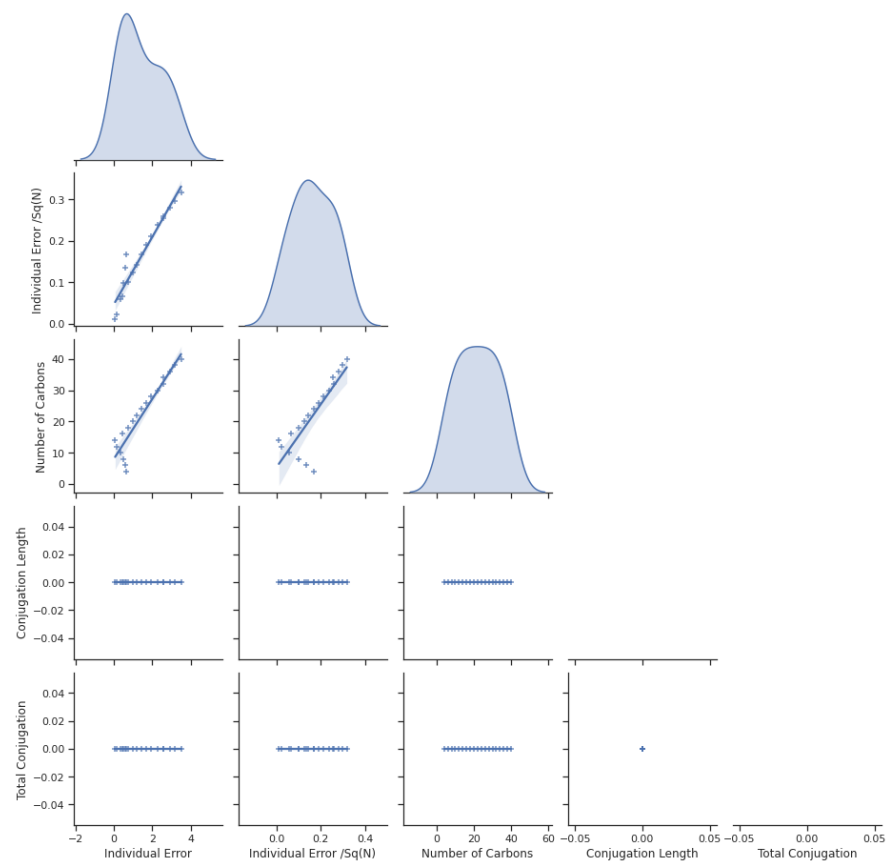
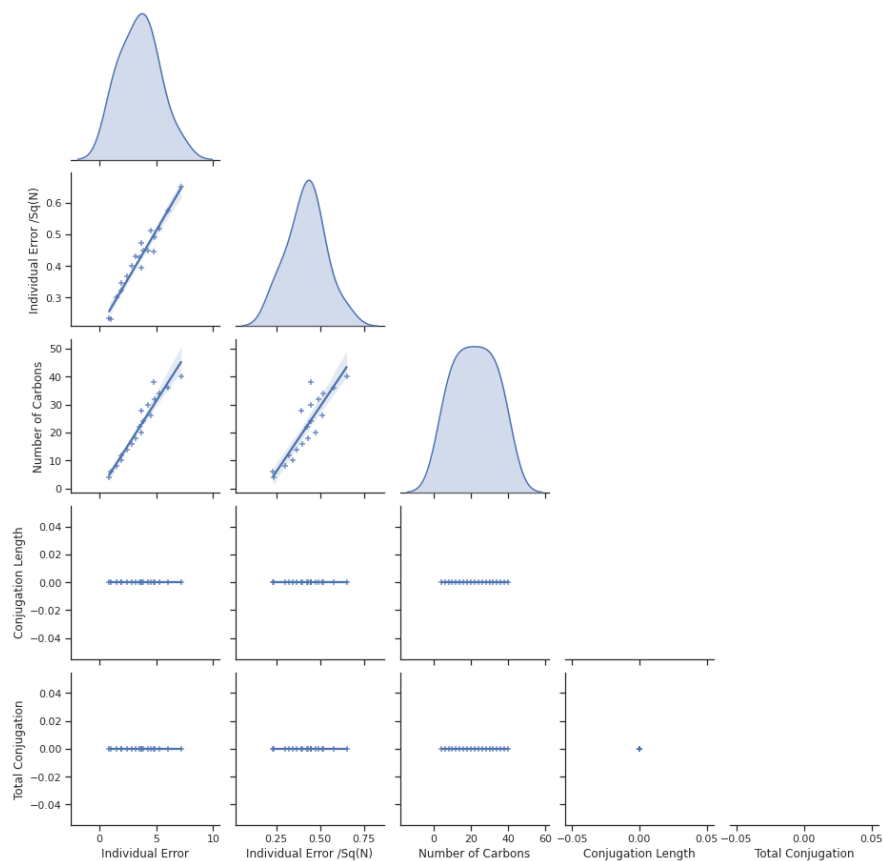
# Conjugated Hydrocarbons

- Alkanes & Alkenes
  - Generated with rdkit from SMILES strings
  - 2 to 32 carbons
  - MMFF94 Optimized
  - Torsion angles frozen at 180°
- Alkynes
  - Generated with rdkit from SMILES strings
  - 2 to 32 carbons
  - MMFF94 Optimized
- Annulenes
  - Generated with rdkit from SMILES strings
  - Up to 20 carbons
  - MMFF94 Optimized
- Conjugation Length
  - Linear Hydrocarbons
  - Increasing conjugation
  - All molecules have 32 Carbons
  - Generated with rdkit from SMILES strings
  - MMFF94 Optimized
  - All molecules have 32 Carbons

# Plot information

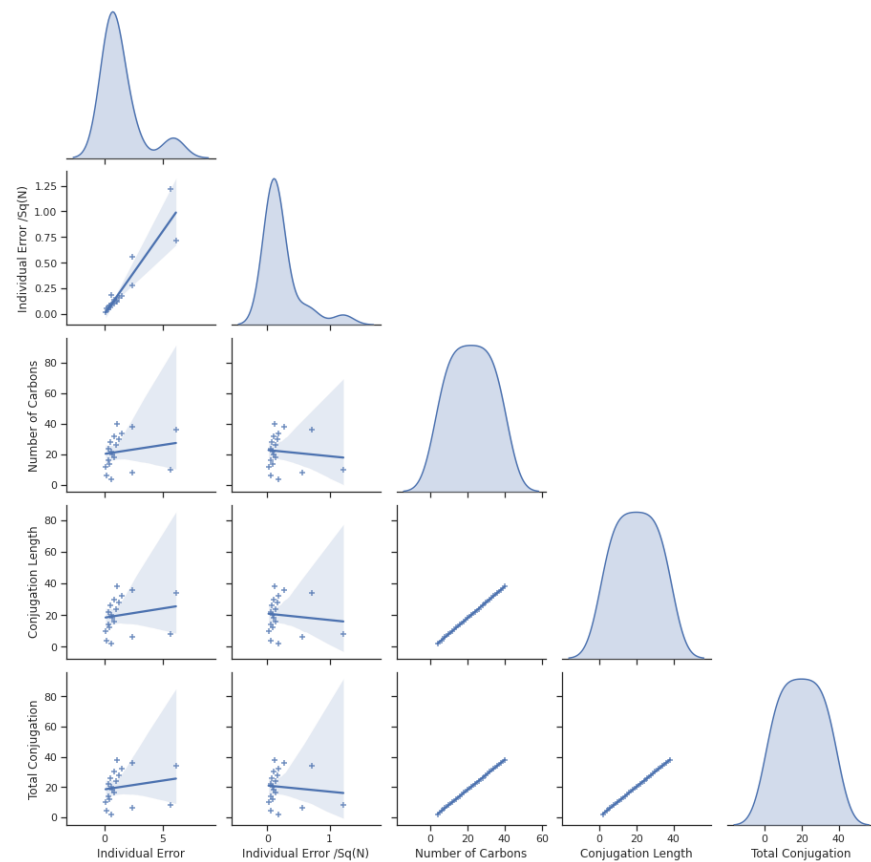
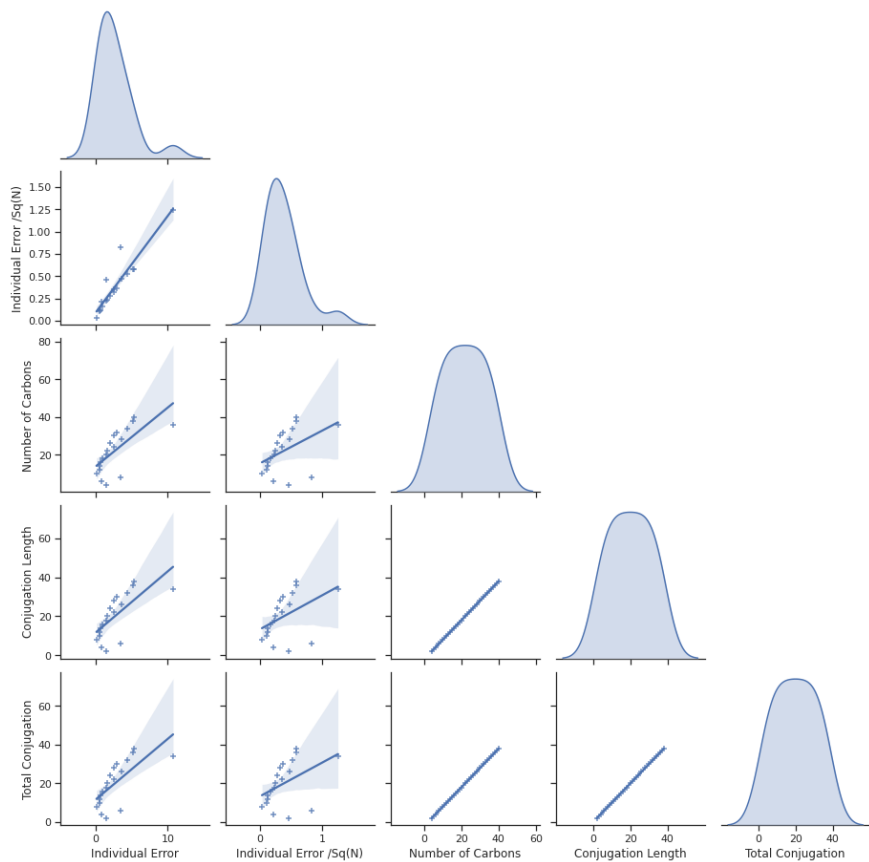
- Individual Error: difference between predicted and calculated energy
- Individual Error /  $\text{Sq}(N)$ : Individual error divided by the square root of the number of atoms in molecule
- Conjugation Length: Longest length of uninterrupted pi bonds (2 or more)
- Total Conjugation: Sum of all the length of uninterrupted pi bonds

# Alkanes

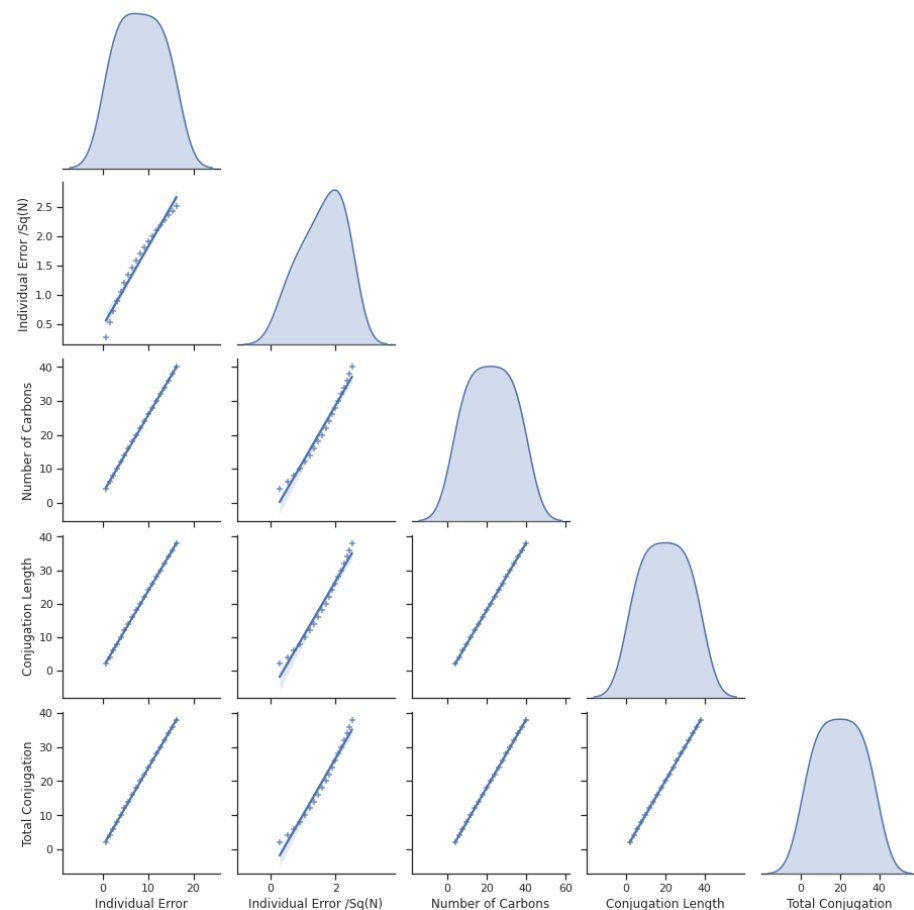
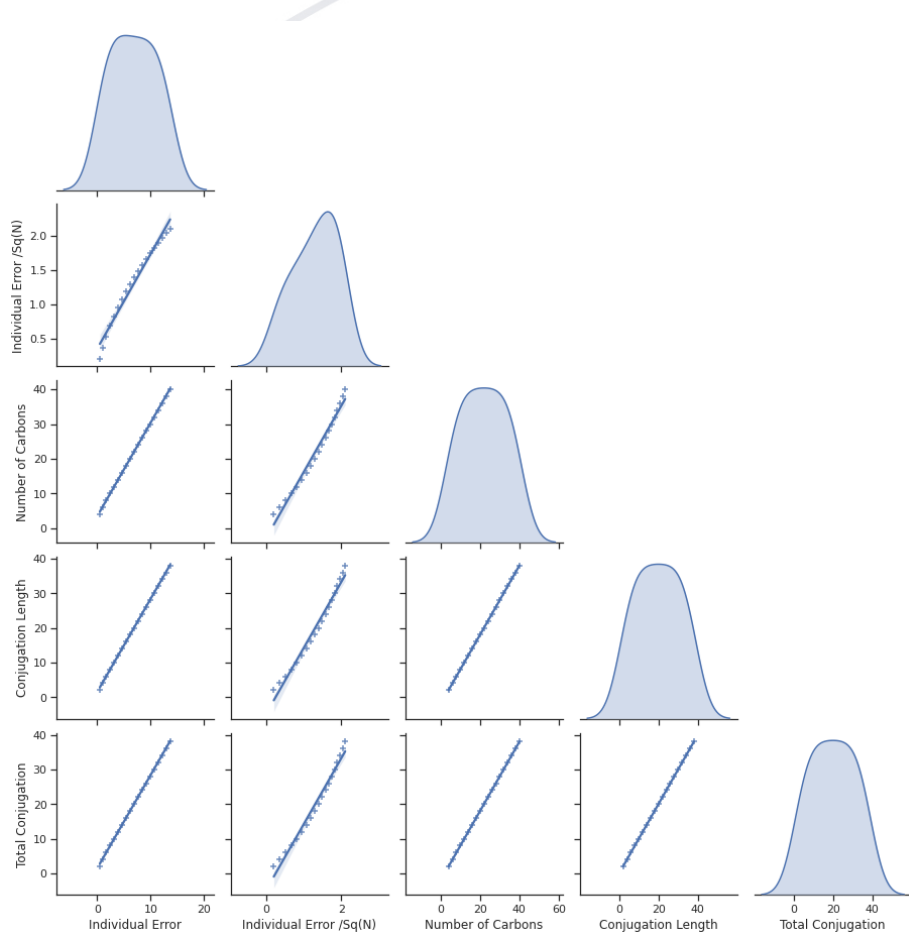




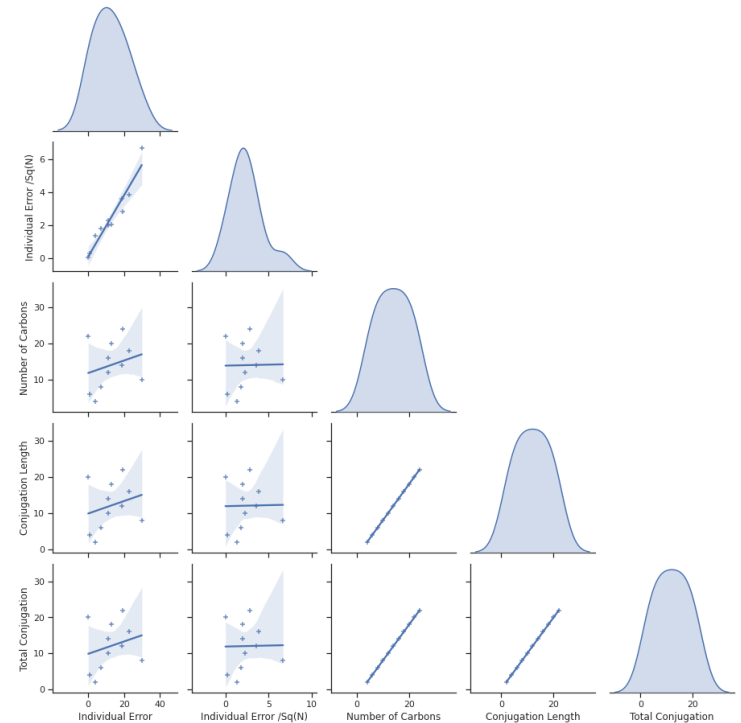
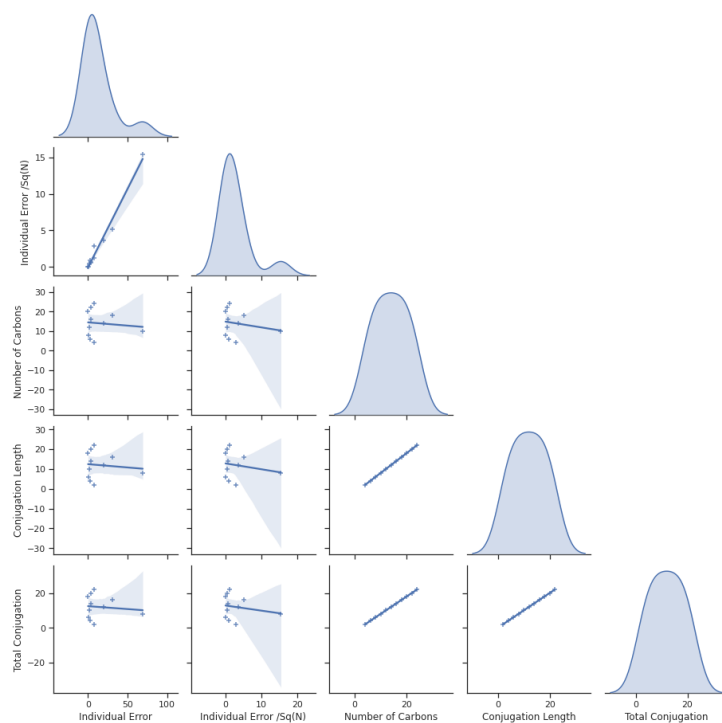
# Alkenes



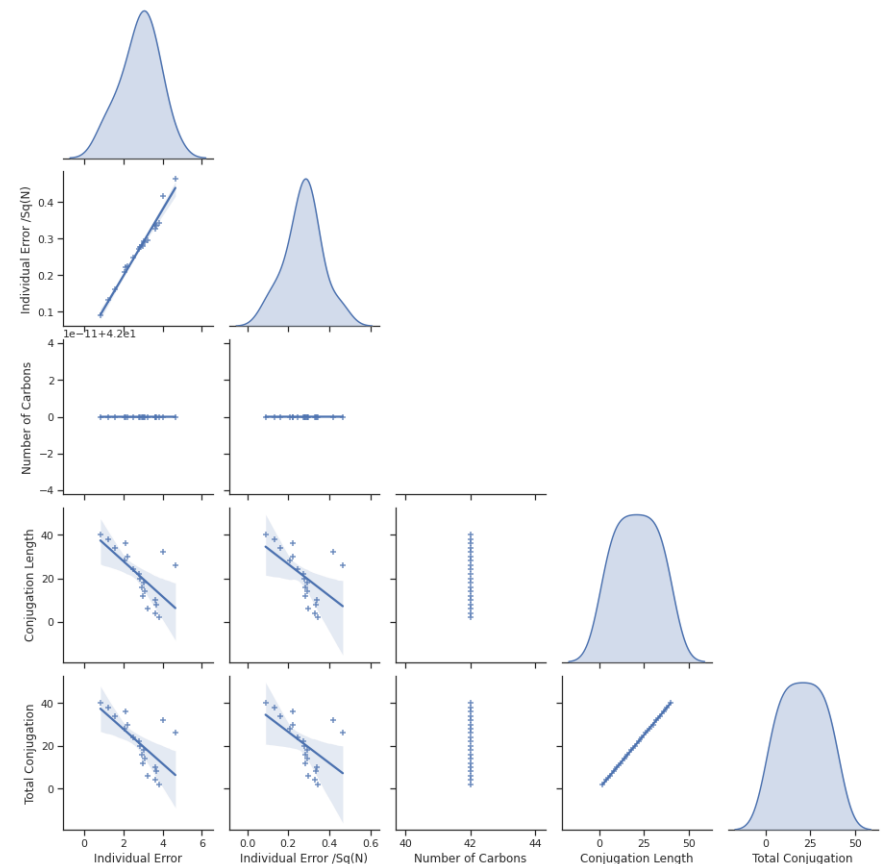
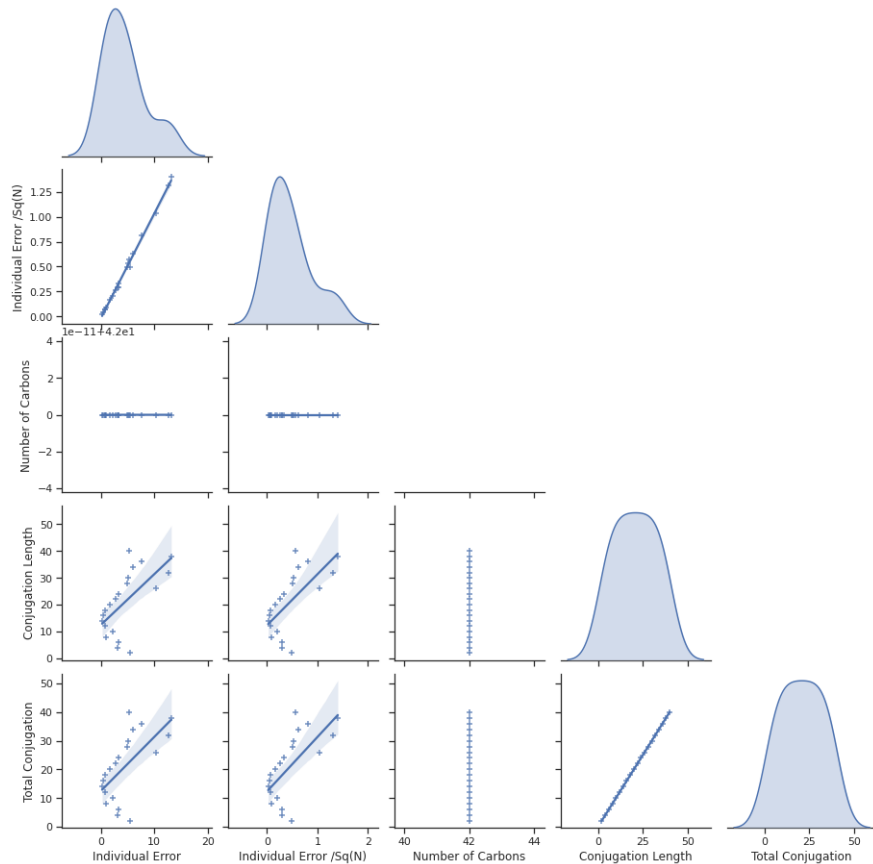
# Alkynes



# Annulenes

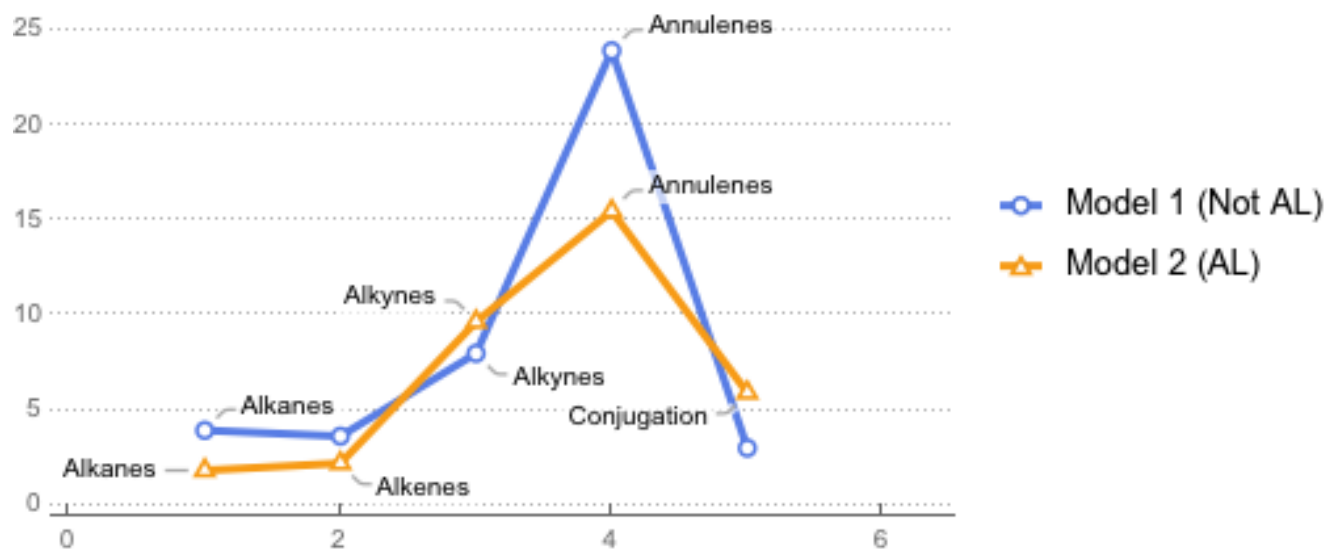


# Conjugation Length



# Model 1 vs Model 2

RMSE Model 1 vs Model 2



# Moving forward

- Add more tests to conjugated hydrocarbons
- Increase training set to improve model
- Active learning to improve model
- Qualitative test



## Lightning Talks

# Nonparametric estimation in a generalized SEIR model

Amanda Patrick, MS

Mentors: Nick Hengartner, PhD and Imelda Trejo, PhD

August 4, 2020

**Theoretical Biology and Biophysics**

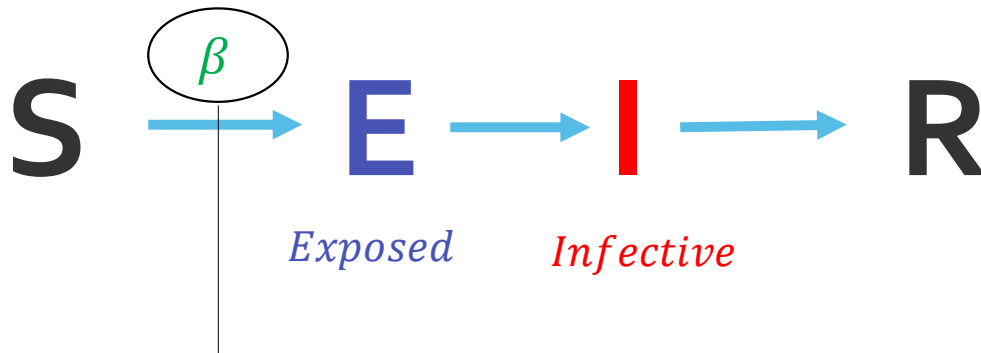
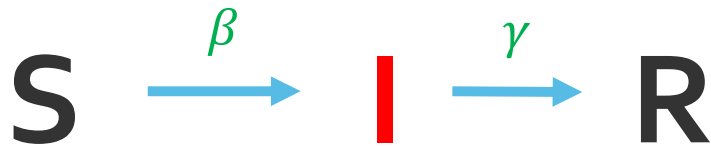




# Compartmental Modeling for Infectious Diseases

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*transmission rate*   *recovery rate*



*What if this is not constant?*

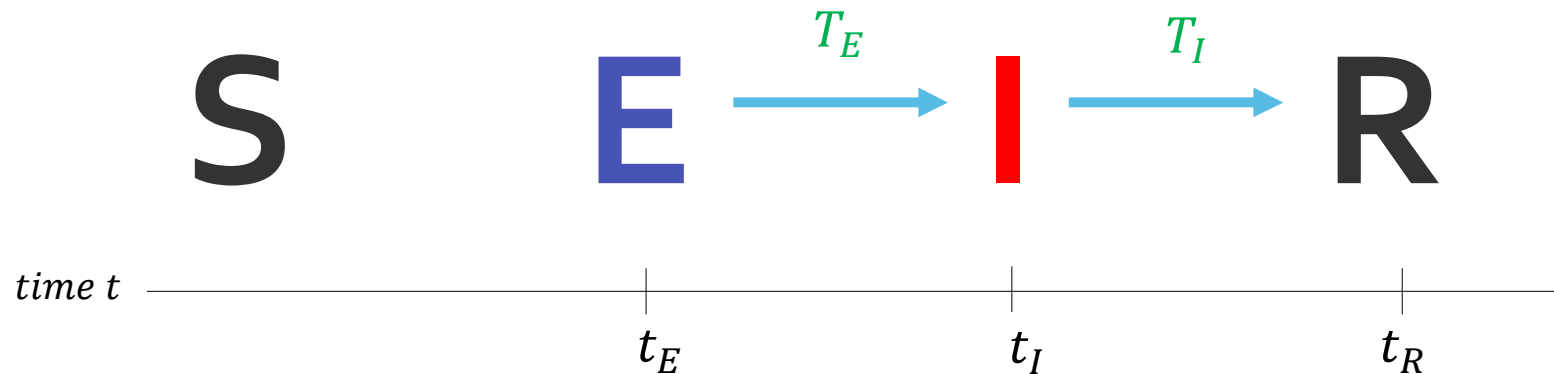
*Force of infection: Per person rate at which a susceptible person will contract the disease*

$$\frac{dS}{dt} = -\beta IS$$

$$\text{Force of infection} = \beta I$$

# Explanation

Random Variables:



Incoming Infected per time  $-S'(t) = \rho(t)I(t)$

$\rho(t)$ : per capita rate at which an infective will transmit the disease

# Generalized SEIR Model

Set of integral-differential equations

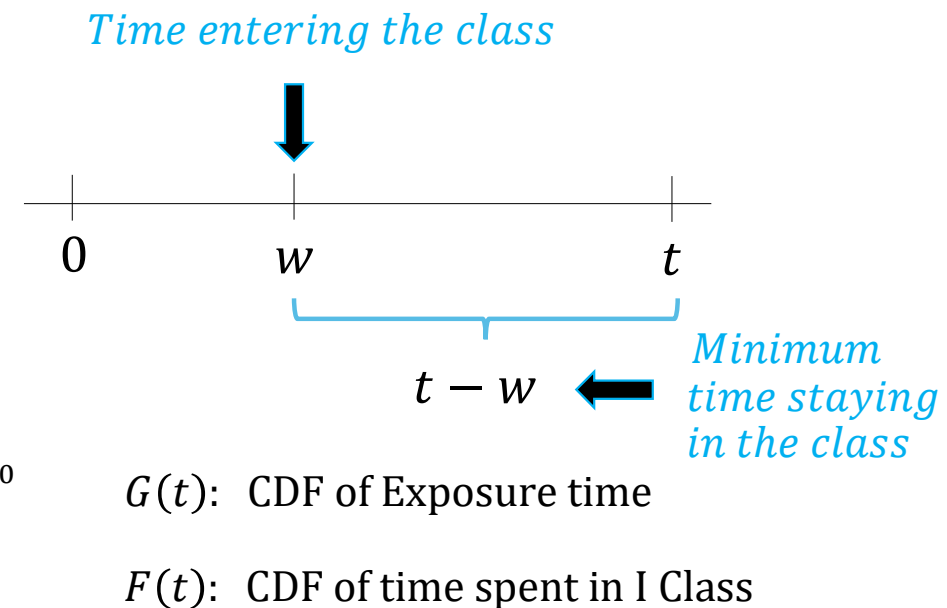
$$S'(t) = -\rho(t)I(t)$$

$$E(t) = \int_0^t -S'(w)(1 - G(t - w))dw + (1 - G(t))E_0$$

$$I(t) = \int_0^t v(w)(1 - F(t - w))dw + (1 - F(t))I_0 + G(t)E_0$$

$$R(t) = \int_0^t v(w) F(t - w)dw + R_0 + F(t)I_0$$

$$v(t) = \int_0^t -S'(w)g(t - w)dw: \text{ Incoming Infective per time}$$



# Two Equations for the Infective Class

*Rate coming into the I class*

$$I(t) = \int_0^t \overbrace{\mathbf{v}(\mathbf{w})}^{P(T_I > t - w)} (1 - F(t - w)) dw + (1 - F(t))I_0 + G(t)E_0 \quad \leftarrow \text{From Model}$$

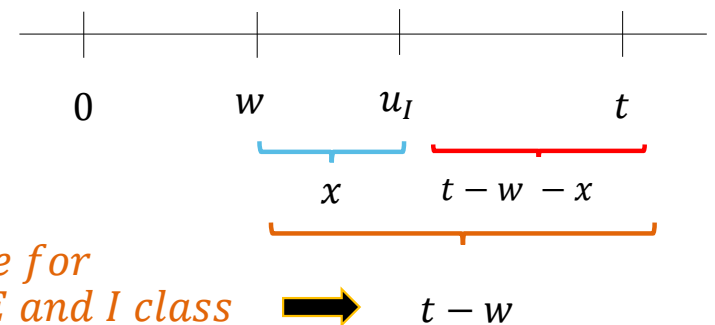
*Rate coming into the E class*

$$= \int_0^t \overbrace{-\mathbf{S}'(\mathbf{w})\boldsymbol{\psi}(t - \mathbf{w})}^{P(T_E + T_I > t - w)} dw + (1 - F(t))I_0 + G(t)E_0, \quad \leftarrow \text{Another Interpretation}$$

$$= \int_0^t \boldsymbol{\rho}(\mathbf{w})I(\mathbf{w})\boldsymbol{\psi}(t - \mathbf{w})dw + (1 - F(t))I_0 + G(t)E_0$$

$$\boldsymbol{\psi}(t) = \int_0^t (1 - F(t - v))\mathbf{g}(v)dv$$

*Minimum time for  
total time in E and I class*



## Two Equations for the Infective Class

$$I(t) = \int_0^t \mathbf{v(w)}(1 - F(t - w))dw + (1 - F(t))I_0 + G(t)E_0 \quad \leftarrow \text{Can be used to get estimation for } I(t)$$

$$I(t) = \int_0^t \mathbf{\rho(w)I(w)\psi(t - w)}dw + (1 - F(t))I_0 + G(t)E_0 \quad \leftarrow \text{Volterra equation of the second type}$$

↑  
*Use to estimate  $\rho(t)$*

# Using 1<sup>st</sup> interpretation of I(t)

*Rate coming into the I class*

$$I(t) = \int_0^t \overbrace{\mathbf{v}(\mathbf{w})}^{\text{Rate coming into the I class}} (1 - F(t - \mathbf{w})) d\mathbf{w} + (1 - F(t))I_0 + G(t)E_0 \quad \leftarrow \text{Can be used to get estimation for } I(t)$$

$\{Y_k\}_k$  : Cumulative number of individuals entering the I class on the interval  $(t_{k-1}, t_k]$

$E[Y_k] = \int_{t_{k-1}}^{t_k} \mathbf{v}(\mathbf{w}) d\mathbf{w}$  : Expected number of people per time entering the I class on the interval  $(t_{k-1}, t_k]$

Let  $s_\tau \in (t_{k-1}, t_k]$

Estimation for  $I(t)$ :  $\hat{I}(s_\tau) = \sum_{j: t_j \leq s_\tau} \frac{Y_j}{t_j - t_{j-1}} \int_{t_{j-1}}^{t_j} (1 - F(t - \mathbf{w})) d\mathbf{w} + \frac{Y_k}{t_k - t_{k-1}} \int_{t_{k-1}}^{s_\tau} (1 - F(t - \mathbf{w})) d\mathbf{w}$

# Using 2<sup>nd</sup> Interpretation of I(t)

$$I(t) = \int_0^t \rho(w) I(w) \psi(t-w) dw + (1 - F(t)) I_0 + G(t) E_0$$

← *Volterra equation of the second type*

*Trapezoidal Method*

$$I(s_k) = \frac{\delta}{2} [I(0)\rho(0)\psi(s_k) + 2 \sum_{j=1}^{k-1} I(s_j)\rho(s_j)\psi(s_k - s_j) + I(s_k)\rho(s_k)\psi(s_0)] + (1 - F(t_k)) I_0 + G(t_k) E_0$$

*Linear in elements of  $\rho$*

$$\delta = s_j - s_{j-1}$$

$$\vec{I} = H \vec{\rho} + \vec{v}, \quad \vec{v} = \overrightarrow{1 - F} I_0 + \vec{G} E_0$$



# Constraints

---

Use estimation for  $I$ :  $\hat{I}$

$$\hat{I} = H\vec{\rho} + \vec{v}, \quad \vec{v} = \overrightarrow{1 - F} I_0 + \vec{G} E_0$$



$$\begin{aligned} H_{i0} &= \frac{\delta}{2} \hat{I}(s_0) \psi(s_i) \\ H_{ij} &= \delta \hat{I}(s_j) \psi(s_i - s_j), & j < i \\ H_{ii} &= \frac{\delta}{2} \hat{I}(s_i) \psi(s_0) \\ H_{ij} &= 0, & j \geq i \end{aligned}$$

# Another Linear Relationship

$$E[Y_k] = \int_{t_{k-1}}^{t_k} v(w) dw :$$

$$= \int_{t_{k-1}}^{t_k} v(w) dw$$

$$= \int_{t_{k-1}}^{t_k} \int_{t_0}^w S'(u) g(w-u) du dw$$

$$= \int_{t_{k-1}}^{t_k} \rho(u) I(u) \int_{\max(u, t_{k-1})}^{t_k} g(w-u) du dw, \quad \forall u \in [0, t_k]$$

*Expected number of people per time entering the I class on the interval  $(t_{k-1}, t_k]$*

*Trapezoidal Method*

$$R_k(u) = G(t_k - u) - G(\max(t_{k-1}, u) - u)$$

$$E[Y_k] = \frac{\delta}{2} [\rho(0)I(0)R_k(0) + 2 \sum_{i=1}^{\frac{k}{\delta}-\delta} \rho(s_i)I(s_i)R_k(s_i) + \rho(\frac{s_k}{\delta})I(\frac{s_k}{\delta})R_k(\frac{s_k}{\delta})]$$

$$\vec{Y} = M \vec{\rho}$$

# Optimization Problem

$$\vec{Y} = M \vec{\rho}$$

Regularization with the 2<sup>nd</sup> difference matrix to control smoothness

$$D_{ij} = \begin{cases} 1 & |i - j| = 1 \\ -2 & i = j \\ 0 & \text{otherwise} \end{cases}$$

$$\hat{I} = H \vec{\rho} + \vec{v},$$

$$\vec{v} = \overrightarrow{1 - F} I_0 + \vec{G} E_0$$

---

Questions?



# Comparing RDX and HMX reduced chemistry models applied to 1-D hot spot simulations

Michael N. Sakano\*, Alejandro Strachan\*, Edward M. Kober†

\*School of Materials Engineering and Birck Nanotechnology Center, Purdue University

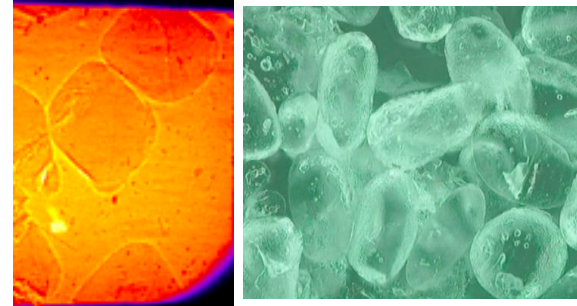
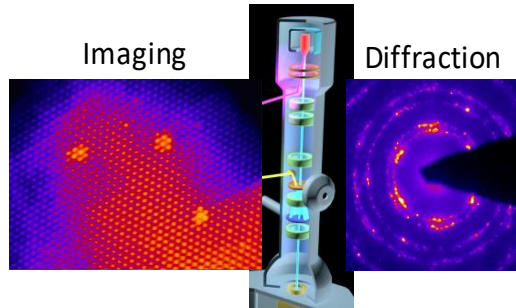
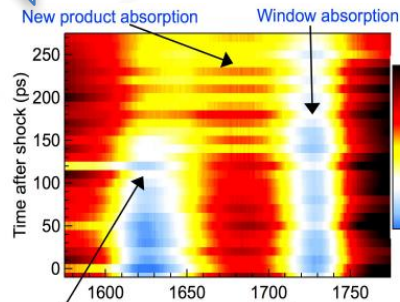
†T-1 Division, Los Alamos National Laboratory

Funding provided by High Explosives Grand Challenge

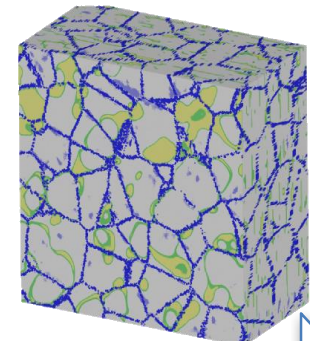
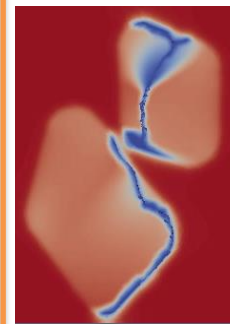
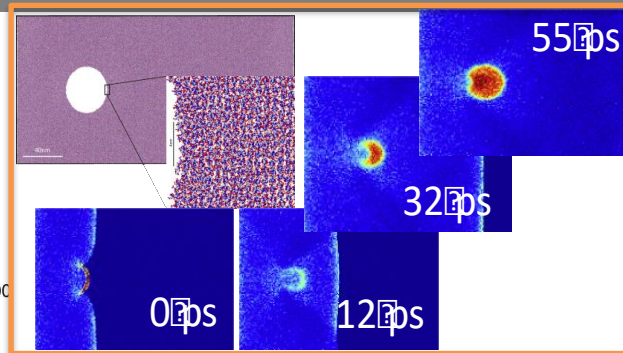
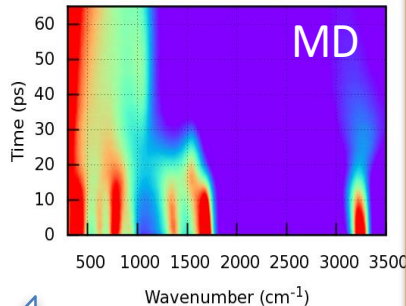
08/19/2020

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# Complex chemical initiation coupled across scales



Laser shocks Laser heating Vibrations Kolsky bar Gas gun

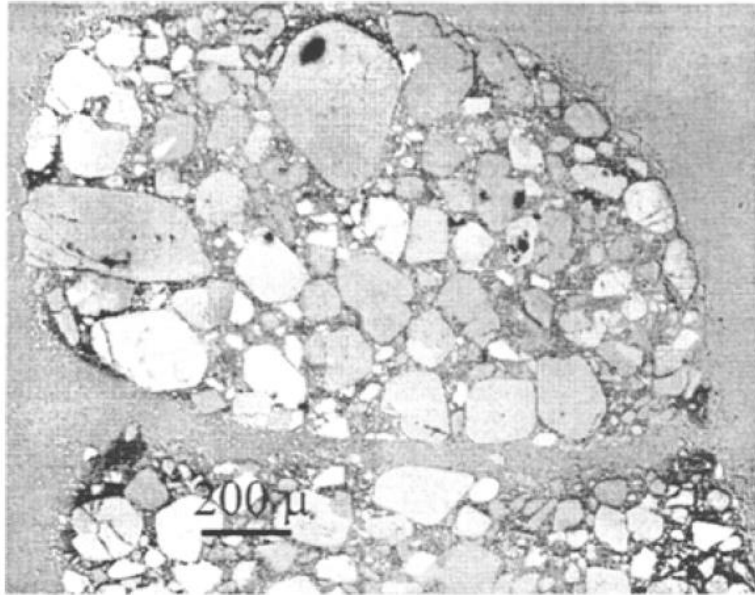


ONR MURI Topic 21 – Purdue, Stanford, LANL – PI: Strachan  
PCP@Xtreme - Predictive Chemistry & Physics at Extreme Temperature and Pressure

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Slide 2

# Discrepancies in initiating single crystal HMX

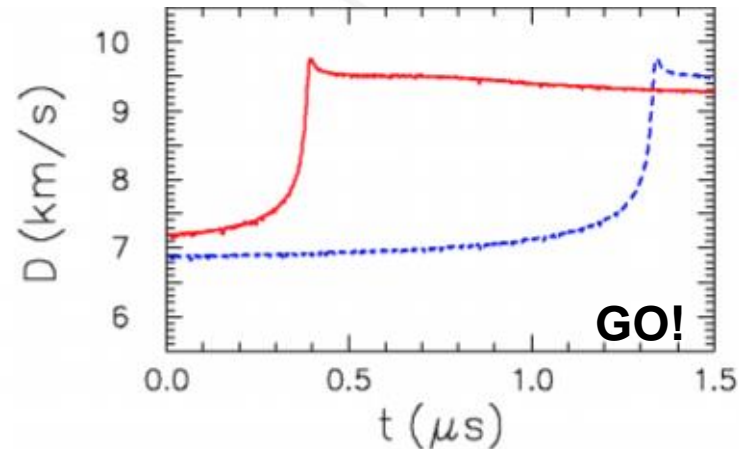


**PBX 9501 Molding Powder**

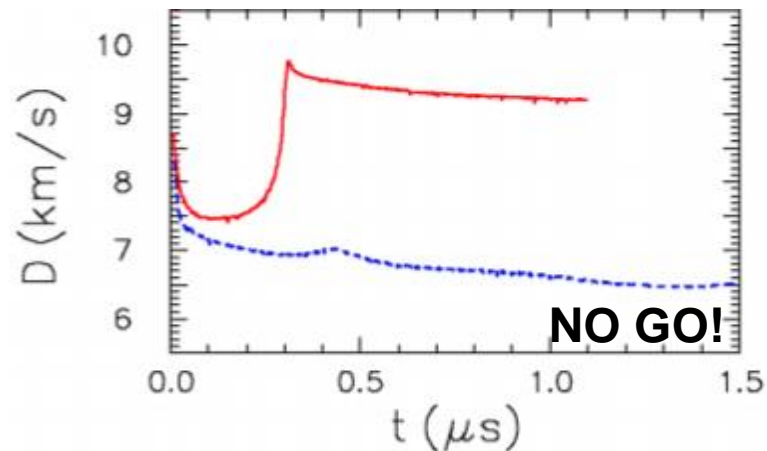
Skidmore, C. B.; Phillips, D. S.; Son, S. F.; Asay, B. W. Characterization of HMX Particles in PBX 9501. *AIP Conf. Proc.* **1998**, 429 (1), 579–582. <https://doi.org/10.1063/1.55666>.

Menikoff, R. Micro-gap Experiments and Insensitive Explosives. *AIP Conf. Proc.* **2007**, 955 (1), 967–970. <https://doi.org/10.1063/1.2833290>.

**Flyer plate**



**PBX 9404 detonation wave**

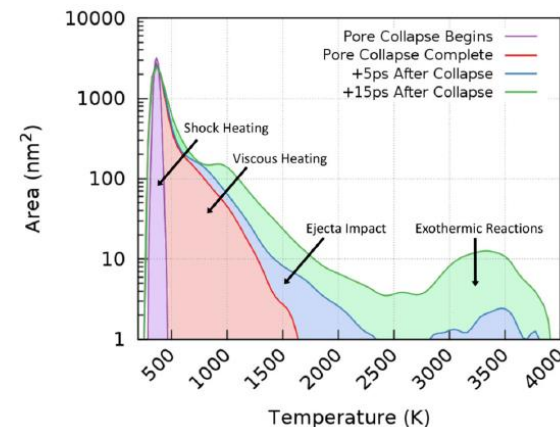
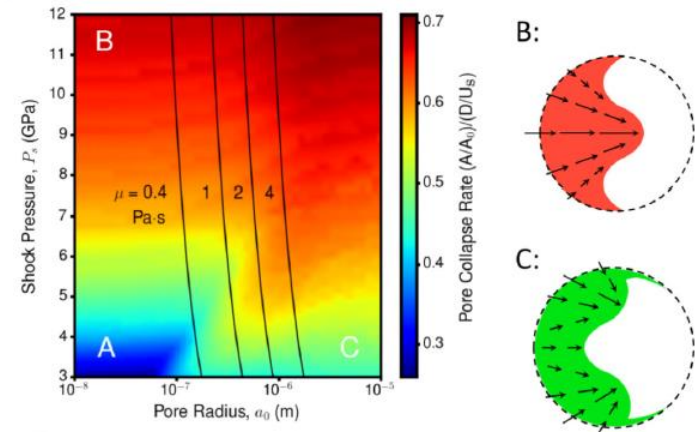
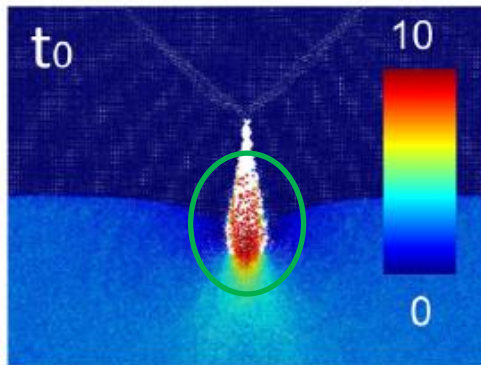
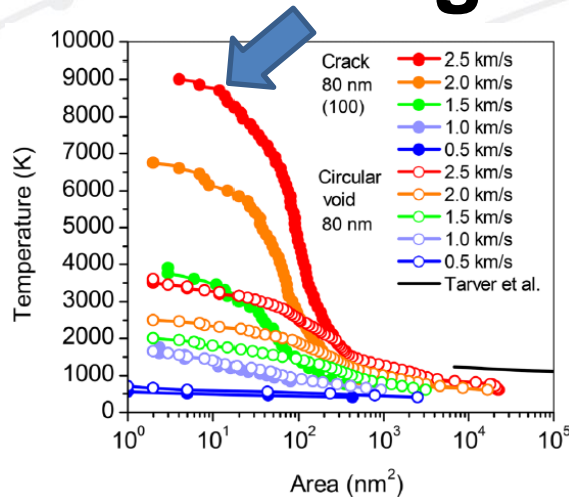


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Slide 3



# Jetting leads to critical hot spots



Li, C.; Hamilton, B. W.; Strachan, A. Hotspot Formation Due to Shock-Induced Pore Collapse in 1,3,5,7-Tetranitro-1,3,5,7-Tetrazoctane (HMX): Role of Pore Shape and Shock Strength in Collapse Mechanism and Temperature. *J. Appl. Phys.* **2020**, 127 (17), 175902. <https://doi.org/10.1063/5.0005872>.

Wood, M. A.; Kittell, D. E.; Yarrington, C. D.; Thompson, A. P. Multiscale Modeling of Shock Wave Localization in Porous Energetic Material. *Phys. Rev. B* **2018**, 97 (1), 014109. <https://doi.org/10.1103/PhysRevB.97.014109>.

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Pore collapse leads to critical hot spots

Slide 4



# Hot spot model of detonation

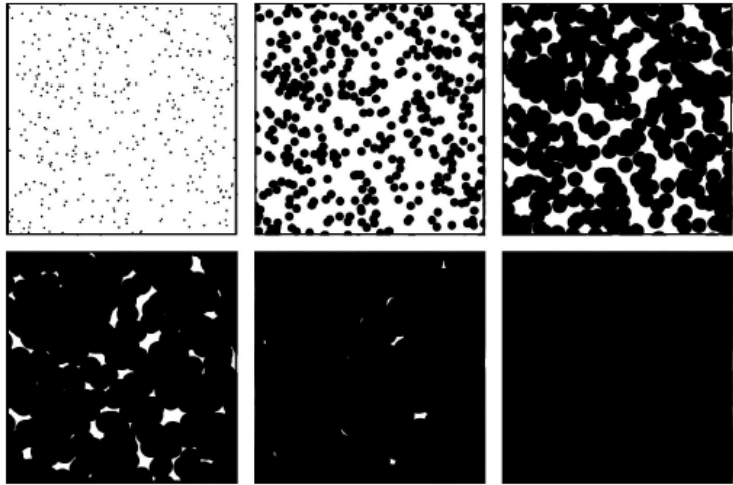


FIGURE 1. 2D rendering of the standard SHSM.

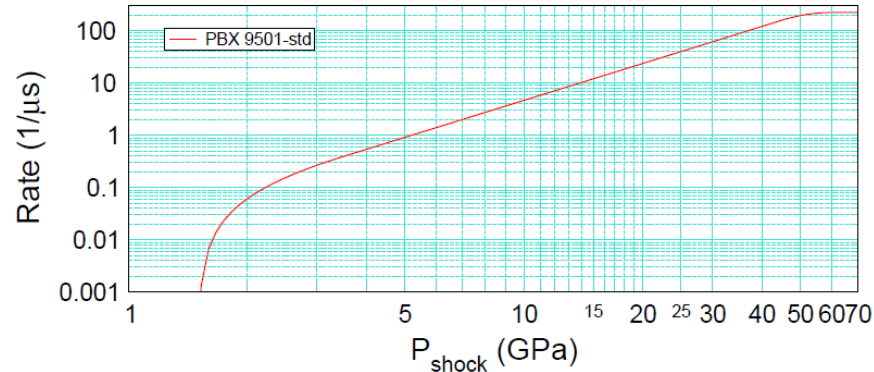


Figure 3: SURF burn rate for PBX 9501.

SURF parameters

$p_{scale}$	1 GPa
$t_{scale}$	1 $\mu s$
$P_0$	1.5
$P_{low}$	3.5
$P_1$	45.
$P_{hi}$	60.
$C$	0.0208
$f_n$	2.35
$n$	1.0
$n_{hi}$	0.0
$P_{burn}$	35.

Critical parameters:

- 1) Number, size and distribution of initial hot spots
- 2) Growth (deflagration) rate of hot spots vs local conditions

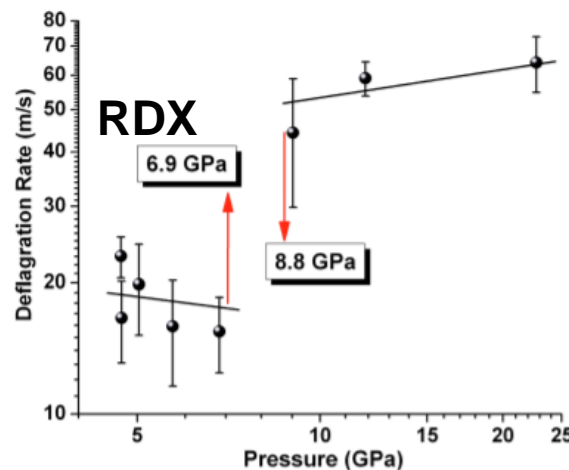
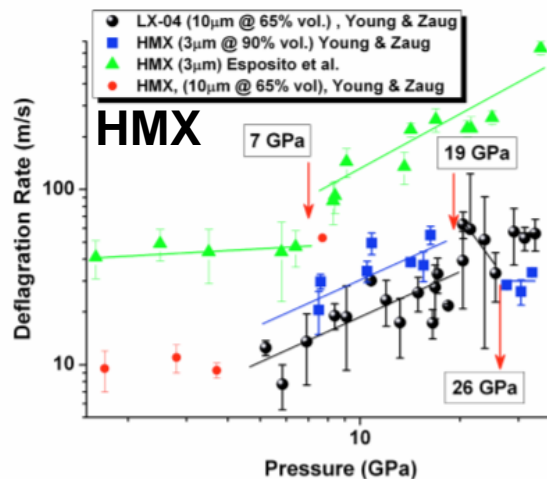
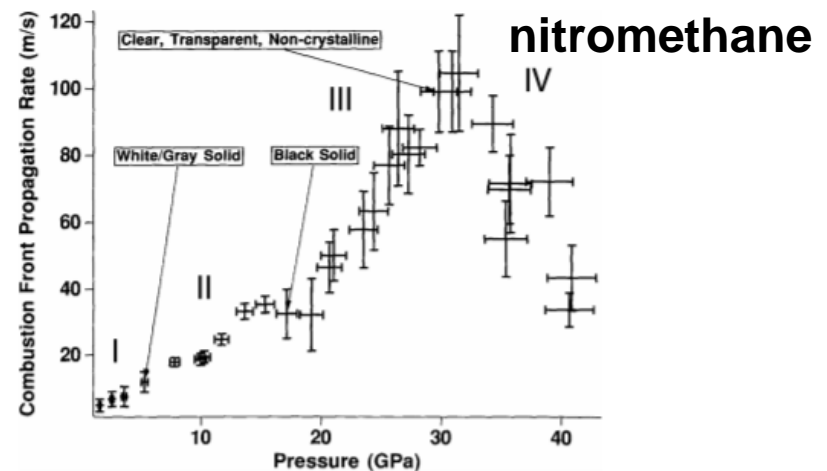
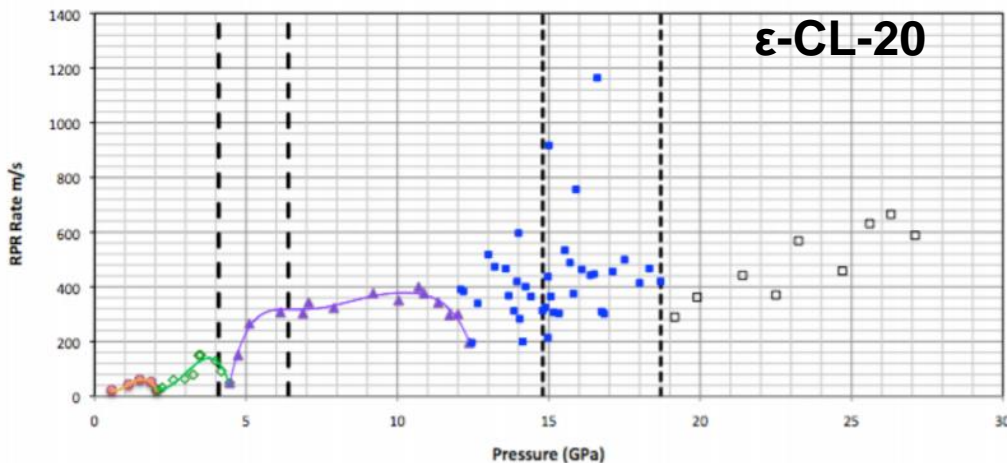
Hill, L. G. The Shock-Triggered Statistical Hot Spot Model. *AIP Conf. Proc.* **2012**, 1426 (1), 307–310. <https://doi.org/10.1063/1.3686280>.

Menikoff, R. Effect of Heterogeneities on the Reaction-zone Profile of a Propagating Detonation Wave. *LA-UR-20-24842*. **2020**.

UNCLASSIFIED

Focus on deflagration rate versus temperature and pressure

# Experimental measurements of deflagration rates

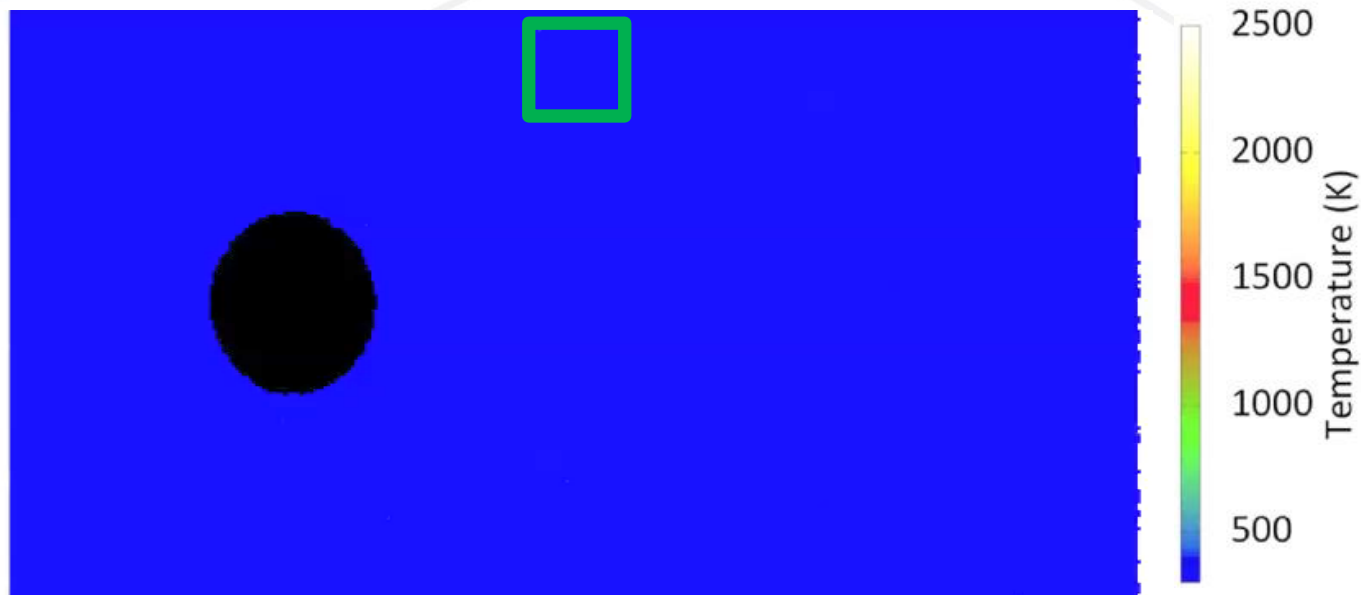


Zaug, J. M.; Young, C. E.; Long, G. T.; Maienschein, J. L.; Glascoe, E. A.; Hansen, D. W.; Wardell, J. F.; Black, C. K.; Sykora, G. B. Deflagration Rates of Secondary Explosives under Static Mpa—Gpa Pressure. *AIP Conf. Proc.* **2009**, 1195 (1), 420–423.  
<https://doi.org/10.1063/1.3295162>.  
 Zaug, J. M.; Foltz, M. F.; Hart, E. *Deflagration Rates and Molecular Bonding Trends of Statically Compressed Secondary Explosives*; LLNL-PROC-425270; **2010**.

UNCLASSIFIED

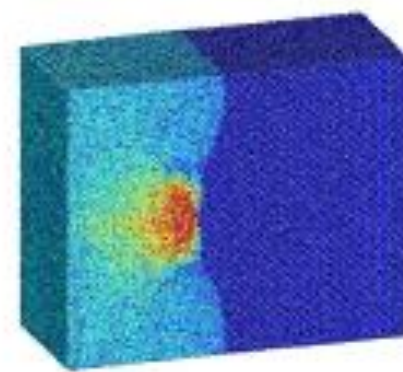
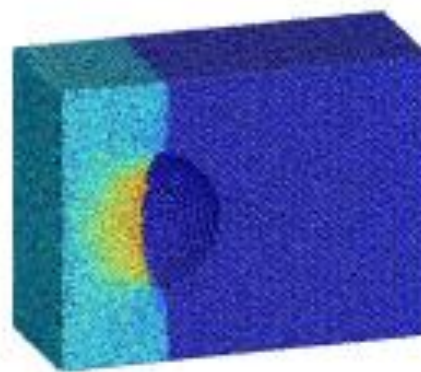
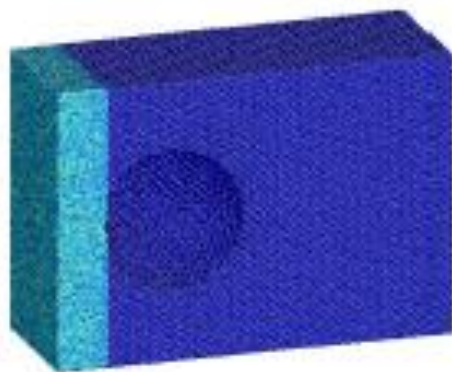
Difficult to determine reaction deflagration rate from experiments

# Deflagration wave from hot spot



Wood, M. A.; Cherukara, M. J.; Kober, E. M.; Strachan, A. Ultrafast Chemistry under Nonequilibrium Conditions and the Shock to Deflagration Transition at the Nanoscale. *J. Phys. Chem. C* **2015**, 119 (38), 22008–22015.

Shan, T.-R.; Thompson, A. P. Shock-Induced Hotspot Formation and Chemical Reaction Initiation in PETN Containing a Spherical Void. *J. Phys. Conf. Ser.* **2014**, 500 (17), 172009.



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Slide 7

# Molecular Dynamics

Initial atomic positions  $\vec{R}_i$  and velocities  $\vec{V}_i$

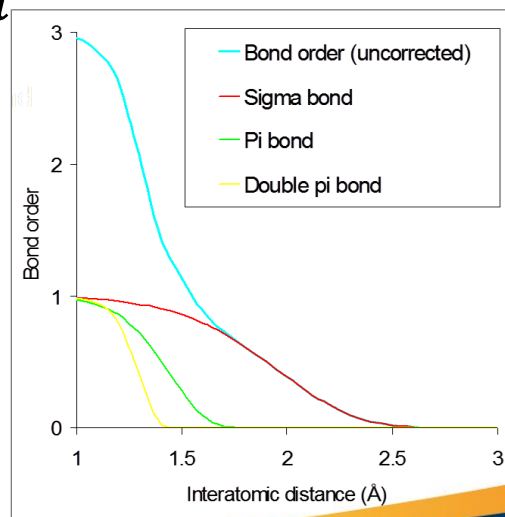
Compute forces on atoms  $\vec{F}_i = -\vec{\nabla}_{R_i} V(\{R_j\})$

Solve Newton's Second Law  $\dot{\vec{V}}_i = \frac{\vec{F}_i}{M_i}$   $\dot{\vec{R}}_i = \vec{V}_i$

Integrate classical equations of motion

Bond order reactive potential

[1] Senftle, T. P. et al. *npj Comp. Mat.* **2**, (2016)



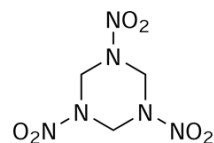
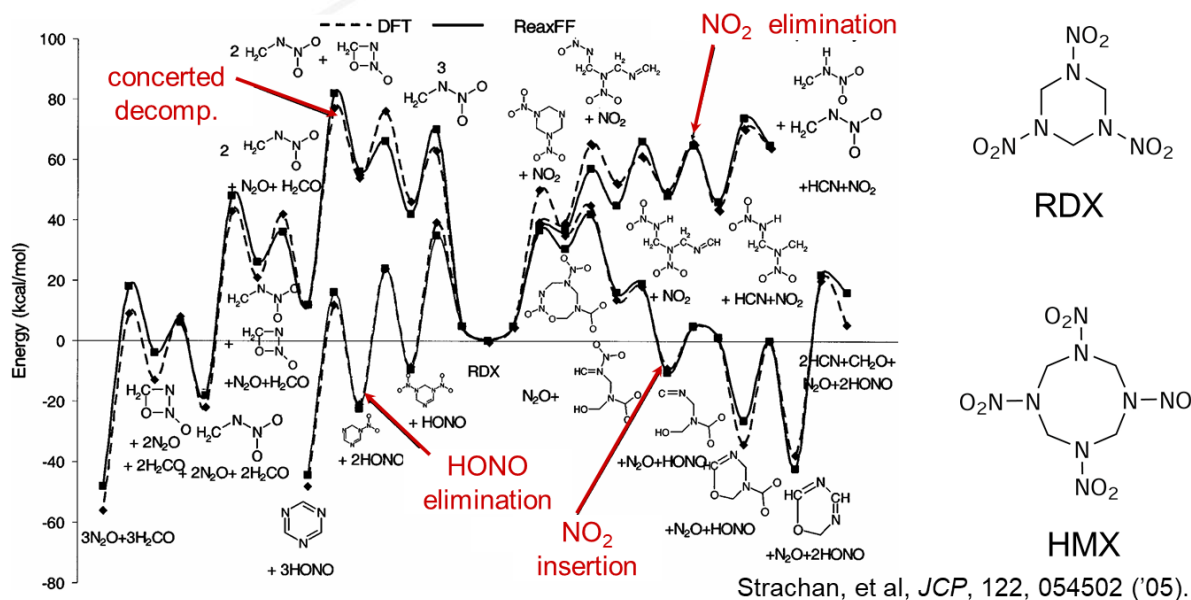
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What is a good potential to use?

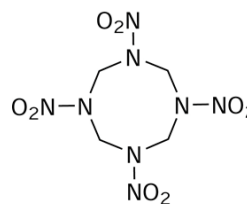


# ReaxFF / ReaxFF-Ig

## Unimolecular Decomposition ...comparison between ReaxFF and B3LYP/6-31G\*\*

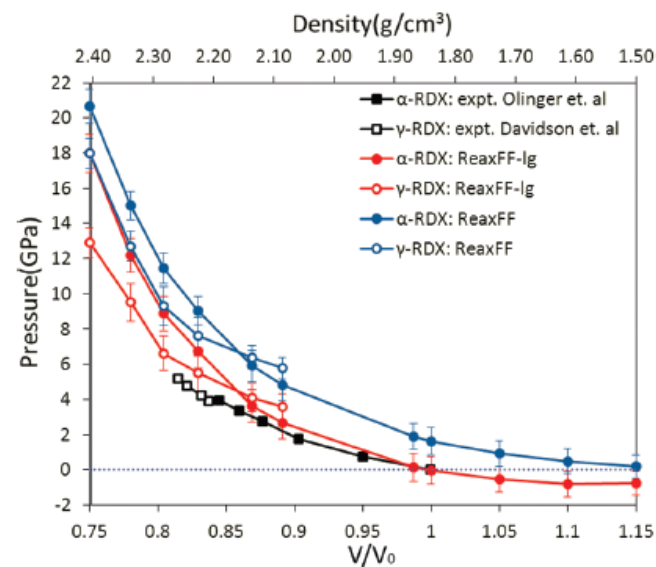


RDX



HMX

L. Lui et al., *J Phys. Chem. A*,  
2011, **115**, 11016



Lack of bi-molecular reaction schemes

Low gradient vdw term: improved solids' densities

Calibrated to ground state and transition state geometries

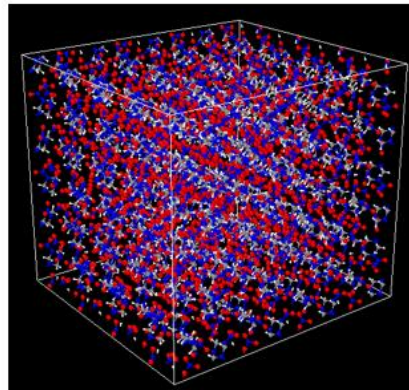
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# Sample Preparation

- **2** systems
  - RDX and HMX
- **10** densities
  - $V/V_0 = 0.65 - 1.50$
- **7** temperatures (NVT)
  - 1200 - 3000 K

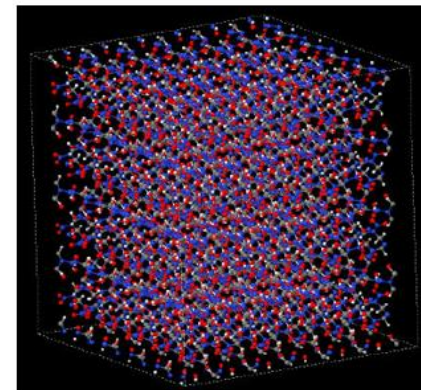
## RDX

- 4x4x3 unit cells
- 8064 atoms
  - 384 molecules
- $V = 76.25 \text{ nm}^3$



## HMX

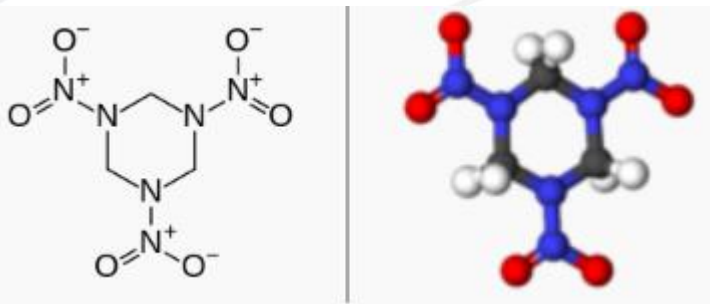
- 6x5x4 unit cells
- 6720 atoms
  - 240 molecules
- $V = 75.45 \text{ nm}^3$



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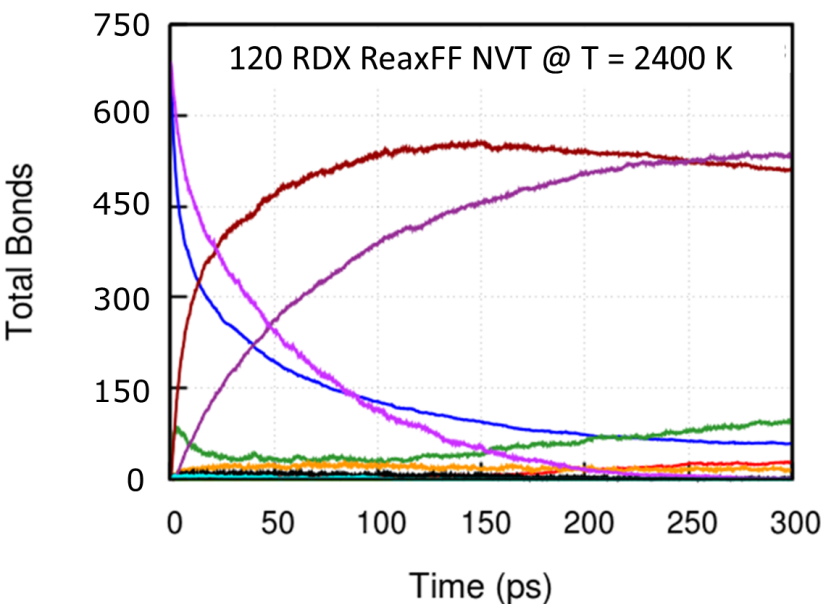
Slide 10

# Coordination Geometry Analysis

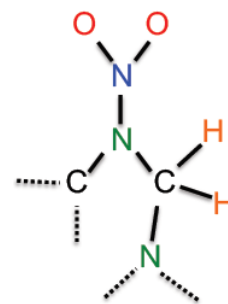


<https://ipfs.io/ipfs/QmXoypizjW3WknFiJnKLwHCnL72vedxjQkDDP1mXWo6uco/wiki/RDX.html>

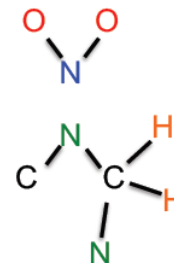
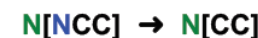
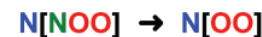
- Assume all atoms have a maximum of four bonds
- 280 different geometry configurations for *HCNO* atom types
- Based on group additivity approach of Benson et al.



— H-H  
— H-C  
— H-N  
— H-O  
— C-C  
— C-N  
— C-O  
— N-N  
— N-O  
— O-O



Nitro group dissociation



Cohen, N.; Benson, S. W. Estimation of Heats of Formation of Organic Compounds by Additivity Methods. *Chem. Rev.* **1993**, 93 (7), 2419–2438. <https://doi.org/10.1021/cr00023a005>.

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Evolution of bond types – coupled chemistry

# Dimensionality Reduction – Non-negative matrix factorization

BONDS

	H-H	H-C	H-N		O-O
Time (ps)	0	720	0	...	0
	0	694	0	...	0
	0	688	0	...	0
	2	680	6	...	0
	⋮	⋮	⋮	...	⋮
	10	60	118	...	16

 $\approx$

Example:  
Decomposition of  
120 RDX molecules

<http://scikit-learn.org/stable/modules/generatord/sklearn.decomposition.NMF.html>

	Components (N)					BONDS					
	H-H	H-C	H-N			H-H	H-C	H-N		O-O	
Time (ps)	0.0	0.0	1.0	0.0	•	0	0	0	0	...	12
	0.0	0.0	1.0	0.0		0	0	0	0	...	10
	0.0	0.0	1.0	0.0		0	680	0	0	...	0
	0.0	0.0	0.9	0.1		10	0	0	0	...	0
	⋮	⋮	⋮	⋮							
	0.9	0.1	0.0	0.0							

(Concentration)

(Composition)

Components (N)

E. M. Kober,  
personal  
communication &  
manuscript in  
preparation

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Matrix decomposition for concentrations

Slide 12

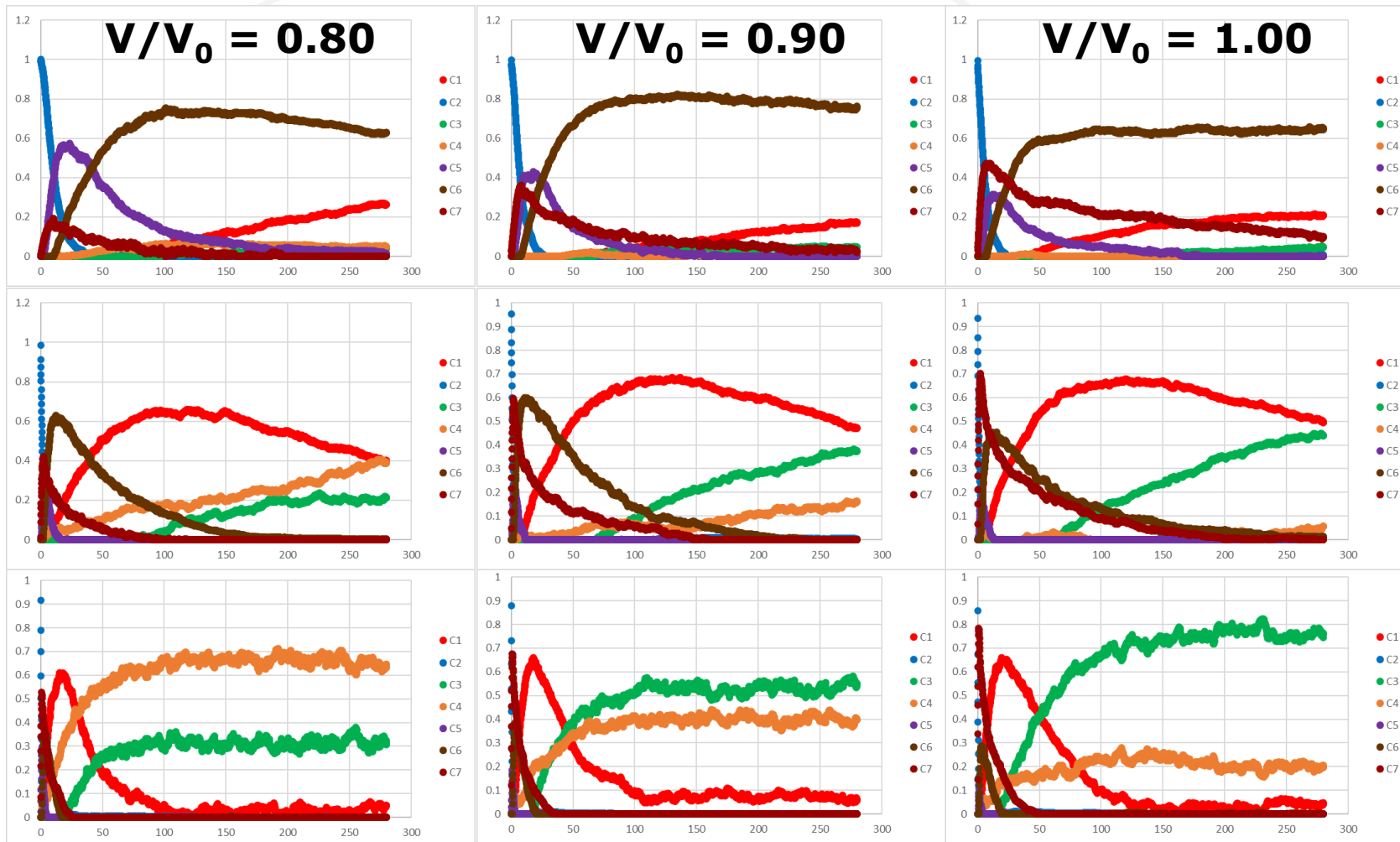


# Application of NMF - RDX

1800 K

2400 K

3000 K



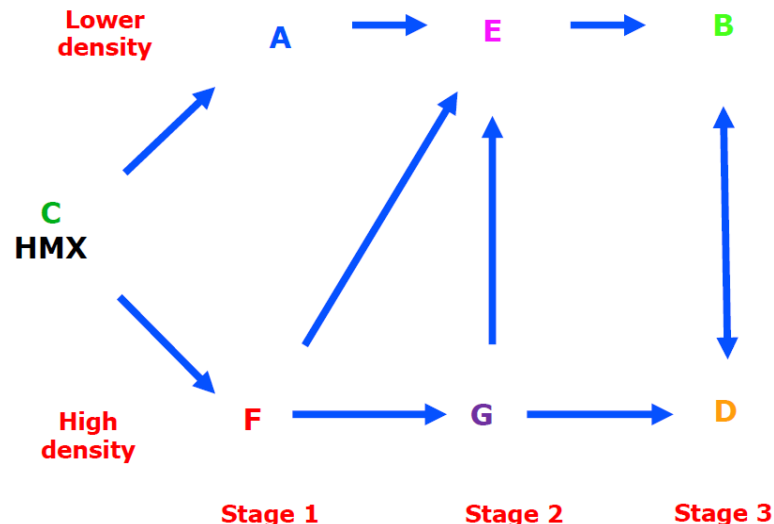
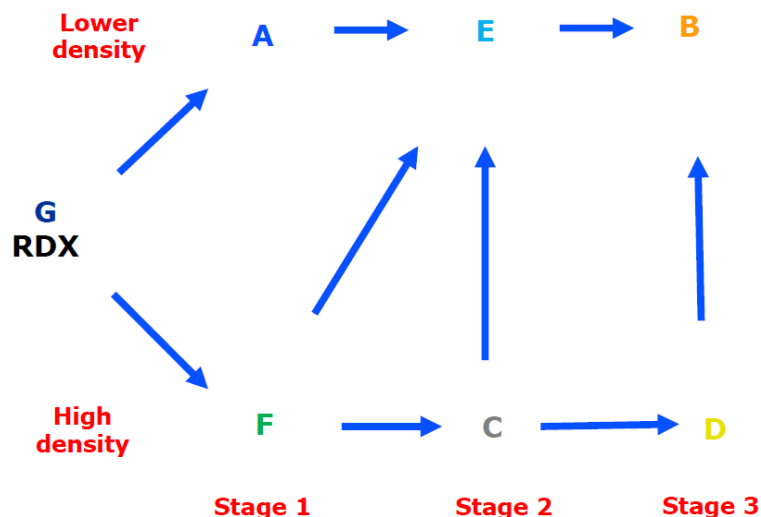
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Concentrations as a function of T and P

Slide 13

# Comparison of 7-Component reduced models

$$V/V_0 = 0.65 - 1.50$$



Similarity in pathways but differences in chemistry:

Lower density path: 1)  $\text{NO}_2$  dissociation and ring break-up; 2)  $\text{N}_2$  and  $\text{H}_2\text{O}$  production; 3) Carbon oxidation

Higher density path: 1) Associative reaction to form  $\text{NO}_3$ , ring stays intact, RDX forms triazine (stable intermediate); 2) RDX forms  $\text{N}_2$  and  $\text{H}_2\text{O}$ , while HMX oxidizes carbon in the ring; 3) RDX oxidizes the triazine, while HMX produces  $\text{N}_2$  and  $\text{H}_2\text{O}$

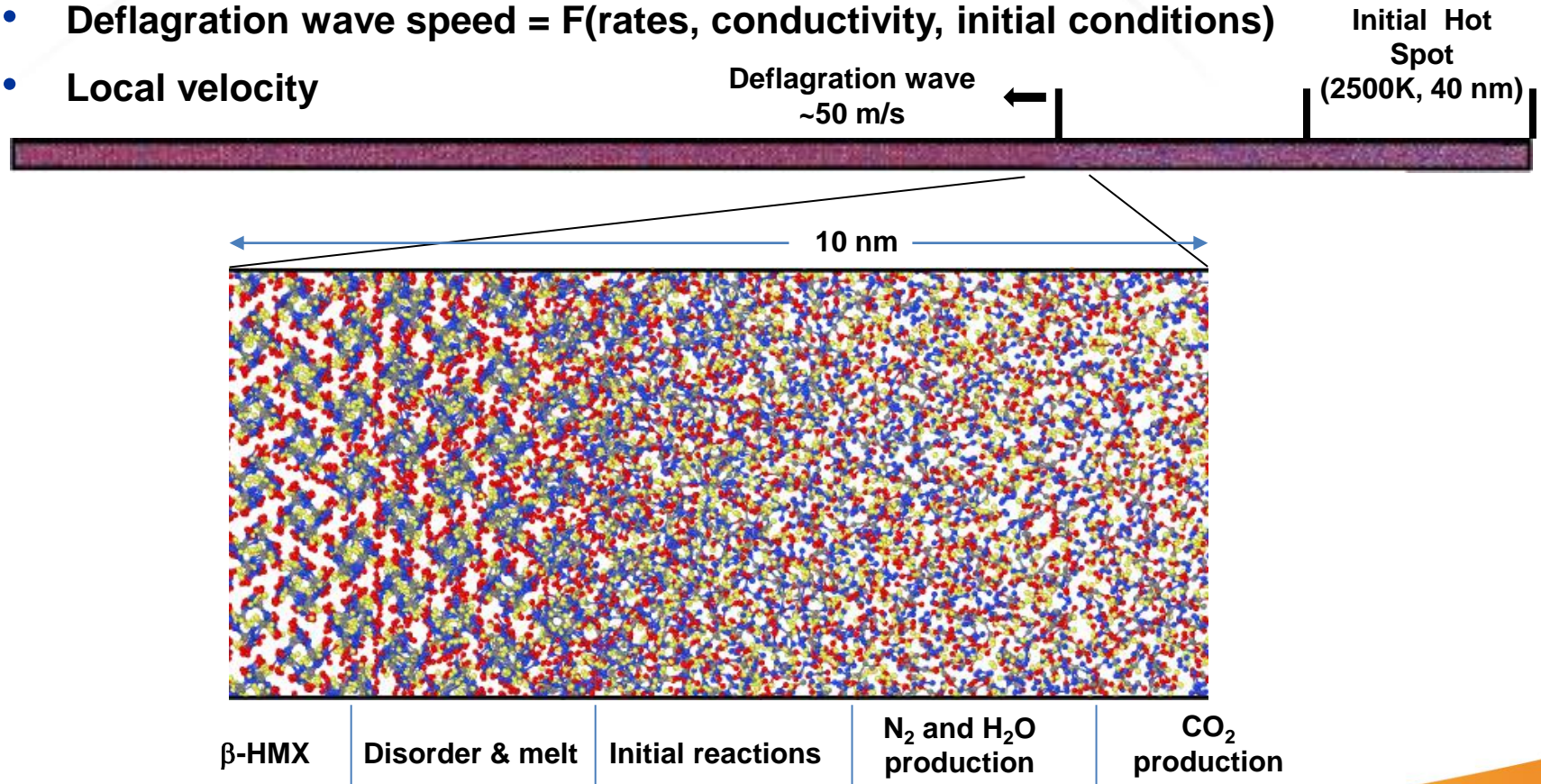
Extracting rate constants from these schemes assuming Arrhenius kinetics

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Slide 14

# 1-D Deflagration Simulations on HPC resources

- Typical simulation: 5 x 5 x 200 nm ~ 1M atoms (20-100 nodes, 20-100 hours)
- Thermalize selected region for ignition: 10-50 nm, 1500-2500 K
- Downstream conditions: ambient, 600K 0.8 GPa
- Deflagration wave speed = F(rates, conductivity, initial conditions)
- Local velocity



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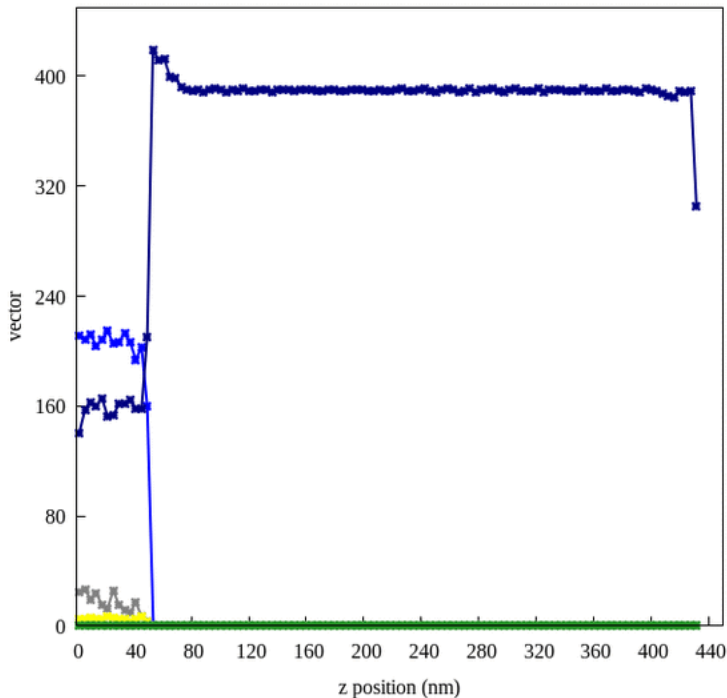
Can track regions of local chemistry

# 1-D Thermal Sticks, 40 nm @ 2500 K

- Fixed wall on both ends
- Projection onto 7 component cook-off model

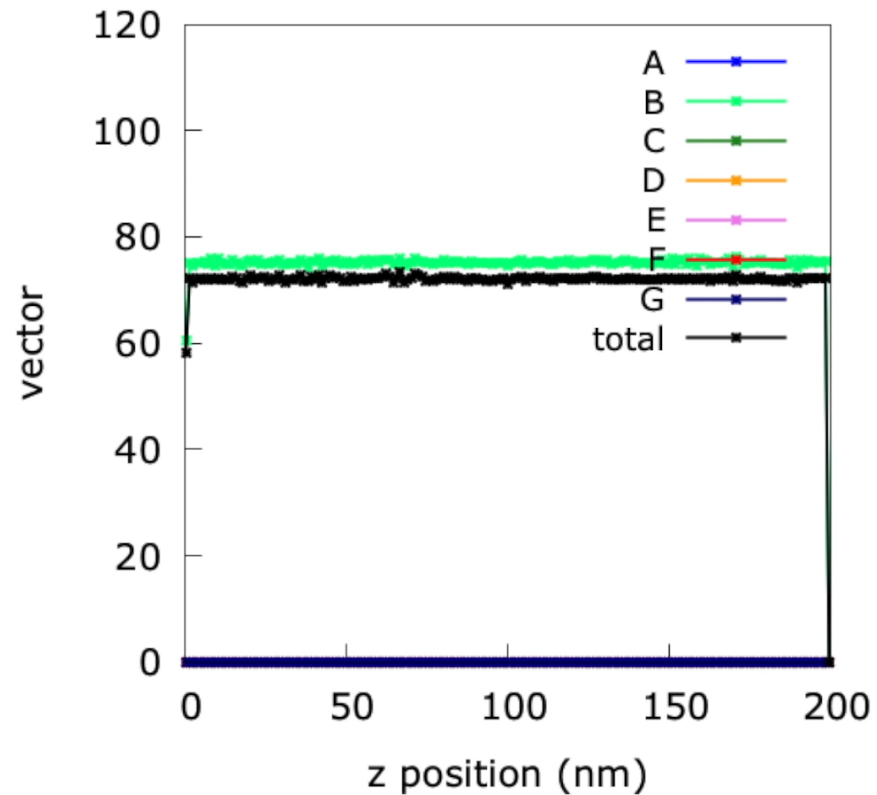
## RDX

Frame 50



## HMX

Frame 0

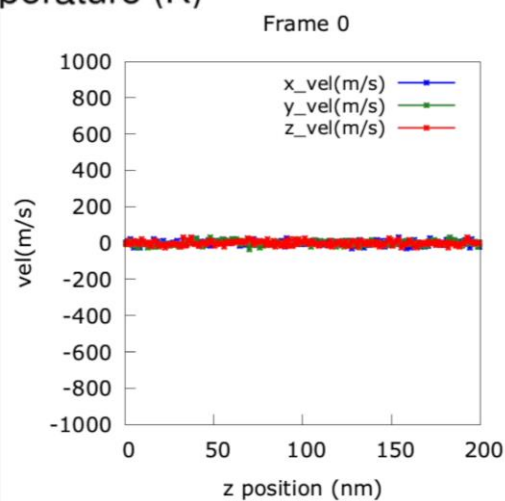
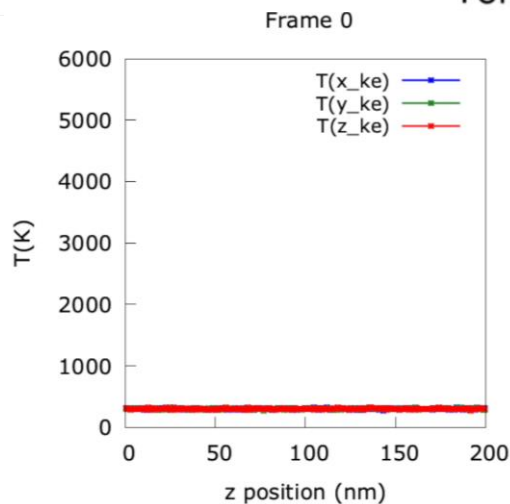
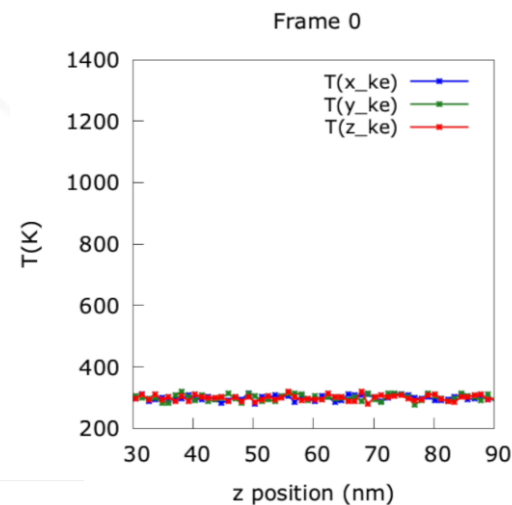
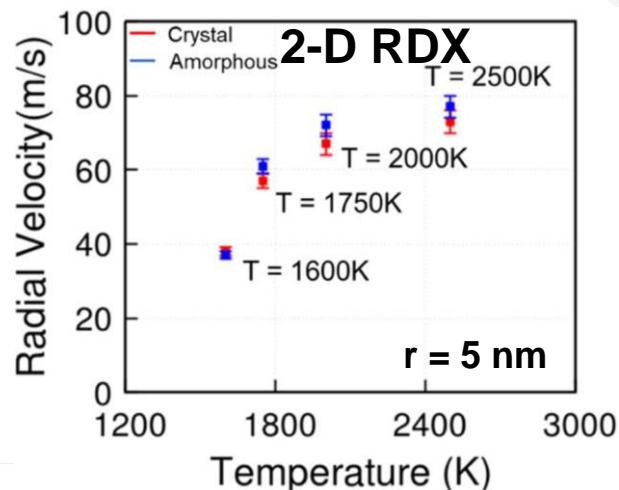
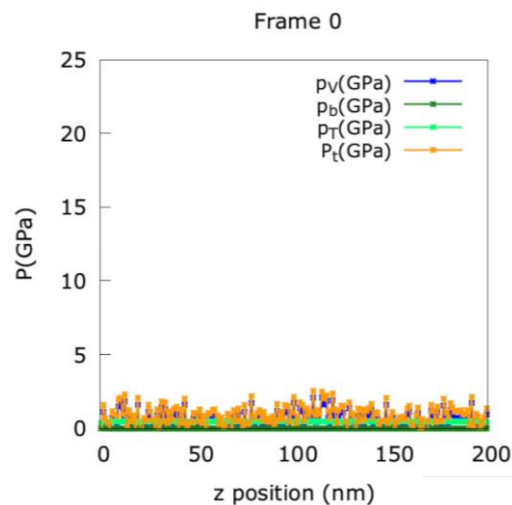


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Apparent differences in local chemistry



# 1-D Thermal Sticks, 40 nm @ 2500 K for HMX



Sakano, M.; Hamilton, B.; Islam, M. M.; Strachan, A. Role of Molecular Disorder on the Reactivity of RDX. *J. Phys. Chem. C* **2018**, 122 (47), 27032–27043. <https://doi.org/10.1021/acs.jpcc.8b06509>.

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Measure reaction front velocity vs temperature and size

# Summary

- Refining reduced-order chemistry model for RDX and HMX (Arrhenius fits, EOS)
- Map small-scale cook-off simulations to HPC runs of 1-D thermal stick variations
- Extracting deflagration wave speeds
- Will map to 1-D gaps (voids) for various sizes and shock speeds: common chemistry to deflagration?
- Improve modeling of chemistry by implementing Quantum Thermal Bath (redistributes energies)

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Slide 18

# Reducing Complexity of Variational Quantum Eigensolvers for Quantum Chemistry

July 14<sup>th</sup>, 2020



James Sud (Post-Bachelor)

University of California, Berkeley  
Quantum Computing Summer  
School



Nikolay Tkachenko (2<sup>nd</sup> year Ph.D student)

Utah State University  
20200056DR Quantum Chemistry with  
Quantum Computers

Mentors: [Lukasz Cincio \(T-4\)](#), [Yu Zhang \(T-1\)](#), [Pavel Dub \(C-IIAC\)](#)

# Agenda

July 14, 2020



- Intro to quantum chemistry and computing
- Quantum chemistry simulations on quantum computers
- The Variational Quantum Eigensolver (VQE)
- ROSALIN optimizer for reducing measurement counts
- Choosing an ansatz
- Qubit permutations for reducing circuit depth and measurement counts
- Preliminary Results



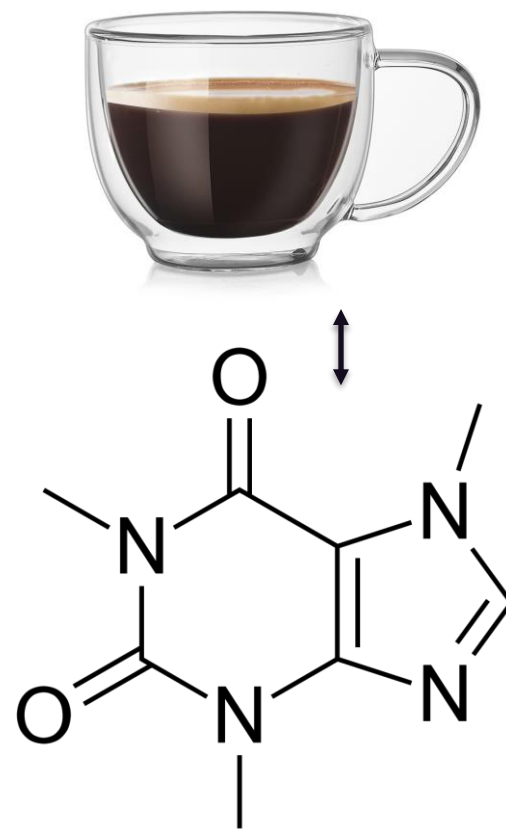
# Introduction

**DOE pushes for quantum computing (2018):** Quantum chemistry is immediate applications of quantum computing directly impacting LANL present and future missions

**Goal:** Find the electronic ground state of chemical systems

**Motivation:** drug discovery, catalysis, materials design, etc.

**Challenge:** The **exact** simulation of the Caffeine molecule in sto-3g basis is a **daunting task for any existing supercomputer** (Diagonalizing and storing a matrix of size  $6 \times 10^{43}$ )



[https://en.wikipedia.org/wiki/Caffeine#/media/File:Caffeine\\_structure.svg](https://en.wikipedia.org/wiki/Caffeine#/media/File:Caffeine_structure.svg)

# Introduction

**Solution:** Use quantum computers  
– each one-electron spin-orbital is encoded in one qubit

**Example:** LiH, 4 electrons, 495 possible configurations (sto-3g basis)

$$|\Psi\rangle = c_0 \begin{array}{c} \text{---} \\ \text{---} \\ \uparrow\downarrow \\ \uparrow\downarrow \end{array} + c_1 \begin{array}{c} \text{---} \\ \uparrow \\ \text{---} \\ \uparrow\downarrow \\ \uparrow\downarrow \end{array} + c_2 \begin{array}{c} \uparrow \\ \text{---} \\ \text{---} \\ \uparrow\downarrow \\ \uparrow\downarrow \end{array} + \dots$$

$$|\Psi\rangle = c_0 |111100\dots 0\rangle + c_1 |111010\dots 0\rangle$$

CPU/GPU



scales exponentially



$$\hat{H}|\Psi\rangle = E|\Psi\rangle$$



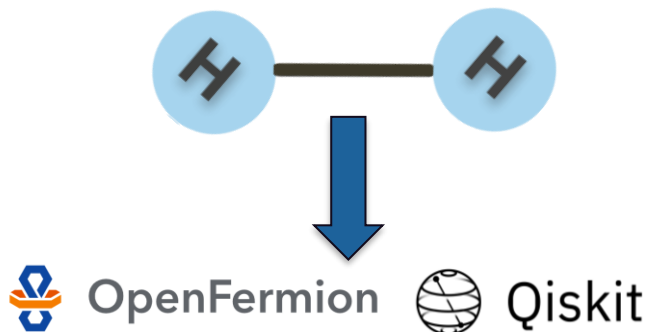
scales polynomially



QPU

# Quantum Chemistry on Quantum Computers

## Start with Molecular Coordinates

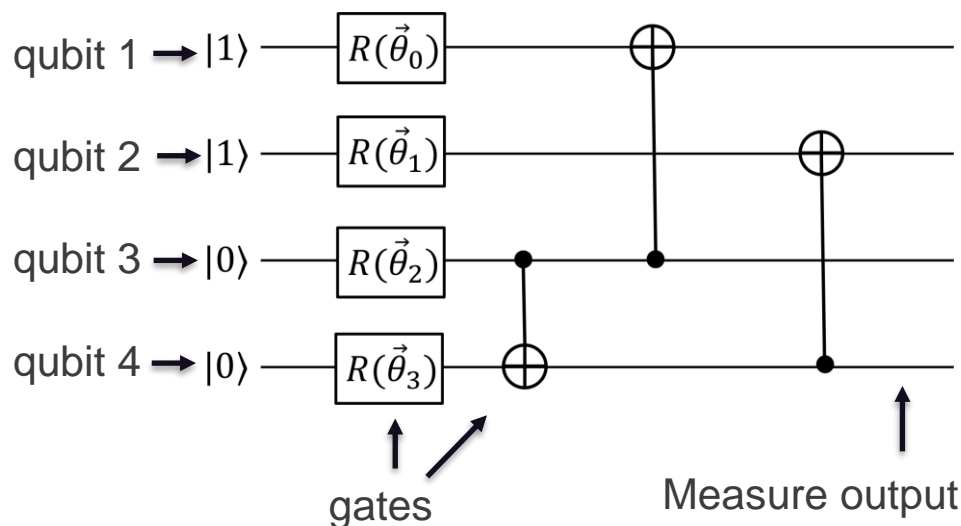


## Convert to qubit Hamiltonian

-0.09064961228803427 [] +  
0.04523355281799211 [X0 Z1 X2] +  
0.04523355281799211 [X0 Z1 X2 Z3] +  
0.04523355281799211 [Y0 Z1 Y2] +  
0.04523355281799211 [Y0 Z1 Y2 Z3] +  
0.17217557977330947 [Z0] +  
0.17217557977330947 [Z0 Z1] +  
0.16614307978715506 [Z0 Z1 Z2] +  
0.16614307978715506 [Z0 Z1 Z2 Z3] +  
0.12090952696916293 [Z0 Z2] +  
0.12090952696916293 [Z0 Z2 Z3] +  
0.16892496212247482 [Z1] +  
-0.22572827882218005 [Z1 Z2 Z3] +  
0.17464093706423592 [Z1 Z3] +  
-0.22572827882218005 [Z2]

## Prepare quantum circuit and measure with respect to qubit Hamiltonian

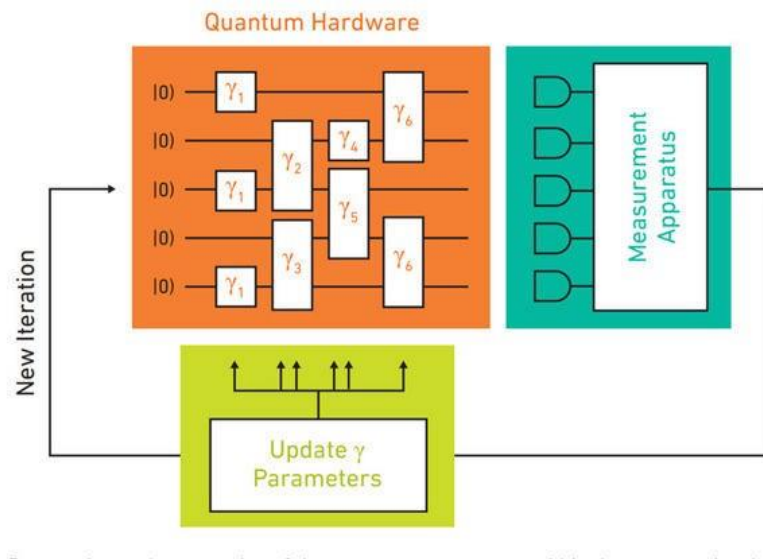
H2 molecule in sto-3g basis



<https://github.com/quantumlib/OpenFermion>  
<https://ai.googleblog.com/2016/07/towards-exact-quantum-description-of.html>  
<https://www.eagleresearchllc.com/products/uncategorized/h2-true-molecular-hydrogen/>  
<https://user-images.githubusercontent.com/50955024/69568348-ac558780-0f78-11ea-9242-7ac8e8ba450a.png>

# ROSALIN Optimizer Designed for Chemistry Hamiltonians

## Variational Quantum Eigensolver (VQE)



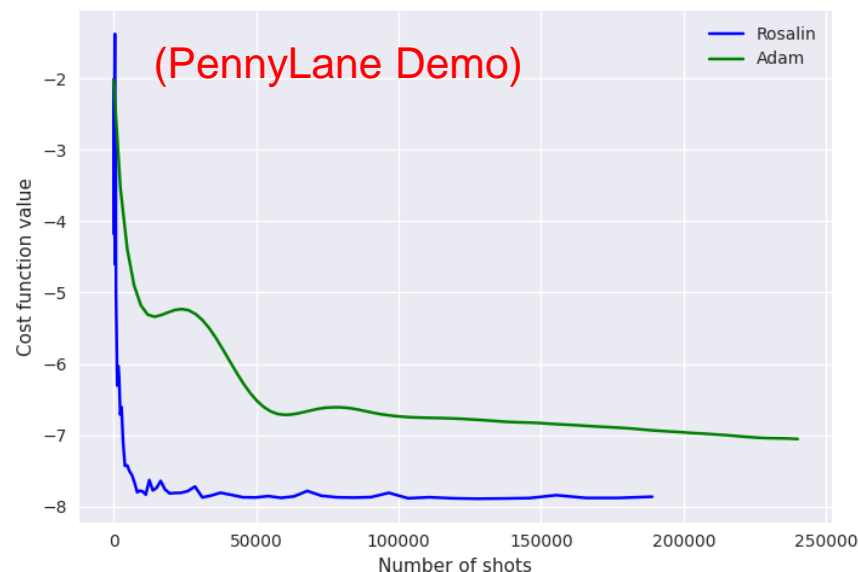
Classical VQE (2014): Full sampling. **Requires many measurements**

\* Arrasmith et. al. <https://arxiv.org/pdf/2004.06252.pdf>

\*\*<https://medium.com/@lana.bozanic/the-variational-quantum-eigensolver-efb8fab14c85>

[https://pennylane.ai/qml/demos/tutorial\\_rosalin.html](https://pennylane.ai/qml/demos/tutorial_rosalin.html)

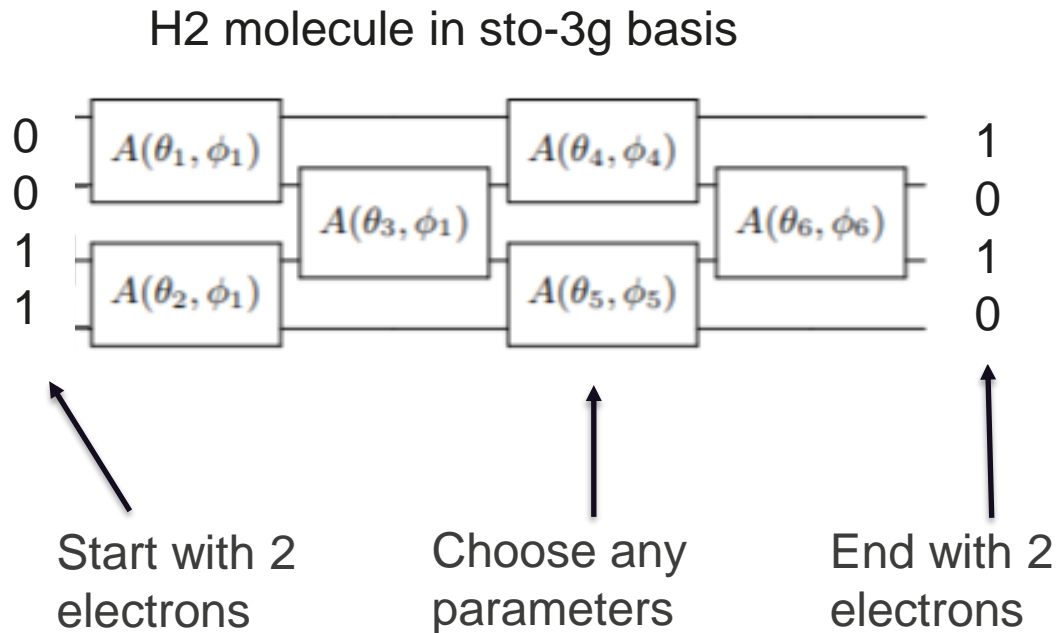
Rosalin Optimizer for VQE (Developed at LANL in this DR project)\*



ROSALIN optimizer: Random sampling from the Hamiltonian. **Reduces number of measurements.**

Implemented in Open-source PennyLane software package

# Creating a Variational Ansatz (Trial Wavefunction)



## Electron number preserving gates:

Ensures that our guess has the correct number of electrons. **Constrains search to subspace of  $2^4$  Hilbert Space**

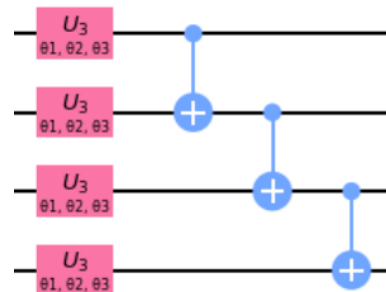
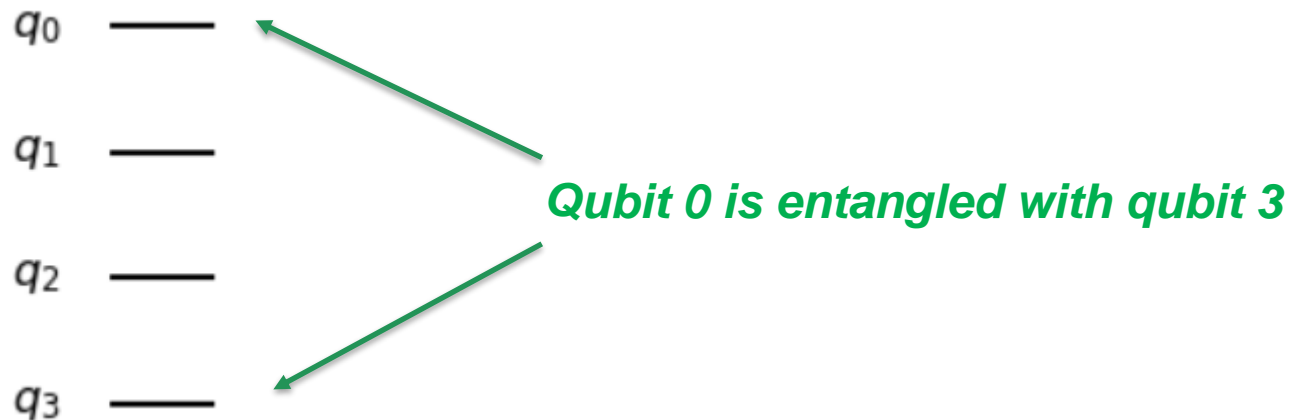
**Advantage:** Sweet spot between circuit depth and number of parameters.

Reduces measurement count from  $>10^8$  to  $10^5$  for Hydrogen molecule

<https://arxiv.org/pdf/1904.10910.pdf>

# Permutations of Hamiltonian

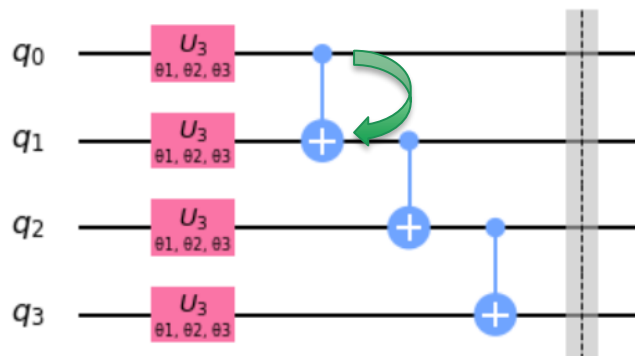
- **Idea:** to place more *entangled* qubits close to each other to reduce the required depth of the circuit and increase the accuracy of the final energy.



*Ansatz layer that corresponds to qubit chain architecture.*

# Permutations of Hamiltonian

- **Idea:** to place more *entangled* qubits close to each other to reduce the required depth of the circuit and increase the accuracy of the final energy.



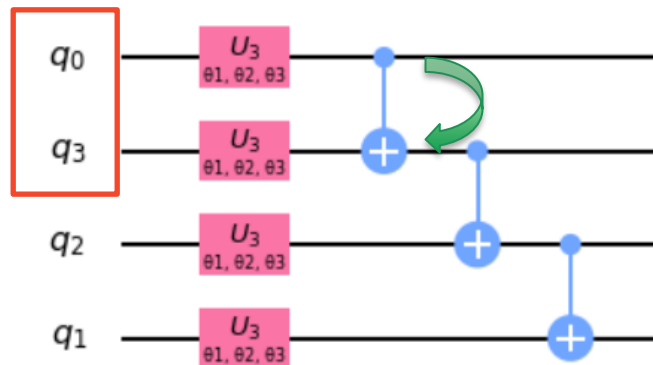
*3 layers are required to model the system*

*Let us swap qubit 1 and qubit 3*

# Permutations of Hamiltonian

- **Idea:** to place more *entangled* qubits close to each other to reduce the required depth of the circuit and increase the accuracy of the final energy.

*Now qubit 0 and 3 are close to each other*

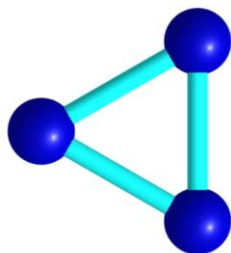


*After the permutation only 1 layer is required*



# Examples

- System:  $\text{H}_3^+$
- Basis: STO-3G
- Geometry:

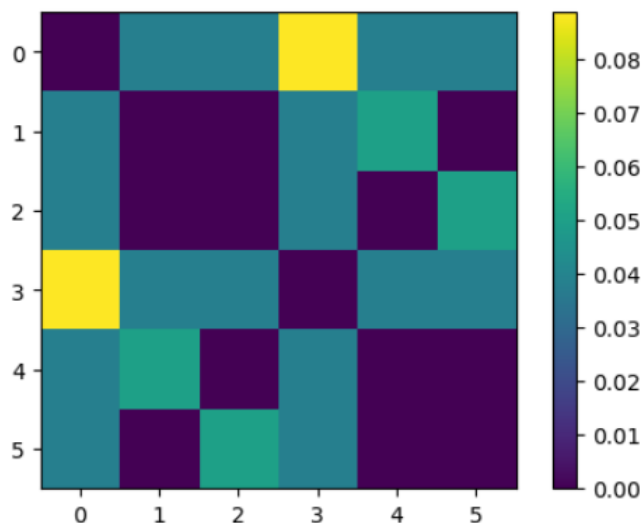


Number of qubits: 6

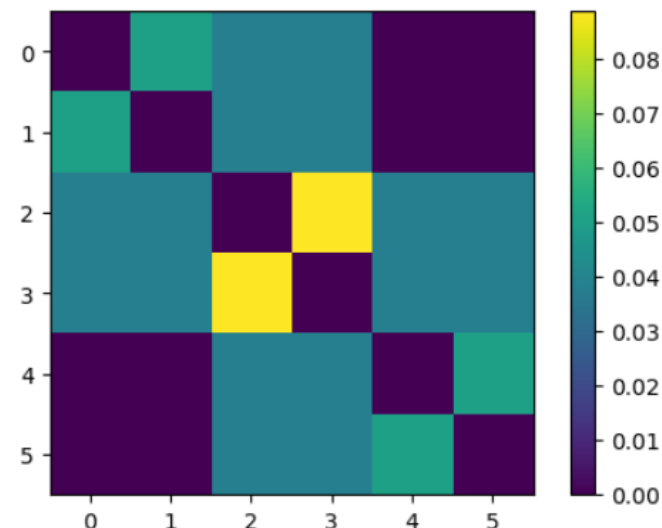
Exact energy: -1.35076

(need to match this energy by quantum computer)

## Before



## After



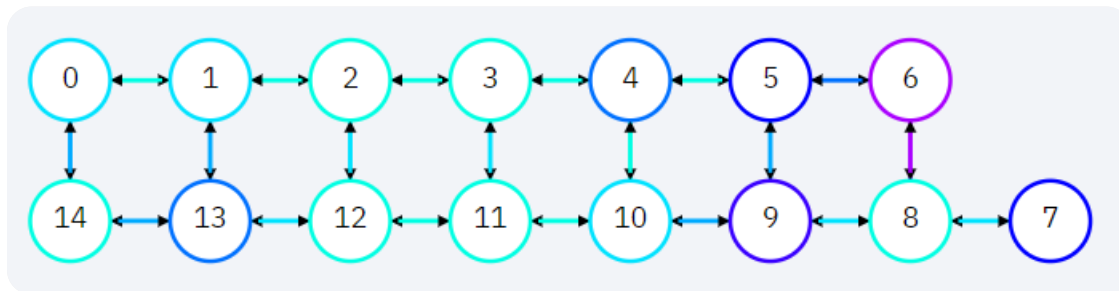
	Not permuted	Permuted
Final optimized energy for ansatz with a fixed depth = 5	-1.33158301 Hartree	<b>-1.35070986 Hartree</b>
Optimized energy - Exact energy	12 kcal/mol	<b>0.03 kcal/mol</b>

# Conclusions

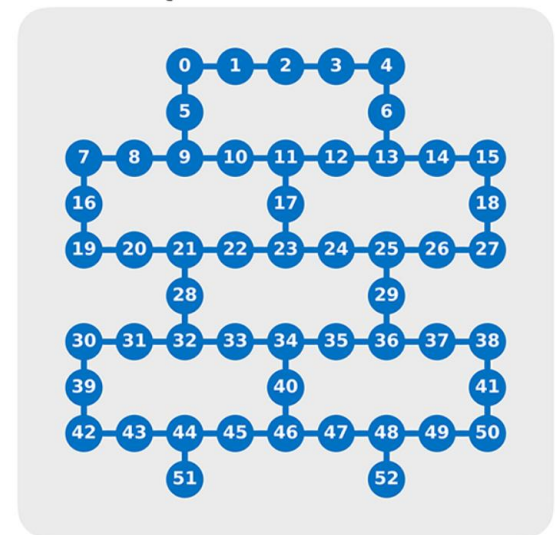
- U1 symmetric ansatz was encoded in LANL made Rosalin optimizer and reduces measurement counts by multiple orders of magnitude
- The algorithm to implement the permutations of Hamiltonian to ensure better entanglement between spin orbitals has been developed.
- Permuted Hamiltonian was found to be more efficient based on the several examples of molecules/encodings.

## Future plans

Take into account the topology of actual devices and run calculations on actual hardware (IBMQ, Rigetti).



53 Qubit Rochester Device



# Thank you for your attention!

## Acknowledgments

- LANL LDRD 20200056DR “Quantum Chemistry using Quantum Computers”
- LANL Quantum Summer School 2020

# Permutations of Hamiltonian

- **Idea:** to place more **entangled** qubits close to each other to reduce the required depth of the circuit and increase the accuracy of the final energy.

- **Entanglement measurement:**

Von Neumann entropy for the reduced system

$$S = - \sum_{\alpha} \omega_{\alpha} \ln(\omega_{\alpha})$$

$\omega_{\alpha}$  is the eigenvalue of the reduced density matrix

$S_i$  - single orbital entropy ( $\rho$  contains only one qubit)

$S_{ij}$  - two - orbital entropy ( $\rho$  contains two qubits)

**Mutual information:**  $I_{ij} = (S_i + S_j - S_{ij})(1 - \delta_{ij})$

# Fiedler Vector

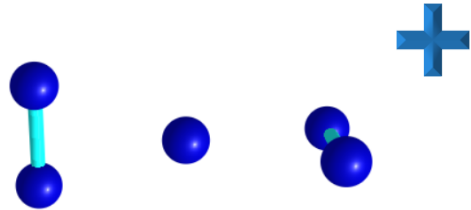
**Cost function:**  $f = \sum_{i < j} I_{ij} (i - j)^2$

We know that Fiedler vector can minimize the following function:

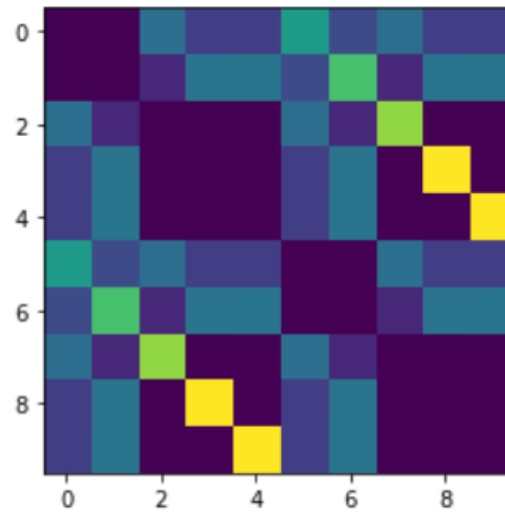
$$F(\mathbf{x}) = \mathbf{x}^T L \mathbf{x} = \sum_{i < j} I_{ij} |x_i - x_j|^2$$

Where  $L$  is a graph Laplacian:  $\mathbf{L} = \mathbf{D} - \mathbf{I}$ , with  $\mathbf{D}_{ii} = \sum_j I_{ij}$

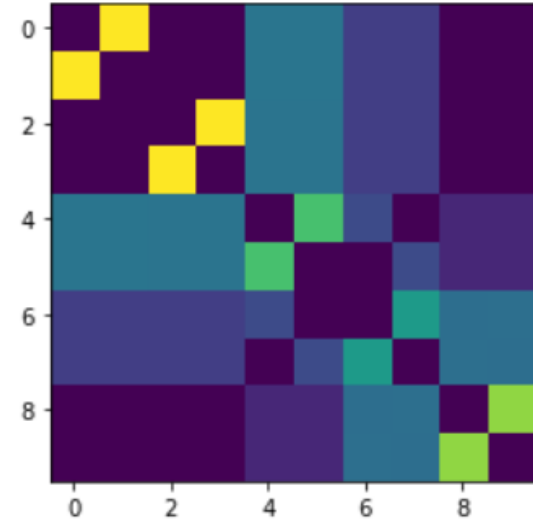
# Examples



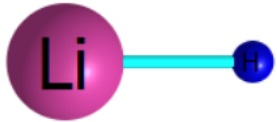
Before



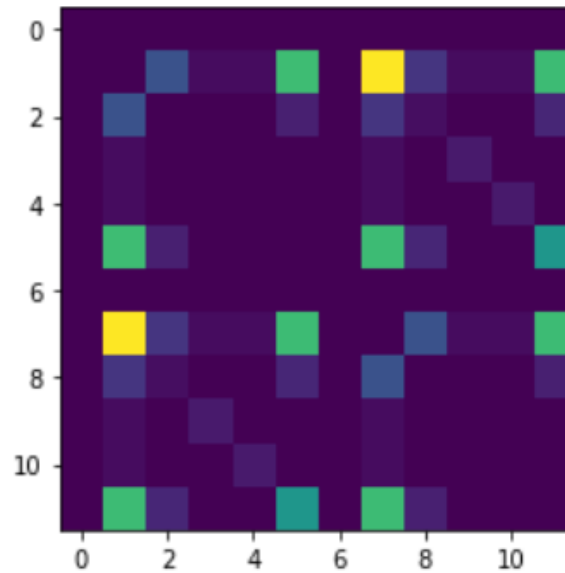
After



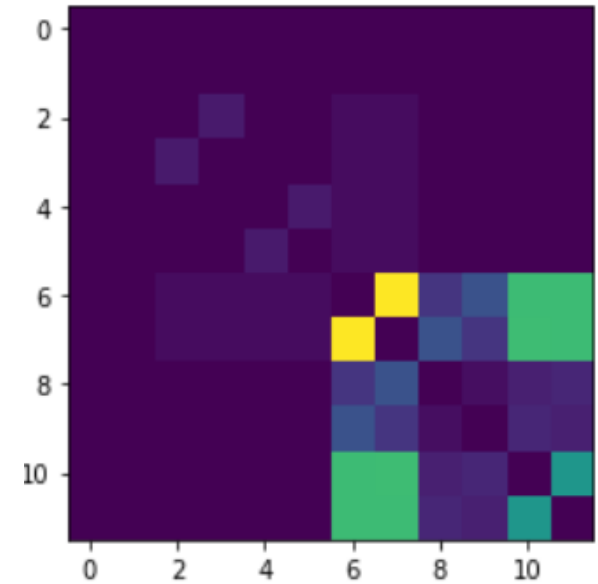
# Examples



Before

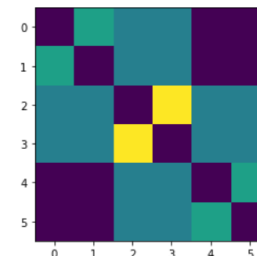
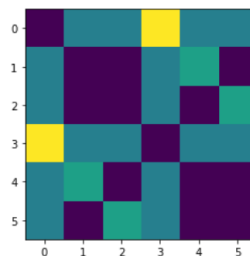


After



# Preliminary Results

- **System:**  $\text{H}_3^+$
- **Optimizer:** COBYLA
- **Exact energy:** -1.3507306 (need to match this energy by quantum computer)
- **Ansatz depth:** 5



	Not permuted	Permuted
Energy	-1.33158301	<b>-1.35070986</b>
Num_evaluations	100000	100000





# **Nonadiabatic molecular dynamics analysis of Hybrid Dion-Jacobson 2D Lead Iodide Perovskites**

Ying Wang

Mentor: Dr. Amanda Neukirch

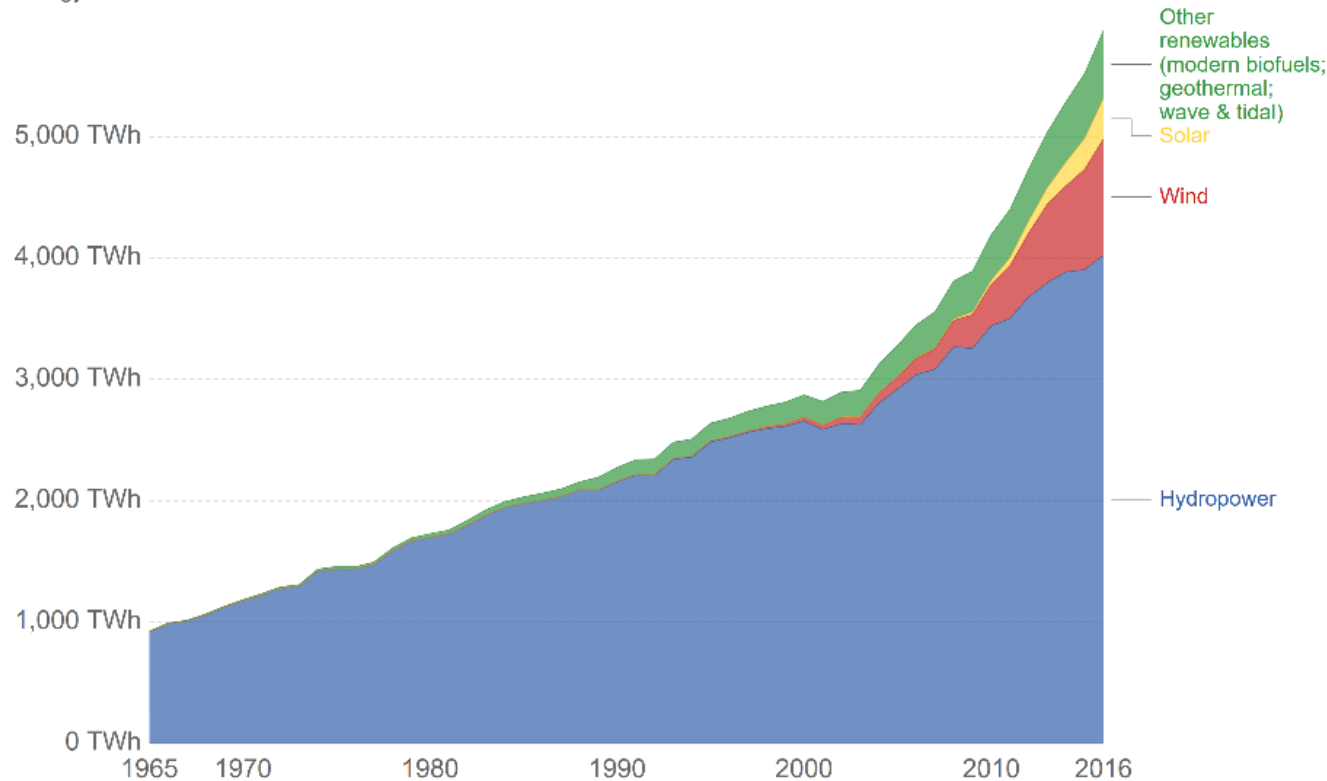
Co-mentor: Dr. Dibyajyoti Ghosh

08/13/2020

# Introduction to Solar Cells

## Modern renewable energy consumption, World

Total renewable energy consumption, measured in terawatt-hours (TWh) per year. This data includes all renewable energy sources with the exclusion of traditional biomass.



Source: BP Statistical Review of Global Energy

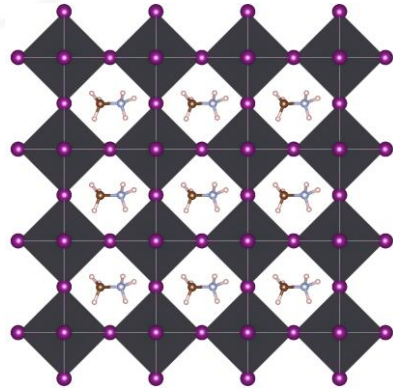
CC BY



crystalline silicon:

- Manufacturing is relatively expensive
- Improvement of efficiency is almost saturated.

# Hybrid Organic-Inorganic Lead Halide Perovskites: Promising Materials as Solar Cells



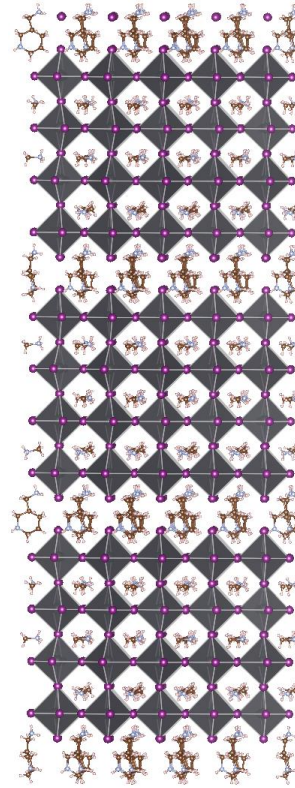
3D hybrid organic-inorganic  
lead halide perovskites

## Pros:

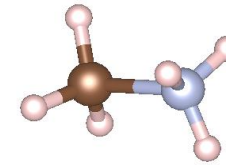
- **Solution-processed, low cost**
- **Direct bandgap**
- **Strong light absorption**
- **High carrier mobility**

## Cons:

- **Instability against moisture, light and heat**

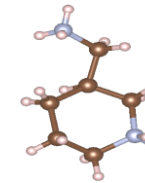


hybrid Dion-Jacobson 2D  
lead iodide perovskites



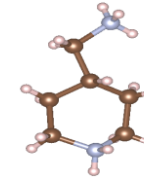
methylammonium

## Interlayer cations



3AMP

3-(aminomethyl)piperidinium

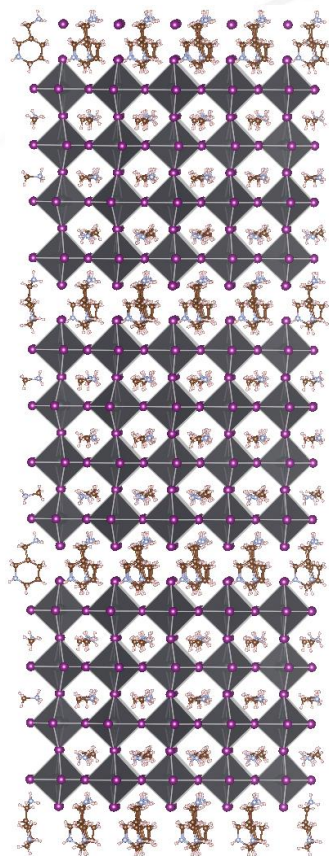


4AMP

4-(aminomethyl)piperidinium

- **High stability**
- **Low but promising power conversion efficiency (PCE)**
- **Highly tunable structure**

# Hybrid Organic-Inorganic Lead Halide Perovskites: Promising Materials as Solar Cells



hybrid Dion-Jacobson 2D  
lead halide perovskites

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JOURNAL OF THE AMERICAN CHEMICAL SOCIETY

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Article  
pubs.acs.org/JACS

## Hybrid Dion–Jacobson 2D Lead Iodide Perovskites

Lingling Mao,<sup>†</sup> Weijun Ke,<sup>†</sup> Laurent Pedesseau,<sup>§</sup> Yilei Wu,<sup>†</sup> Claudine Katan,<sup>||</sup> Jacky Even,<sup>§</sup>  
Michael R. Wasielewski,<sup>†,‡</sup> Constantinos C. Stoumpos,<sup>\*,†</sup> and Mercuri G. Kanatzidis<sup>\*,†,‡</sup>

<sup>†</sup>Department of Chemistry, and <sup>‡</sup>Argonne-Northwestern Solar Energy Research (ANSER) Center, Northwestern University, Evanston, Illinois 60208, United States

<sup>§</sup>Univ Rennes, INSA Rennes, CNRS, Institut FOTON – UMR 6082, Rennes F-35000, France

<sup>||</sup>Univ Rennes, ENSCR, INSA Rennes, CNRS, ISCR (Institut des Sciences Chimiques de Rennes) – UMR 6226, Rennes F-35000, France

- $(3\text{AMP})(\text{MA})_3\text{Pb}_4\text{I}_{13}$  with the narrowest bandgap and least distorted structure, longer lifetime
- 3AMP-PCE ~7%, 4AMP-PCE ~5%

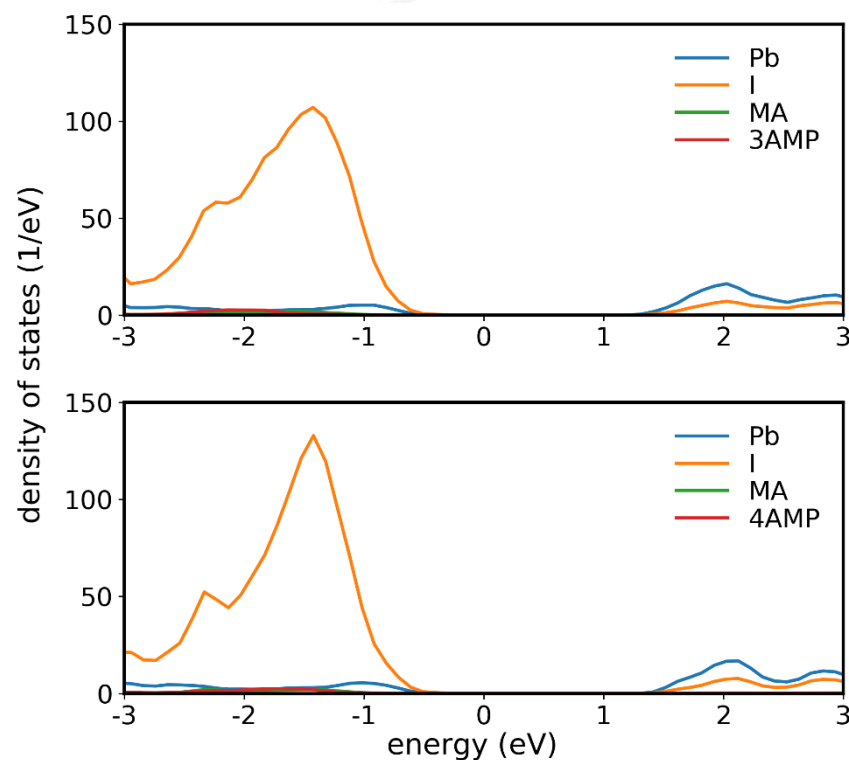
Prediction:

$(3\text{AMP})(\text{MA})_3\text{Pb}_4\text{I}_{13}$  has obviously shorter nonradiative lifetime than  $(4\text{AMP})(\text{MA})_3\text{Pb}_4\text{I}_{13}$

Rationalize the reason of the lifetime difference



# Main composition of CBM & VBM



VBM

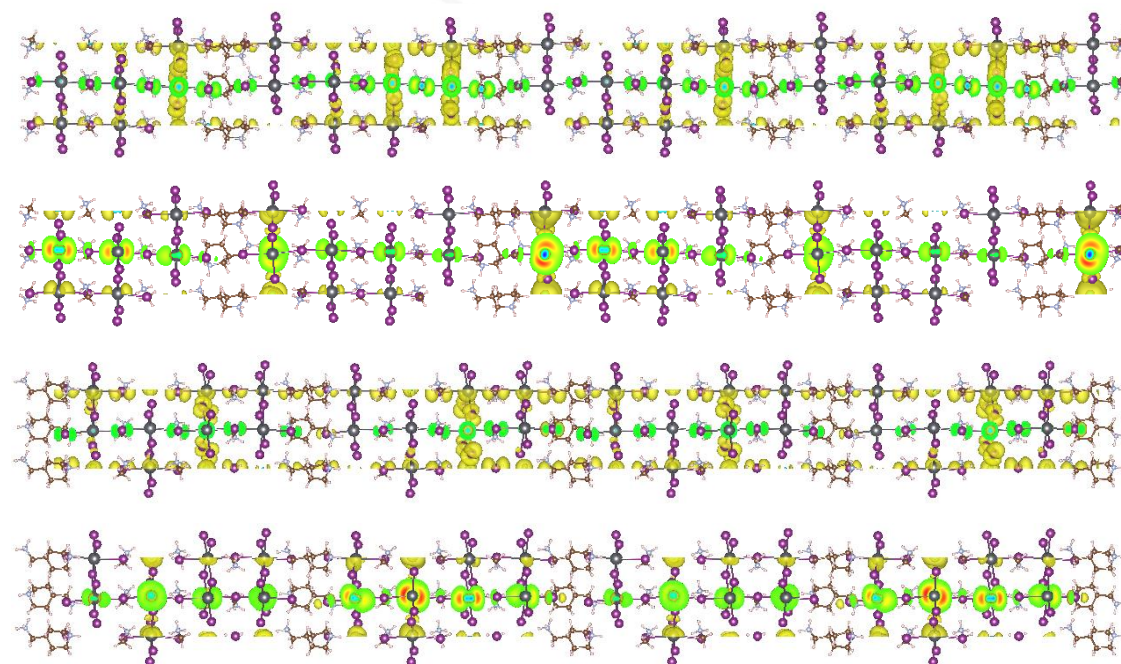
3AMP

CBM

VBM

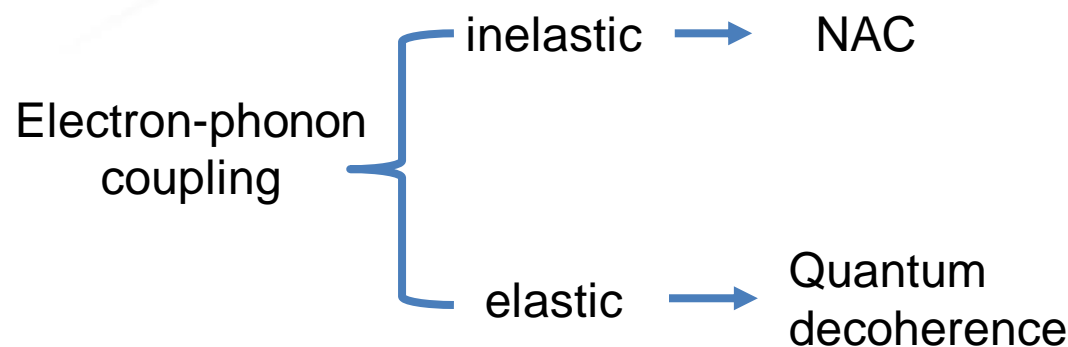
4AMP

CBM



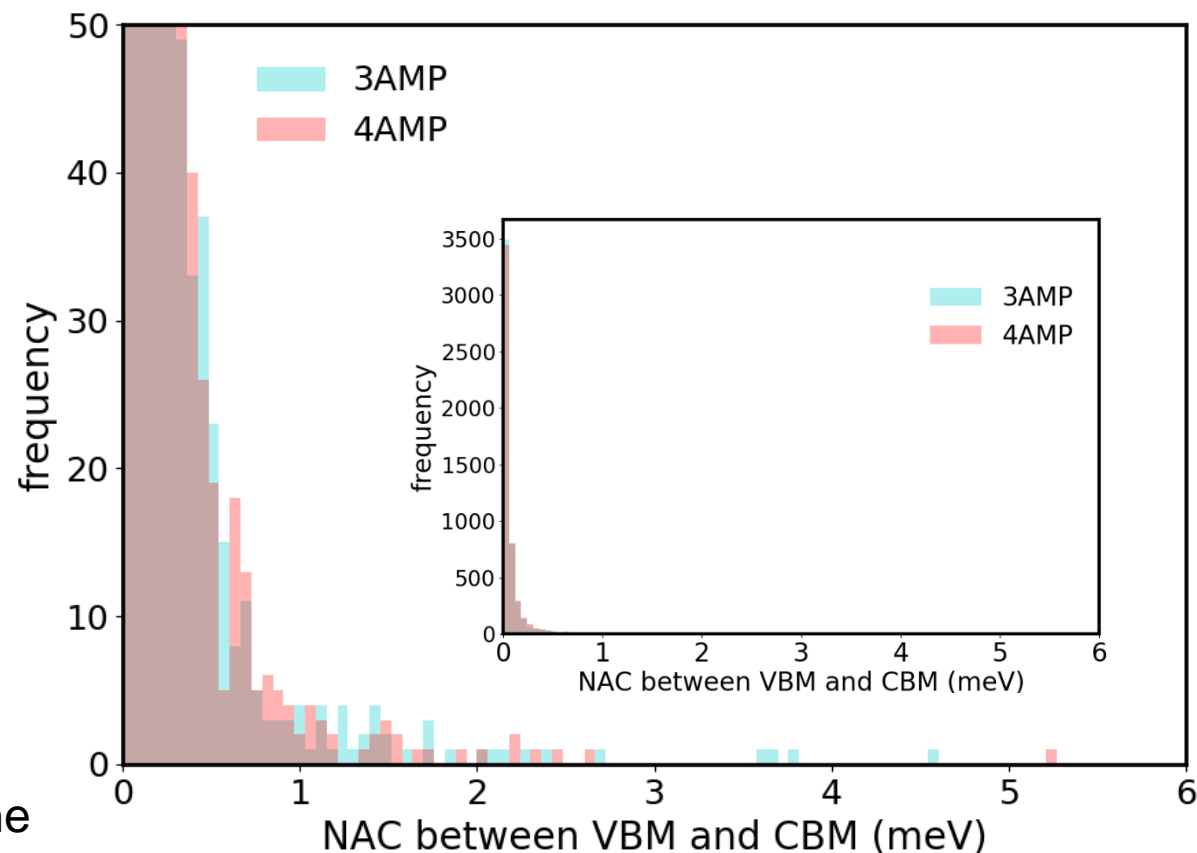
I-5p and Pb-6s dominate VBM (valence band maximum)  
Pb-6p and I-5p dominate CBM (conduction band minimum)

# Nonadiabatic coupling (NAC) doesn't play an important role in determining nonradiative lifetime



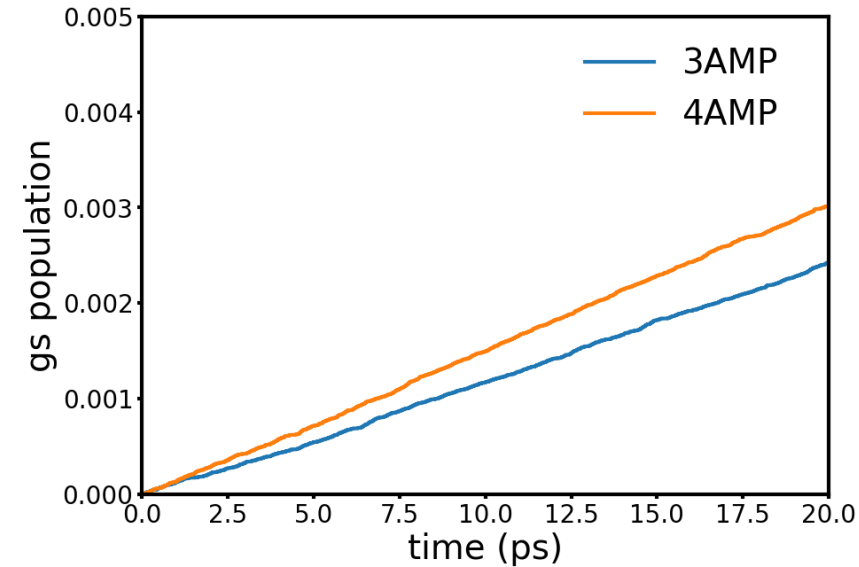
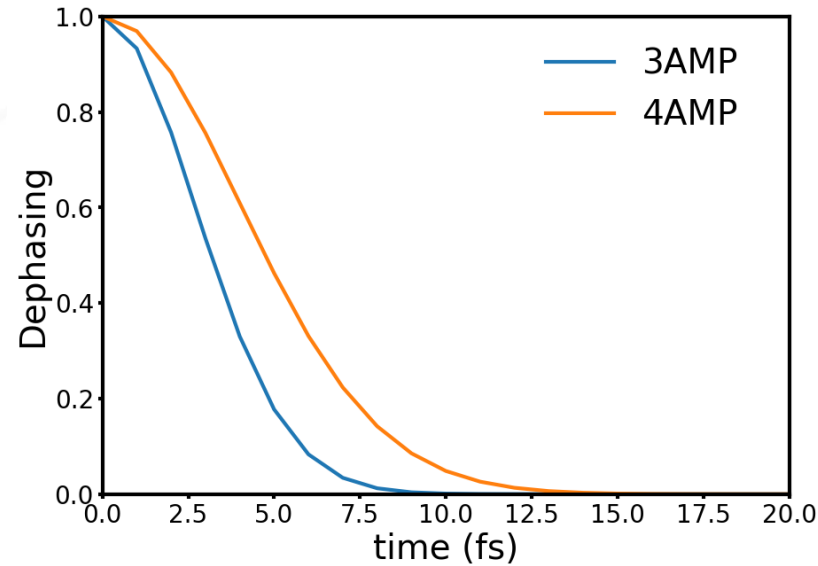
the stronger the NAC, the faster the nonradiative carrier recombination is

When decoherence time is short enough, the shorter decoherence time induces the slower nonradiative carrier recombination



Areas when NAC > 0.1 meV of 3AMP 4AMP is almost the same

# Dephasing time decides the nonradiative lifetime difference



species	Bandgap (eV, 0K)	Bandgap (eV, 300K)	NAC (meV, rms)	Dephasing time (fs)	Radiative lifetime ( $\mu$ s)	Nonradiative lifetime (ns)
3AMP	1.82	1.40	0.21	2.69	1.03	4.12
4AMP	1.89	1.52	0.19	4.03	1.14	3.32



# Conclusion

- Nonadiabatic coupling (NAC) doesn't play an important role in determining nonradiative lifetime
- The shorter dephasing time leads to the longer nonradiative lifetime of 3AMP, better solar cell performance

## Future:

- Relation between the structures and optoelectronic properties.
- Analyzing the vibrational modes which facilitate electron-phonon coupling



Thank you for your attention

Q & A

# References

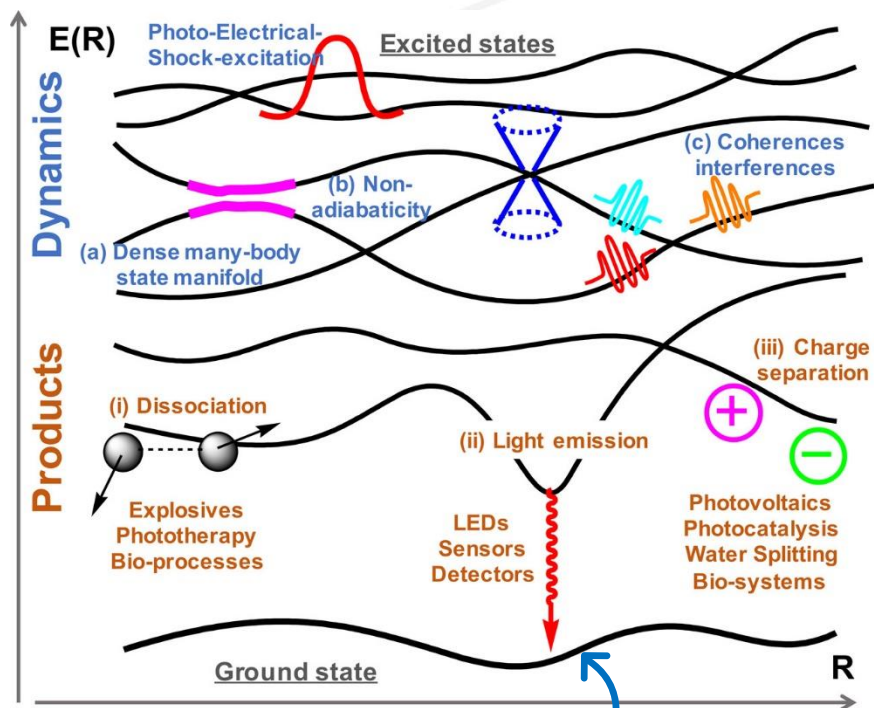
- [1] Mao, Lingling, et al. "Hybrid Dion–Jacobson 2D lead iodide perovskites." *Journal of the American Chemical Society* 140.10 (2018): 3775-3783.
- [2] Pedesseau, Laurent, et al. "Advances and promises of layered halide hybrid perovskite semiconductors." *ACS nano* 10.11 (2016): 9776-9786.
- [3] Yan, Jieli, et al. "Recent progress in 2D/quasi-2D layered metal halide perovskites for solar cells." *Journal of Materials Chemistry A* 6.24 (2018): 11063-11077.
- [4] Zhang, Zhaosheng, et al. "Rapid decoherence suppresses charge recombination in multi-layer 2D halide perovskites: Time-domain ab initio analysis." *Nano letters* 18.4 (2018): 2459-2466.
- [5] Zhou, Liujiang, et al. "Density of states broadening in CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> hybrid perovskites understood from ab initio molecular dynamics simulations." *ACS Energy Letters* 3.4 (2018): 787-793.
- [6] Filip, Marina R., Carla Verdi, and Feliciano Giustino. "GW band structures and carrier effective masses of CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> and hypothetical perovskites of the type APbI<sub>3</sub>: A= NH<sub>4</sub>, PH<sub>4</sub>, AsH<sub>4</sub>, and SbH<sub>4</sub>." *The Journal of Physical Chemistry C* 119.45 (2015): 25209-25219.



# Appendix

# Methodology-nonadiabatic molecular dynamics (NAMD)

## Mixed quantum-classical method



Nuclei are treated classically  
Molecular dynamics (MD) trajectory  
is predetermined and independent  
of electronic evolution

Multielectron adiabatic  
representation

$$\Psi(r, R, t) = \sum_{i=0}^{N_b-1} C_i(t) \Phi_i(r|R)$$

Time-dependent Schrodinger eq

$$i\hbar \frac{\partial \Psi(r, R, t)}{\partial t} = H(r, R, t) \Psi(r, R, t)$$

Nonadiabatic coupling (NAC)

$$i\hbar \frac{dc_i}{dt} = \sum_{j=0}^{N_b-1} (\epsilon_i \delta_{ij} - i\hbar \vec{d}_{ij}) C_j$$

$$\vec{d}_{ij} = \left\langle \Phi_i \left| \frac{\partial \Phi_j}{\partial R} \right. \right\rangle, -i\hbar \vec{d}_{ij} \equiv \text{NAC}$$

Decoherence-induced surface hopping  
(DISH)

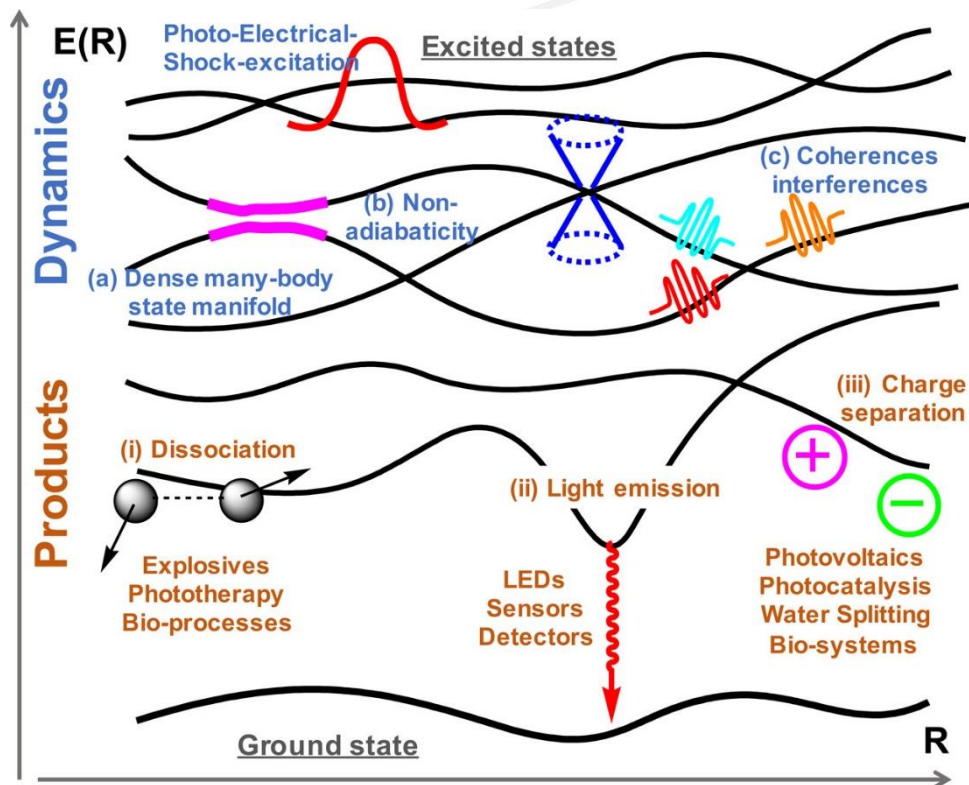
$t > \tau_\alpha$ , the system **either hops to that state with probability  $|c_\alpha|^2$  or the state is projected out with probability  $1 - |c_\alpha|^2$**

$\tau_\alpha$  -- decoherence time

$$g_{i \rightarrow j} = p_{i \rightarrow j} b_{i \rightarrow j}, b_{i \rightarrow j} = \begin{cases} \exp\left(-\frac{\epsilon_j - \epsilon_i}{k_B T}\right) & \epsilon_j > \epsilon_i \\ 1 & \epsilon_j \leq \epsilon_i \end{cases}$$

py **x**aiD

Chem. Rev. 2020, 120, 4, 2215-2287



When decoherence time is shorter than the time of coherent electronic transition—including decoherence in surface hopping

$$C_{un}(t) = \langle \delta E(t) \delta E(0) \rangle_T$$

Decoherence function

$$D(t) = \exp\left(-\frac{1}{\hbar^2} \int_0^t dt' \int_0^{t'} dt'' C(t'')\right)$$

Gaussian fitting

Decoherence function characterizes the dephasing time.





# Long-timescale simulations of morphological changes in $\text{ThO}_2$ nanoparticles

Shanshan Wu

08/13/2020

Mentors: Danny Perez & Ping Yang



# Motivation:

## Properties:

Thorium is a naturally radioactive element

Thorium has abundant reserves (3-4 times of U)

ThO<sub>2</sub> is the highest melting point oxide (3390°C)



## Applications:

Nuclear reactor fuel pellets

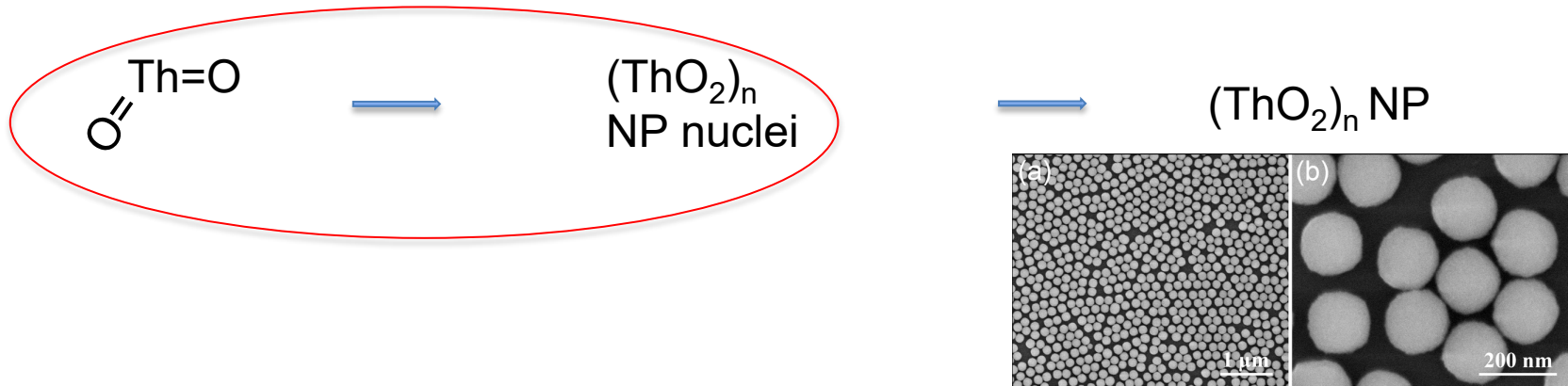
Catalysts, optical glass

Sinter fine high-temperature ceramics

<div>89 <b>Ac</b> (227)</div> <div>ACTINIDES</div>													
Thorium	Protactinium	Uranium	Neptunium	Plutonium	Americium	Curium	Berkelium	Californium	Einsteinium	Fermium	Mendelevium	Nobelium	Lawrencium
90 <b>Th</b> (232)	91 <b>Pa</b> (231)	92 <b>U</b> (238)	93 <b>Np</b> (237)	94 <b>Pu</b> (244)	95 <b>Am</b> (243)	96 <b>Cm</b> (247)	97 <b>Bk</b> (247)	98 <b>Cf</b> (251)	99 <b>Es</b> (252)	100 <b>Fm</b> (257)	101 <b>Md</b> (258)	102 <b>No</b> (259)	103 <b>Lr</b> (260)

# Goal

- Understand the dynamical and morphological properties of ultrasmall  $\text{ThO}_2$  nanoparticles (NPs)
- Using accelerated molecular dynamics (AMD) simulations in actinide systems

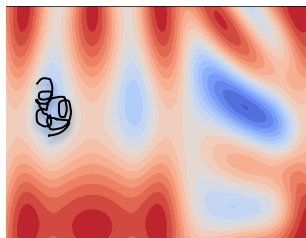
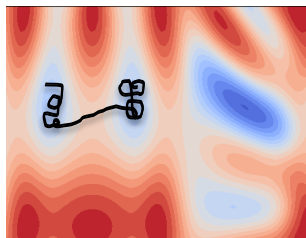
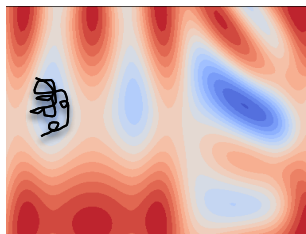
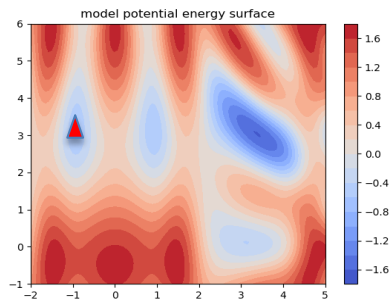


L. Wang, R. Zhao, X. Wang, L. Mei, L. Yuan, S. Wang, Z. Chai, W. Shi.  
*CrystEngComm* 16, no. 45 (2014): 10469-10475.

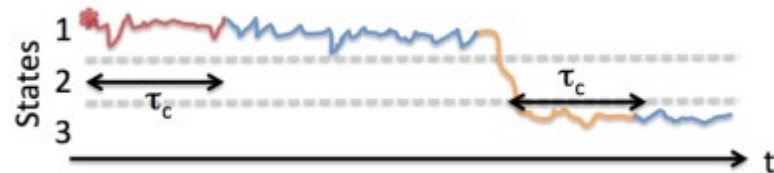
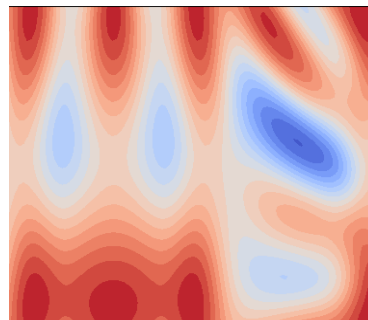
# Parallel trajectory splicing (ParSplice) – An accelerated MD method

- MD simulation: numerical propagation of Newton's equation  $m\ddot{x} = f = -\nabla U$ 
  - Provides fully-resolved atomistic trajectories
  - But the timescale of traditional MD is limited to ns
- ParSplice: parallelize MD simulation in the time domain to extend timescales to  $\mu\text{s}$ , ms

# ParSplice



Segment database



Spliceable segments  
 $T_c$ : decorrelation time

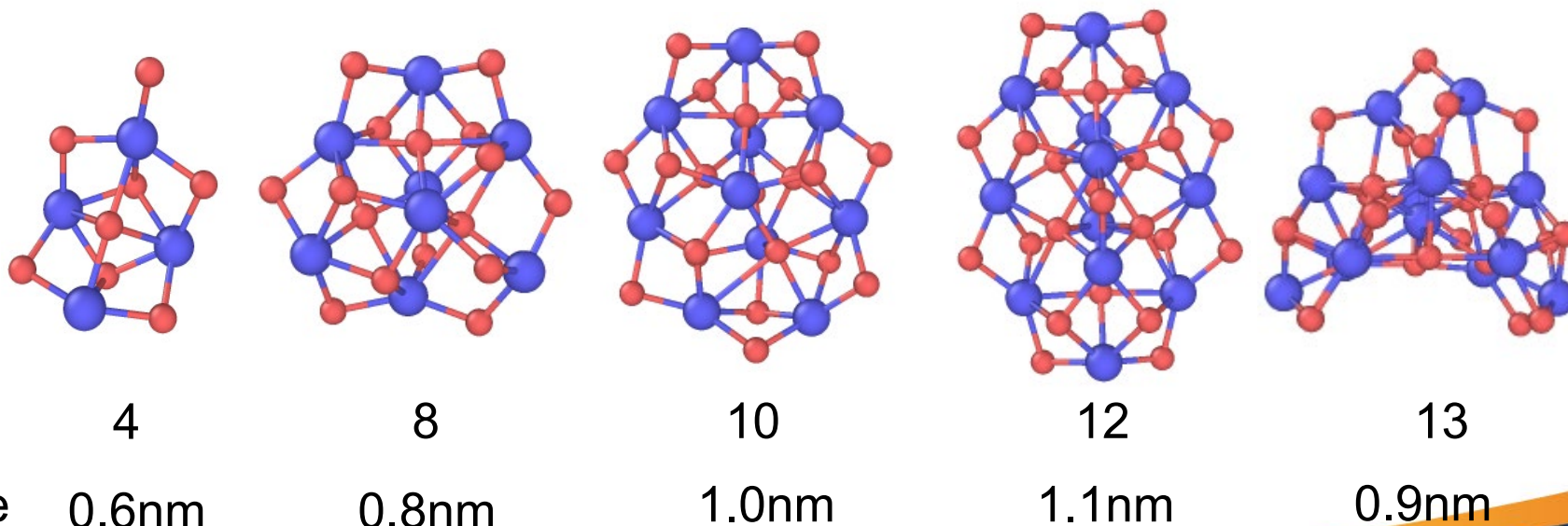
- Lost strict time continuity
- Generate statistically correct state-to-state trajectories
- Use HPC efficiently

D Perez, ED Cubuk, A Waterland, J. Chem Theory Comput 12.1 (2016): 18-28.

# AMD simulation of $(\text{ThO}_2)_n$ in vacuum

Force field: EAM/alloy

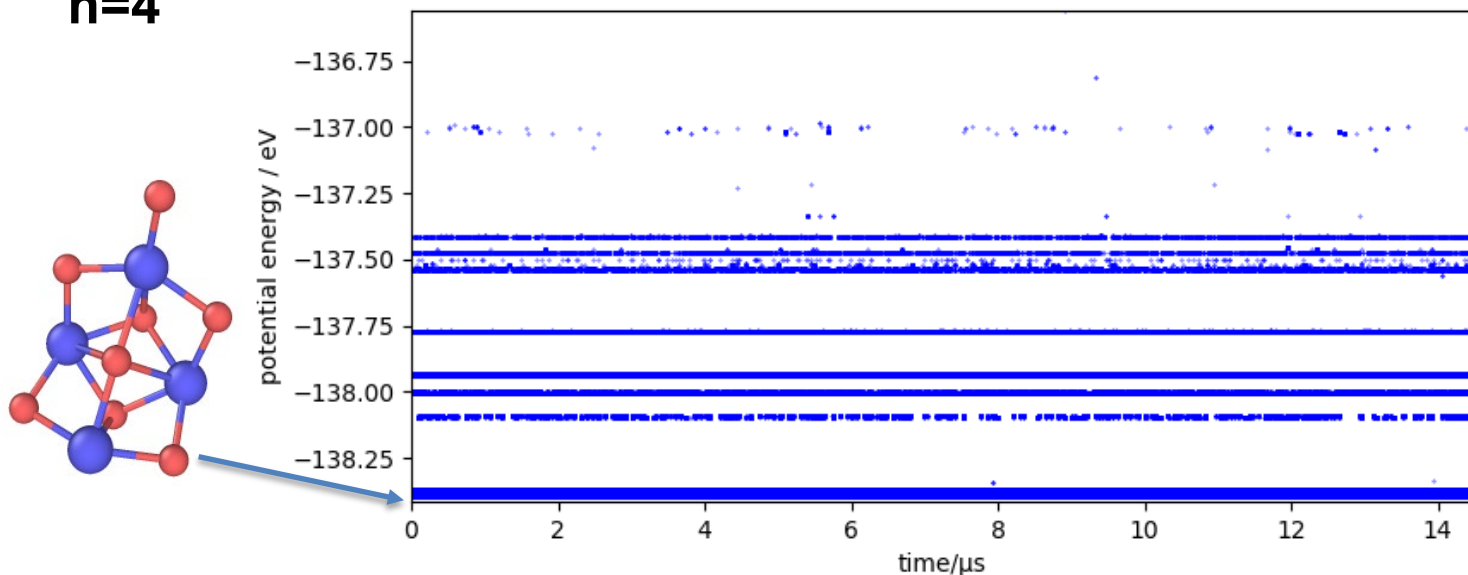
Systems being studied:



M.W.D. Cooper, M.J.D Rushton, R.W. Grimes, *J. Phys. Condens. Matter* **26** (2014) 105401.

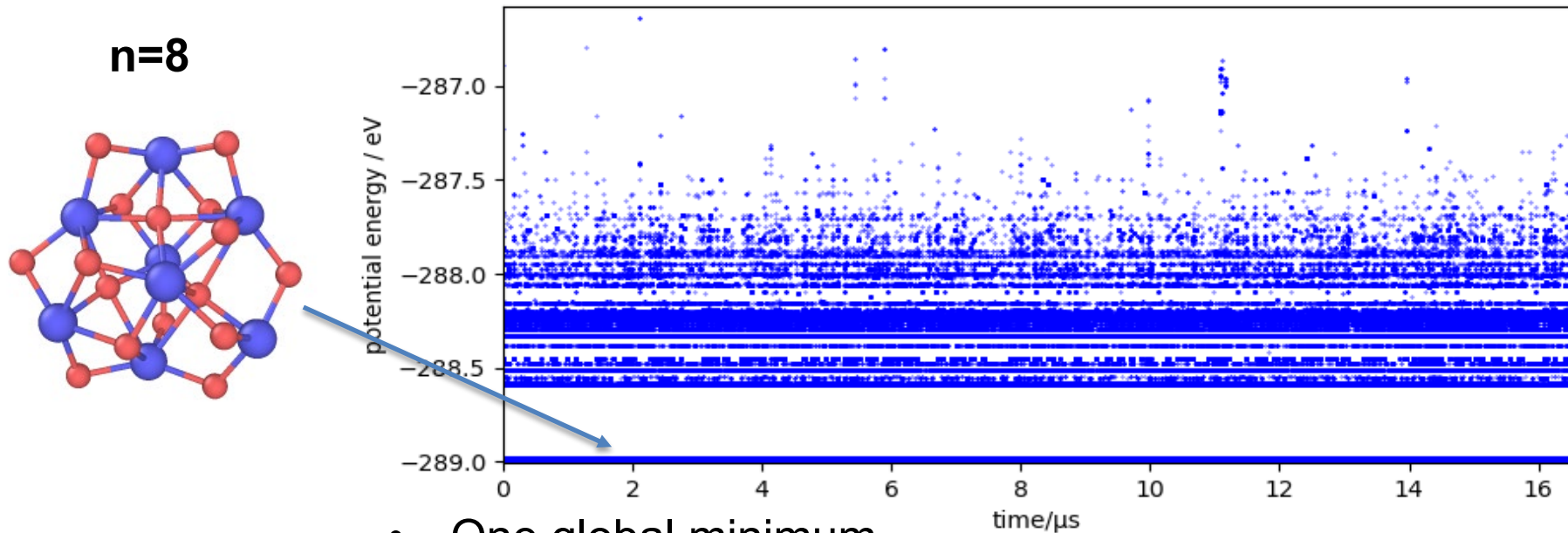
# AMD simulation of $(\text{ThO}_2)_n$ in vacuum at 1000K

$n=4$



- One global minimum
- Frequent transitions among local minima

# AMD simulation of $(\text{ThO}_2)_n$ in vacuum at 1000K

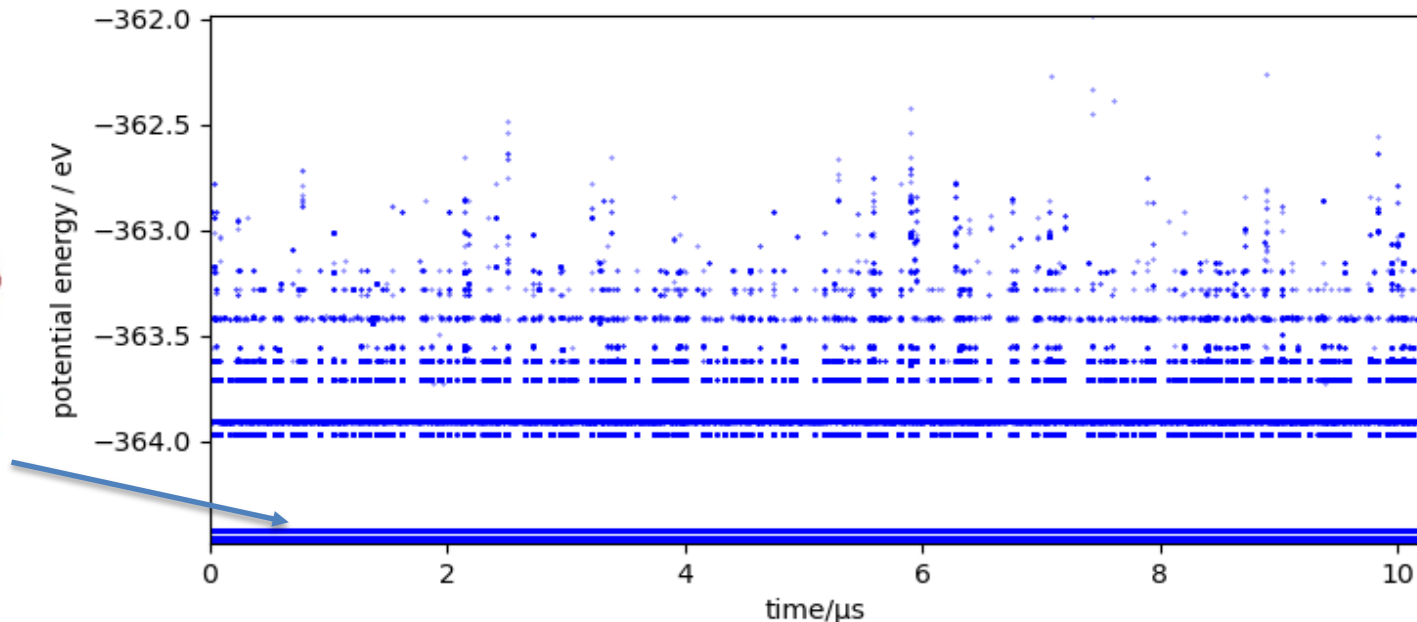
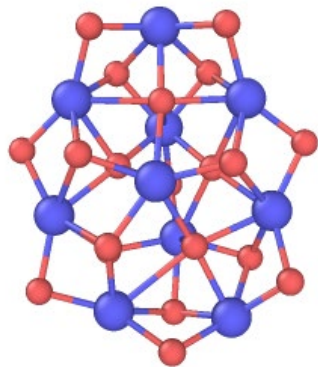


- One global minimum
- Frequent transitions among local minima



# AMD simulation of $(\text{ThO}_2)_n$ in vacuum at 1000K

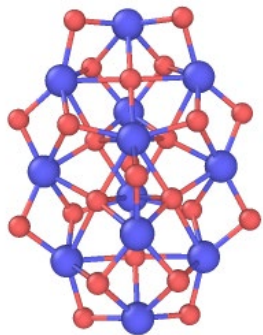
$n=10$



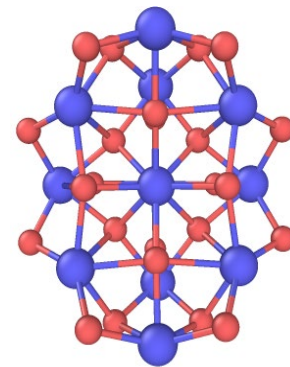
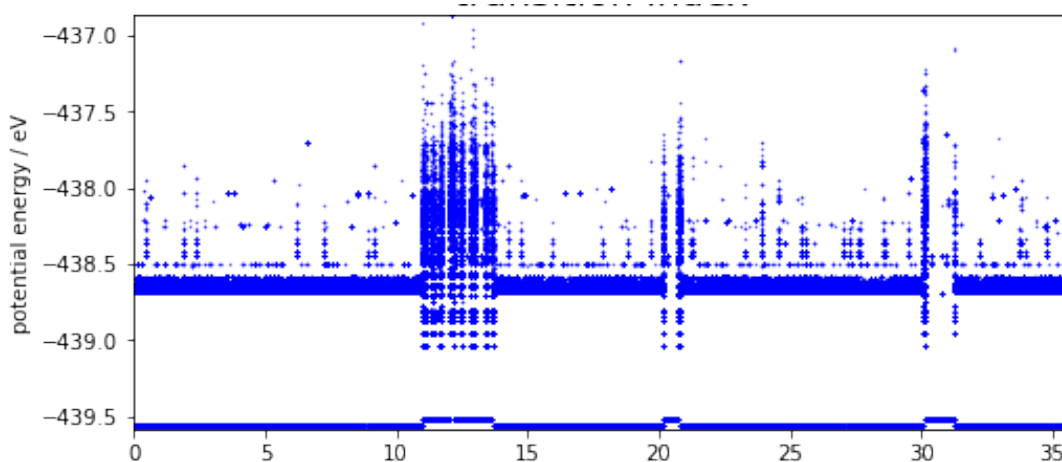
- One superbasin with two very close minima
- Frequent transitions among local minima

# AMD simulation of $(\text{ThO}_2)_n$ in vacuum at 1000K

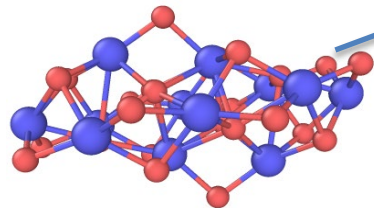
$n=12$



Front view



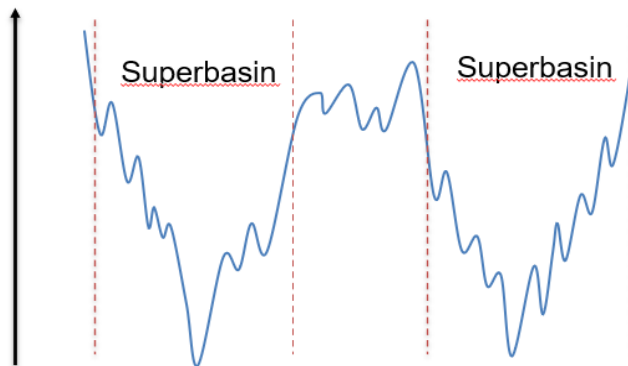
Front view



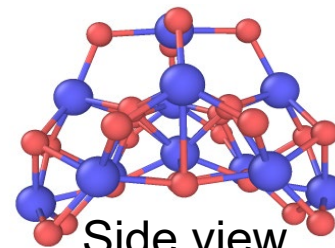
Side view

- T
- A
- S

Potential energy



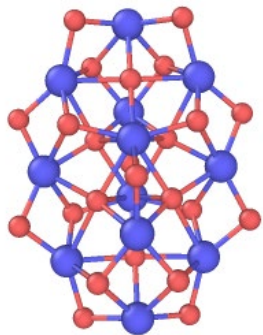
cale  
e energy



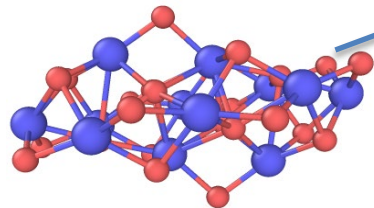
Side view

# AMD simulation of $(\text{ThO}_2)_n$ in vacuum at 1000K

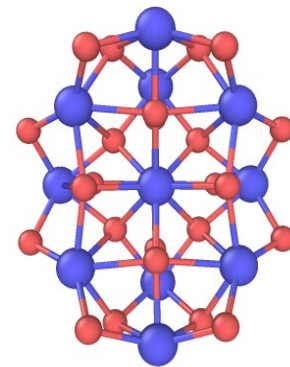
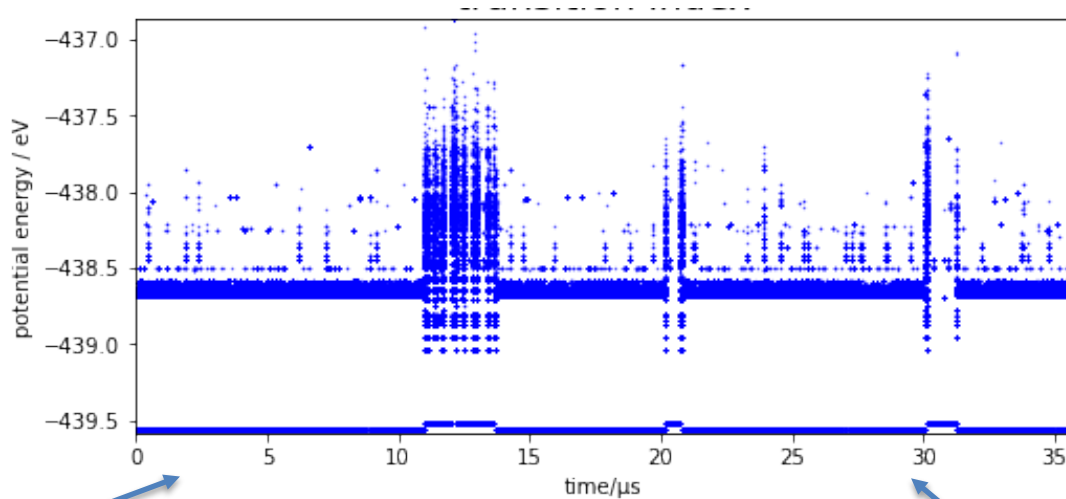
$n=12$



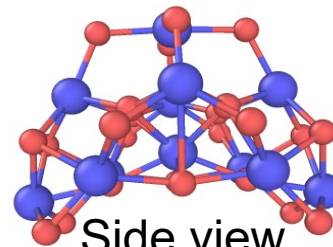
Front view



Side view



Front view

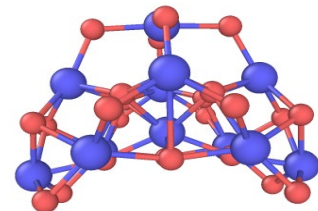
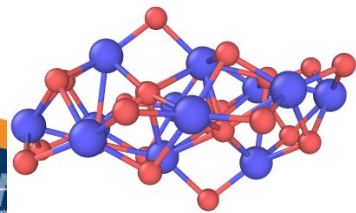
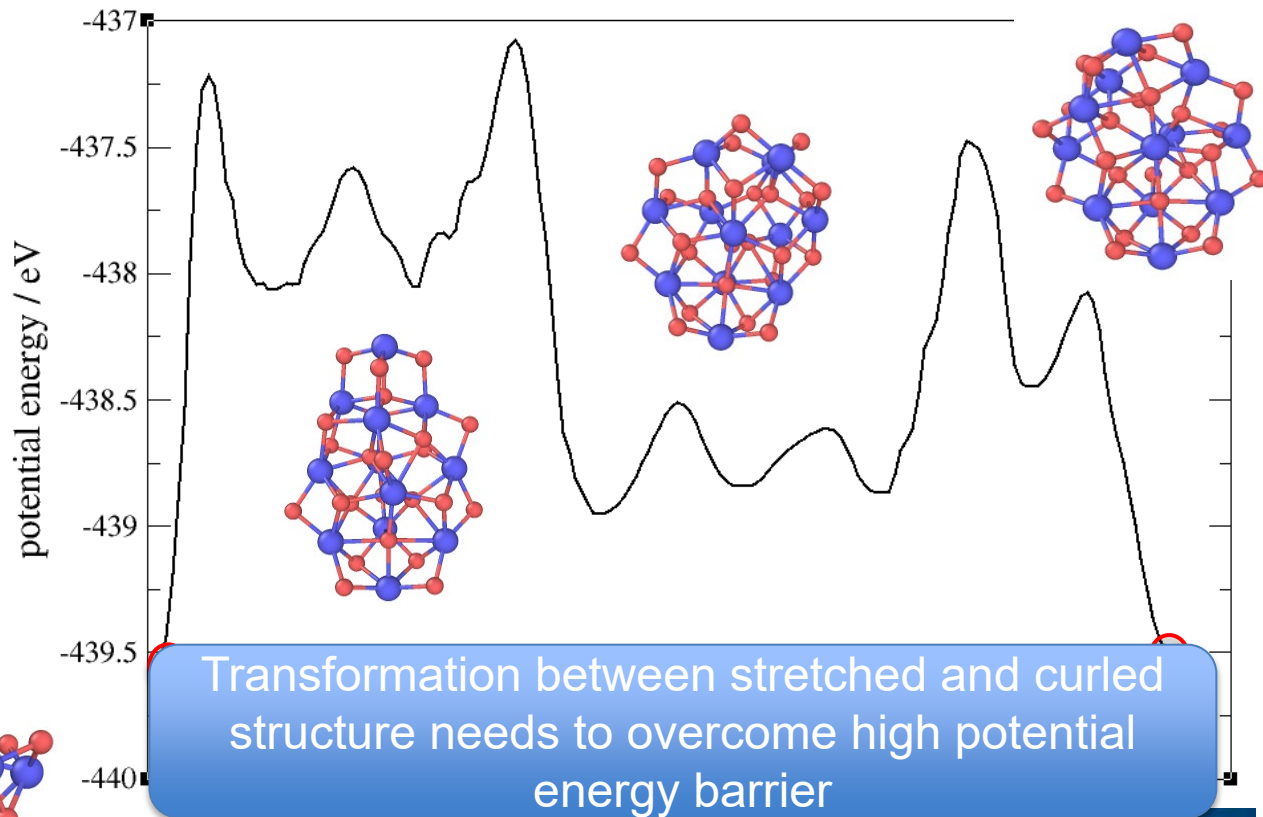


Side view

- Two superbasins were visited
- Activated process occur on  $\mu\text{s}$  timescale
- Stretched structure has minimum free energy

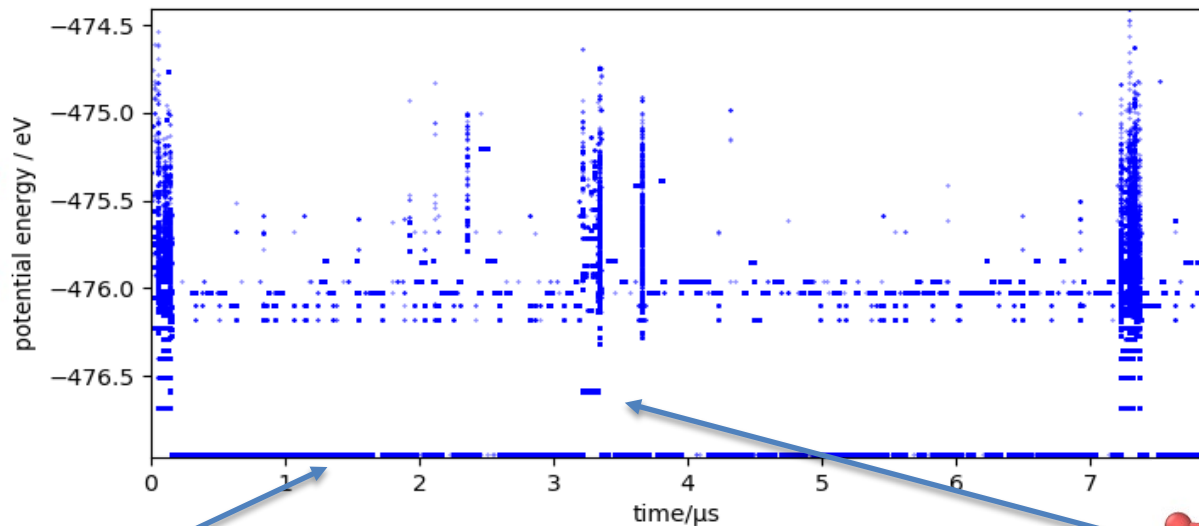
# Minimum energy pathway of n=12 transition

2.3eV PE  
barrier



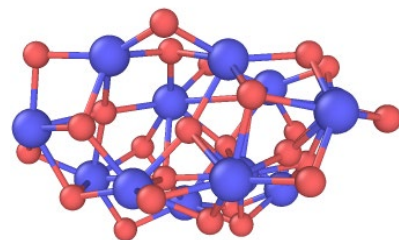
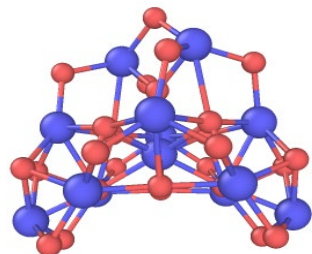
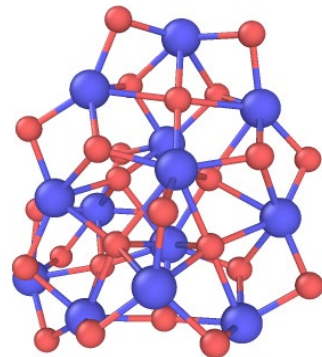
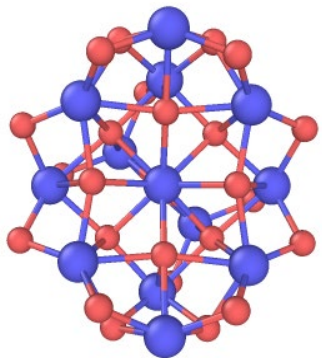
# AMD simulation of $(\text{ThO}_2)_n$ in vacuum at 1000K

$n=13$



Front view

Front view



Side view

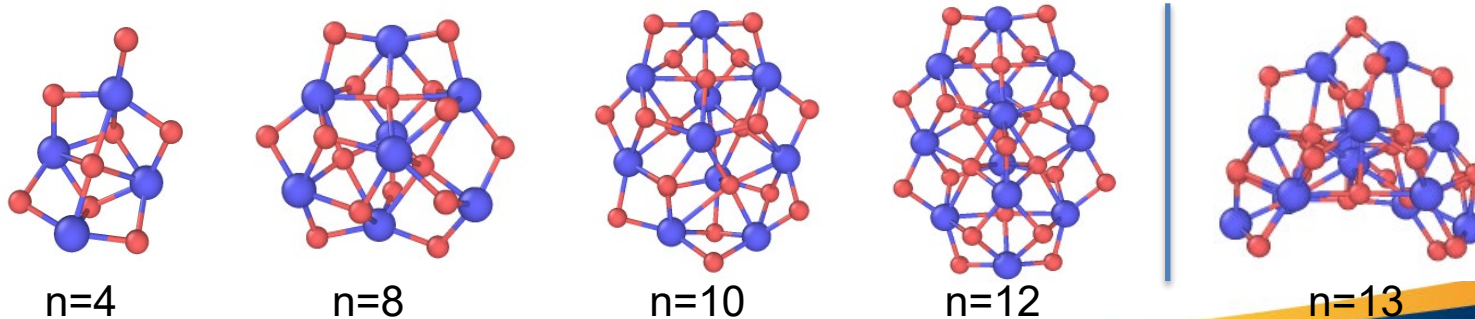
Side view

Two superbins were visited  
Curled structure has minimum free energy



# Conclusion

- This is the first  $\text{ThO}_2$  NP system to be studied by the long-timescale dynamics.
- $\text{ThO}_2$  NPs  $\leq 10$  units have one superbasin. Local minima will relax to global minimum quickly. At lower temperature, one global conformation may be observed
- $\text{ThO}_2$  NPs with 12 and 13 units have two superbasins. Two types of configurations may be observed. The slow time scale of transition from stretched to curled conformation may affect the growth of small NPs.



# Future work

- Accelerated Quantum MD simulations of  $\text{ThO}_2$  NPs
- AMD simulation of larger size clusters to understand the nucleation process



# Acknowledgments:

## Mentors:

Dr. Danny Perez

Dr. Ping Yang

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DOE-ECP project

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CASCADE at PNNL

grizzly at LANL

Thank you  
Questions ?